Second Order Elliptic PDE: Higher Interior Regularity and Boundary Regularity

MATH 612, Texas A&M University

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Higher Interior Regularity

If we assume additional smoothness on the coefficients of our elliptic operator L, and likewise assume f is in a higher regularity space, then we can conclude additional local regularity on our weak solution u.

We'll summarize two theorems along these lines, and then move on to boundary regularity.

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Higher Interior Regularity

Theorem 6.3.2. Suppose $U \subset \mathbb{R}^n$ is open and bounded, a^{ij} , b^i , $c \in C^{m+1}(U)$ ($\forall i, j \in \{1, 2, ..., n\}$) for some $m \in \{0, 1, 2, ...\}$, and L is uniformly elliptic. Also, suppose $f \in H^m(U)$, and $u \in H^1(U)$ (not necessarily $H_0^1(U)$) is a weak solution of

Lu = f in U.

Then $u \in H^{m+2}_{loc}(U)$, and for each $V \subset \subset U$, there exists a constant C, depending only on V, U, m, and the coefficients of L, so that

 $||u||_{H^{m+2}(V)} \leq C(||f||_{H^m(U)} + ||u||_{L^2(U)}).$

Higher Interior Regularity

Notes. 1. The proof proceeds by induction on m, noting that Theorem 6.3.1 is (with minor adjustments) the case m = 0.

2. Since Reg $(H^{m+2}) = m + 2 - \frac{n}{2}$, we can conclude that for each $V \subset \subset U$, $u \in C^{\ell^*,\gamma^*}(\bar{V})$, where we recall that

$$\ell^* = m + 2 - [\frac{n}{2}] - 1,$$

and

$$\gamma^* = \begin{cases} \left[\frac{n}{2}\right] + 1 - \frac{n}{2} & \text{if } \frac{n}{2} \text{ is not an integer} \\ \text{any value } \in (0, 1) & \text{if } \frac{n}{2} \text{ is an integer.} \end{cases}$$

Since $V \subset C$ is arbitrary, $u \in C^{\ell^*}(U)$. For example, if n = 3 and m = 2, then for each $V \subset C$ U, $u \in C^{2,\frac{1}{2}}(\overline{V})$, so u is a classical solution (i.e., $u \in C^2(U)$).

Theorem 6.3.3. Suppose $U \subset \mathbb{R}^n$ is open and bounded, a^{ij} , b^i , $c \in C^{\infty}(U)$ ($\forall i, j \in \{1, 2, ..., n\}$), and L is uniformly elliptic. Also, suppose $f \in C^{\infty}(U)$, and $u \in H^1(U)$ (not necessarily $H^1_0(U)$) is a weak solution of

$$Lu = f$$
 in U .

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Then $u \in C^{\infty}(U)$.

Note. Notice that $f \in C^{\infty}(U) \implies f \in H^{\infty}(W)$ for any $W \subset U$, and this allows us to apply Theorem 6.3.2 (with W replacing U) for all $m \in \{0, 1, 2, ...\}$.

Boundary Regularity

In order to extend interior regularity to the boundary (i.e., to remove the local nature of the last three theorems), we need to assume some smoothness on ∂U .

Theorem 6.3.4. Suppose $U \subset \mathbb{R}^n$ is open and bounded with C^2 boundary, $a^{ij} \in C^1(\overline{U})$, b^i , $c \in L^{\infty}(U)$ ($\forall i, j \in \{1, 2, ..., n\}$), and L is uniformly elliptic. Also, suppose $f \in L^2(U)$, and $u \in H^1_0(U)$ is a weak solution of

$$Lu = f \quad \text{in } U \tag{*}$$
$$u = 0, \quad \text{on } \partial U.$$

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Then $u \in H^2(U)$, and there exists a constant *C*, depending only on *U* and the coefficients of *L*, so that

$$||u||_{H^2(U)} \leq C(||f||_{L^2(U)} + ||u||_{L^2(U)}).$$

Boundary Regularity

Note. If $u \in H_0^1(U)$ is the unique weak solution of (*), then according to Theorem 6.2.6 (boundedness of the resolvent), we have

$$||u||_{L^2(U)} \leq \tilde{C} ||f||_{L^2(U)},$$

for some constant \tilde{C} that depends only on U and the coefficients of L. It follows that there exists a constant $\tilde{\tilde{C}}$, depending only on U and the coefficients of L, so that

$$||u||_{H^2(U)} \leq \tilde{\tilde{C}} ||f||_{L^2(U)}.$$

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1. We'll start by working with a locally flat boundary and show in Step 6 how to map the general case into that setting. We'll work locally near a point $\vec{x_0} \in \partial \mathbb{R}^n_+$, and for notational convenience, we'll shift coordinates so that $\vec{x_0} = 0$.

We set

$$U := B^o(0,1) \cap \mathbb{R}^n_+$$
 (i.e., this is U in the first steps)
 $V := B^o(0, \frac{1}{2}) \cap \mathbb{R}^n_+,$

and we'll introduce a cut-off function $\zeta \in C^\infty_c(\mathbb{R}^n)$ so that

$$\zeta(\vec{x}) = \begin{cases} 1 & \vec{x} \in B(0, \frac{1}{2}) \\ \in [0, 1] & \vec{x} \in B(0, \frac{3}{4}) \setminus B(0, \frac{1}{2}) \\ 0 & \vec{x} \in \mathbb{R}^n \setminus B(0, \frac{3}{4}). \end{cases}$$

2. Since $u \in H_0^1(U)$ is a weak solution of (*), we have

$$B[u,v] = (f,v), \quad \forall v \in H^1_0(U),$$

and precisely as in Step 2 of the proof of Theorem 6.3.1 we can express this equation as

$$\int_{U} \sum_{i,j=1}^{n} a^{ij} u_{x_i} v_{x_j} d\vec{x} = (\tilde{f}, v), \quad \forall v \in H_0^1(U), \quad (**)$$

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where

$$\tilde{f}:=f-\sum_{i=1}^n b^i u_{x_i}-cu\in L^2(U).$$

3. Fix any $k \in \{1, 2, ..., n-1\}$ and any $0 < |h| < \frac{1}{8}$, and as in the proof of Theorem 6.3.1, set

$$v = -D_k^{-h}(\zeta^2 D_k^h u).$$

We would like to substitute v into (**) as in the proof of Theorem 6.3.1, but we need to verify that in the current setting we have $v \in H_0^1(U)$. For this, notice that

$$\begin{split} v(\vec{x}) &= -D_{k}^{-h} \Big(\zeta(\vec{x})^{2} \frac{u(\vec{x}+h\hat{e}_{k})-u(\vec{x})}{h} \Big) \\ &= -\Big\{ \zeta(\vec{x}-h\hat{e}_{k})^{2} \frac{u(\vec{x})-u(\vec{x}-h\hat{e}_{k})}{-h^{2}} - \zeta(\vec{x})^{2} \frac{u(\vec{x}+h\hat{e}_{k})-u(\vec{x})}{-h^{2}} \Big\} \\ &= \frac{1}{h^{2}} \Big\{ \zeta(\vec{x}-h\hat{e}_{k})^{2} (u(\vec{x})-u(\vec{x}-h\hat{e}_{k})) - \zeta(\vec{x})^{2} (u(\vec{x}+h\hat{e}_{k})-u(\vec{x})) \Big\}. \end{split}$$

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Since $k \in \{1, 2, ..., n-1\}$, the points $\vec{x} \pm h\hat{e}_k$ never leave $\overline{\mathbb{R}^n_+}$ (for $\vec{x} \in \overline{\mathbb{R}^n_+}$). Since ζ is 0 in $\mathbb{R}^n \setminus B(0, \frac{3}{4})$, it's clear that $v(\vec{x}) = 0$ for \vec{x} near $\partial B(0, 1) \cap \mathbb{R}^n_+$, and since u = 0 in the trace sense on $B(0, 1) \cap \partial \mathbb{R}^n_+$, we see that v = 0 in the trace sense on ∂U .

Otherwise, v is a sum of terms in $H^1(U)$ (as in the proof of Theorem 6.3.1), so $v \in H^1_0(U)$.

We're now justified in substituting v into (**), and as in the proof of Theorem 6.3.1, we'll denote the left-hand side A and the right-hand side B.

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4. Proceeding as in Steps 4-6 of the proof of Theorem 6.3.1, we obtain the inequality

$$\int_{V} |D_{k}^{h} Du|^{2} d\vec{x} \leq C_{1} \int_{U} f^{2} + u^{2} + |Du|^{2} d\vec{x}, \quad k \in \{1, 2, \dots, n-1\},$$

for some constant C_1 (which was C_8 in the proof of Theorem 6.3.1). From another slight restatement of Theorem 5.8.3 (ii), we can conclude that that

$$u_{x_k} \in H^1(V) \quad \forall k \in \{1, 2, \ldots, n-1\},$$

with the estimate

$$\sum_{\substack{k,l=1\\k+l<2n}}^{n} \|u_{x_kx_l}\|_{L^2(V)} \leq C_2(\|f\|_{L^2(U)} + \|u\|_{H^1(U)}).$$

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In this case, the restatement is because we don't have $V \subset \subset U$.

Recall that in Step 7 of the proof of Theorem 6.3.1 (along with a homework problem), we saw that we can replace $||u||_{H^1(U)}$ on the right-hand side of this last inequality with $||u||_{L^2(U)}$. This gets us to

$$\sum_{\substack{k,l=1\\k+l<2n}}^{n} \|u_{x_kx_l}\|_{L^2(V)} \leq C_3(\|f\|_{L^2(U)} + \|u\|_{L^2(U)}).$$

5. We still need an estimate on $||u_{x_nx_n}||_{L^2(V)}$. For this recall from a note following our statement of Theorem 6.3.1 that interior regularity allows us to work with the strong form of our equation, Lu = f for a.e. $\vec{x} \in U$. To take advantage of this, let's first write our equation in the non-divergence form

$$-\sum_{i,j=1}^{n}a^{ij}u_{x_ix_j}+\sum_{i=1}^{n}\tilde{b}^{i}u_{x_i}+cu=f,\quad \tilde{b}^{i}:=b^{i}-\sum_{j=1}^{n}a^{ij}_{x_j}.$$

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We can now isolate $u_{x_nx_n}$ as

$$a^{nn}u_{x_nx_n} = -\sum_{\substack{i,j=1\i+j<2n}}^n a^{ij}u_{x_ix_j} + \sum_{i=1}^n \tilde{b}^i u_{x_i} + cu - f.$$

Our uniform ellipticity condition is

$$\sum_{i,j=1}^n a^{ij}(\vec{x})\xi_i\xi_j \ge \theta |\vec{\xi}|^2, \quad \forall \, \vec{\xi} \in \mathbb{R}^n,$$

for some $\theta > 0$ and all $\vec{x} \in U$. In particular, if we take $\vec{\xi} = \hat{e}_n$, we see that

$$a^{nn}(\vec{x}) \geq \theta, \quad \forall \, \vec{x} \in U.$$

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This allows us to divide our relation for $u_{x_nx_n}$ by $a^{nn}(\vec{x})$ to obtain an inequality

$$|u_{x_nx_n}| \leq C_4 \Big(\sum_{\substack{i,j=1\\i+j<2n}}^n |u_{x_ix_j}| + |Du| + |u| + |f|\Big),$$

for a.e. $\vec{x} \in U$. If we square this inequality and integrate both sides over V, we obtain

$$\|u_{x_nx_n}\|_{L^2(V)} \leq C_5\Big(\sum_{\substack{i,j=1\\i+j<2n}}^n \|u_{x_ix_j}\|_{L^2(V)} + \|Du\|_{L^2(V)} + \|u\|_{L^2(V)} + \|f\|_{L^2(V)}\Big).$$

Combining this inequality with our previous observations, we see that

$$||u||_{H^{2}(V)} \leq C_{6} \Big(||u||_{L^{2}(U)} + ||f||_{L^{2}(U)} \Big),$$

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for some constant C_6 .

6. In the case of a general C^2 boundary, we fix any $\vec{x_0} \in \partial U$, and we let $\vec{\Phi}(\vec{x})$ denote our usual straightening map, noting that for some r > 0 sufficiently small $\vec{\Phi}$ is a C^2 function on $B^o(\vec{x_0}, r)$ with C^2 inverse $\vec{\Psi}(\vec{y})$. We'll label the range of $\vec{\Phi}$ so that $\vec{\Phi}(\vec{x_0}) = 0$. (I.e., $\vec{y_0} = 0$.)

7. We choose s > 0 sufficiently small so that

$$U':=B^o(0,s)\cap\{y_n>0\}\subset\vec{\Phi}(U\cap B(\vec{x_0},r)),$$

and correspondingly we set

$$V' := B^o(0, \frac{s}{2}) \cap \{y_n > 0\}.$$

We also set

$$u'(\vec{y}) := u(\Psi(\vec{y})), \quad \vec{y} \in U'.$$

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See figure on the next slide.

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We'll check the following claims in the homework:

- 1. $u' \in H^1(U')$
- 2. u' = 0 on $\partial U' \cap \{y_n = 0\}$ in the trace sense
- 8. In the new variables, our elliptic PDE can be expressed as

$$L'u'=f'$$
 in U' ,

where

$$f'(ec{y}) := f(ec{\Psi}(ec{y})),$$

and

$$L'u' := -\sum_{k,l=1}^{n} (a'^{kl}u'_{y_k})_{y_l} + \sum_{k=1}^{n} b'^{k}u'_{y_k} + c'u'.$$

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The coefficients a'^{kl} , b'^{k} , and c' are given below.

The coefficients a'^{kl} , b'^{k} , and c' are obtained directly by expressing the original operator L in the variable $\vec{y} = \vec{\Phi}(\vec{x}) (\iff \vec{x} = \vec{\Psi}(\vec{y}))$. Clearly,

$$c'(ec{y})=c(ec{\Psi}(ec{y})).$$

For the first-order term (in original variables)

$$\sum_{r=1}^n b^r(\vec{x}) u_{x_r}(\vec{x}),$$

we need to carry out a short calculation, and we'll do that on the next slide. For this, we'll denote by Φ^k the k^{th} component of $\vec{\Phi}$.

We have

$$\begin{split} \sum_{r=1}^{n} b^{r}(\vec{\Psi}(\vec{y})) \frac{\partial}{\partial x_{r}} u(\vec{\Psi}(\vec{y})) &= \sum_{r=1}^{n} b^{r}(\vec{\Psi}(\vec{y})) \frac{\partial}{\partial x_{r}} u'(\vec{y}) \\ &= \sum_{r=1}^{n} b^{r}(\vec{\Psi}(\vec{y})) D_{y} u'(\vec{y}) \frac{\partial}{\partial x_{r}} \vec{\Phi}(\vec{x}) \\ &= \sum_{r=1}^{n} b^{r}(\vec{\Psi}(\vec{y})) \sum_{k=1}^{n} \Phi_{x_{r}}^{k}(\vec{\Psi}(\vec{y})) u'_{y_{k}}(\vec{y}) \\ &= \sum_{k=1}^{n} \left\{ \sum_{r=1}^{n} b^{r}(\vec{\Psi}(\vec{y})) \Phi_{x_{r}}^{k}(\vec{\Psi}(\vec{y})) \right\} u'_{y_{k}}(\vec{y}). \end{split}$$

We see that

$$b'^{k}(\vec{y}) = \sum_{r=1}^{n} b^{r}(\vec{\Psi}(\vec{y})) \Phi_{x_{r}}^{k}(\vec{\Psi}(\vec{y})).$$

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Proceeding similarly for the second-order term, we find that

$$a'^{kl}(\vec{y}) = \sum_{r,s=1}^{n} a^{rs}(\vec{\Psi}(\vec{y})) \Phi_{x_r}^k(\vec{\Psi}(\vec{y})) \Phi_{x_s}^l(\vec{\Psi}(\vec{y})).$$

Claim. $u'(\vec{y}) = u(\vec{\Psi}(\vec{y}))$ is a weak solution of

$$L'u' = f'$$
 in U' .

In order to see this, we take any $v' \in H_0^1(U')$ and let B'[u', v'] denote the bilinear form associated with L',

$$B'[u',v'] = \int_{U'} \Big\{ \sum_{k,l=1}^n a'^{kl} u'_{y_k} v'_{y_l} + \sum_{k=1}^n b'^k u'_{y_k} v' + c'u'v' \Big\} d\vec{x}.$$

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Also, we set

$$v(\vec{x}) = v'(\vec{\Phi}(\vec{x})),$$

and observe that similarly as in the homework problem above, $v \in H_0^1(\vec{\Psi}(U'))$. In addition, it will be convenient below to extend v as 0 on $U \setminus \vec{\Psi}(U')$.

This will allow us to express B'[u', v'] in terms of u and v, which is what we do next.

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For the first-order term in B'[u', v'], we can write

$$\begin{split} \sum_{k=1}^{n} b'^{k} u'_{y_{k}} v' &= \sum_{k=1}^{n} b'^{k} (\frac{\partial}{\partial y_{k}} u(\vec{\Psi}(\vec{y}))) v(\vec{\Psi}(\vec{y})) \\ &= \sum_{k=1}^{n} b'^{k} (Du) (\vec{\Psi}(\vec{y})) \vec{\Psi}_{y_{k}}(\vec{y}) v(\vec{\Psi}(\vec{y})) \\ &= \sum_{k=1}^{n} b'^{k} \sum_{i=1}^{n} u_{x_{i}} (\vec{\Psi}(\vec{y})) \Psi_{y_{k}}^{i}(\vec{y}) v(\vec{\Psi}(\vec{y})) \end{split}$$

For notational brevity, we'll write this last expression as

$$\sum_{i=1}^{n} \sum_{k=1}^{n} b'^{k} u_{x_{i}} \Psi_{y_{k}}^{i} v.$$

Proceeding similarly for the other two terms, we obtain the relationship on the next slide.

We have

$$B'[u',v'] = \sum_{i,j=1}^{n} \sum_{k,l=1}^{n} \int_{U'} a'^{kl} u_{x_i} \Psi^{i}_{y_k} v_{x_j} \Psi^{j}_{y_l} d\vec{y} \\ + \sum_{i=1}^{n} \sum_{k=1}^{n} \int_{U'} b'^{k} u_{x_i} \Psi^{i}_{y_k} v d\vec{y} + \int_{U'} c' u v d\vec{y}.$$

According to our definition of a'^{kl} , we can write

$$\sum_{k,l=1}^{n} a'^{kl} \Psi_{y_k}^{i} \Psi_{y_l}^{j} = \sum_{k,l=1}^{n} \sum_{r,s=1}^{n} a^{rs} \Phi_{x_r}^{k} \Phi_{x_s}^{l} \Psi_{y_k}^{i} \Psi_{y_l}^{j}$$
$$= \sum_{k,l=1}^{n} \{ (D\Phi^k) A (D\Phi^l)^T \} \Psi_{y_k}^{i} \Psi_{y_l}^{j}$$
$$= (D\Psi^i) \{ (D\vec{\Phi}) A (D\vec{\Phi})^T \} (D\Psi^j)^T.$$

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Recall from our construction of the maps $\vec{\Phi}$ and $\vec{\Psi}$ last semester that $D\vec{\Psi} = (D\vec{\Phi})^{-1}$ (just differentiate the relation $\vec{\Psi}(\vec{\Phi}(\vec{x})) = \vec{x}$). We see that

$$(D\Psi^i)D\vec{\Phi} = \hat{e}_i^T \quad \forall i \in \{1, 2, \dots, n\}.$$

In this way, we see that

$$\sum_{k,l=1}^n a'^{kl} \Psi^i_{y_k} \Psi^j_{y_l} = \hat{e}_i^T A \hat{e}_j = a^{ij}.$$

Similarly,

$$\sum_{k=1}^{b} b'^{k} \Psi_{y_{k}}^{i} = \sum_{k=1}^{n} \sum_{r=1}^{n} b^{r} \Phi_{x_{r}}^{k} \Psi_{y_{k}}^{i} = \sum_{r=1}^{n} b^{r} \sum_{k=1}^{n} (D\vec{\Psi})_{ik} (D\vec{\Phi})_{kr}$$
$$= \sum_{r=1}^{n} b^{r} \{ (D\vec{\Psi}) (D\vec{\Phi}) \}_{ir} = b^{i}.$$

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Combining these observations, and using the change of variables $\vec{y} = \vec{\Phi}(\vec{x})$ (recall that det $D\vec{\Phi}(\vec{x}) = 1$)

$$B'[u', v'] = \sum_{i,j=1}^{n} \sum_{k,l=1}^{n} \int_{U'} a'^{kl} u_{x_i} \Psi_{y_k}^{i} v_{x_j} \Psi_{y_l}^{j} d\vec{y} + \sum_{i=1}^{n} \sum_{k=1}^{n} \int_{U'} b'^{k} u_{x_i} \Psi_{y_k}^{i} v d\vec{y} + \int_{U'} c' uv d\vec{y} \vec{y} = \vec{\Phi}(\vec{x}) \int_{U} \Big\{ \sum_{i,j=1}^{n} \Big[\sum_{k,l=1}^{n} a'^{kl} \Psi_{y_k}^{i} \Psi_{y_l}^{j} \Big] u_{x_i} v_{x_j} + \sum_{i=1}^{n} \Big[\sum_{k=1}^{n} b'^{k} \Psi_{y_k}^{i} \Big] u_{x_i} v + cuv \Big\} d\vec{x} = \int_{U} \Big\{ \sum_{i,j=1}^{n} a^{ij} u_{x_i} v_{x_j} + \sum_{i=1}^{n} b^{i} u_{x_i} v + cuv \Big\} d\vec{x} = B[u, v].$$

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The latter integrals can be expressed over U because v is taken to be 0 on $U \setminus \Psi(U')$. We see that for all $v' \in H_0^1(U')$,

$$B'[u',v'] = B[u,v] = (f,v) = (f',v'),$$

and so u' is a weak solution of L'u' = f'.

9. We will proceed by applying Steps 1-5 to L', and for this, we need to verify that L' satisfies our assumptions on L in the theorem.

Recalling the relation

$$a'^{kl}(\vec{y}) = \sum_{r,s=1}^{n} a^{rs}(\vec{\Psi}(\vec{y})) \Phi_{x_r}^k(\vec{\Psi}(\vec{y})) \Phi_{x_s}^l(\vec{\Psi}(\vec{y})),$$

we see that our assumption of a C^2 boundary ensures us that $a'^{kl} \in C^1(\overline{U'})$ for all $k, l \in \{1, 2, ..., n\}$. It's clear that $b'^k, c \in L^{\infty}(U')$, for all $k \in \{1, 2, ..., n\}$.

We also need to check that L' is uniformly elliptic in U'. For this, we take any $\vec{y} \in U'$ and any $\vec{\xi} \in \mathbb{R}^n$ and we compute

$$\sum_{k,l=1}^{n} a'^{kl}(\vec{y})\xi_k\xi_l = \sum_{k,l=1}^{n} \sum_{r,s=1}^{n} a^{rs}(\vec{\Psi}(\vec{y}))\Phi_{x_r}^k(\vec{\Psi}(\vec{y}))\Phi_{x_s}^l(\vec{\Psi}(\vec{y}))\xi_k\xi_l$$
$$= \sum_{r,s=1}^{n} a^{rs}(\vec{\Psi}(\vec{y}))\eta_r\eta_s \ge \theta |\vec{\eta}|^2,$$

where we've set (with $\vec{\xi}$ viewed as a row vector)

$$ec{\eta}(ec{y}):=ec{\xi}Dec{\Phi}(ec{\Psi}(ec{y}))\implies ec{\xi}=ec{\eta}(ec{y})(Dec{\Phi})^{-1}=ec{\eta}(ec{y})Dec{\Psi}(ec{y}).$$

We see that

$$|\vec{\xi}| \leq C_7 |\vec{\eta}(\vec{y})|,$$

for some constant C_7 .

In particular,

$$\theta |\vec{\eta}|^2 \ge \frac{\theta}{C_7^2} |\vec{\xi}|^2,$$

and this gives uniform ellipticity with constant θ/C_7^2 .

10. We are now justified in applying Steps 1-5 to u' as a solution of L'u' = f', and this provides the inequality

$$||u'||_{H^2(V')} \leq C_8(||f'||_{L^2(U')} + ||u'||_{L^2(U')}).$$

Returning to original variables, we can express this as

$$||u||_{H^2(V)} \leq C_9(||f||_{L^2(U)} + ||u||_{L^2(U)}),$$

where $V = \vec{\Psi}(V')$ and on the right-hand side we've extended the domain of integration from $\vec{\Psi}(U')$ to U.

We now put these local estimates together in the usual way. Since ∂U is compact, we can find finitely many sets $\{V_i\}_{i=1}^N$ as described above, along with one additional set $V_0 \subset \subset U$ so that

$$U=\bigcup_{i=0}^N V_i.$$



From Steps 1-9 of the current proof, we have

$$||u||_{H^2(V_i)} \le K_i(||f||_{L^2(U)} + ||u||_{L^2(U)}), \quad i = 1, 2, \dots N,$$

for some constants $\{K_i\}_{i=1}^N$, while from Theorem 6.3.1 we have

$$||u||_{H^2(V_0)} \le K_0(||f||_{L^2(U)} + ||u||_{L^2(U)}),$$

for some constant K_0 . Finally,

$$||u||^2_{H^2(U)} \le \sum_{i=0}^N ||u||^2_{H^2(V_i)}$$

(inequality because of overlap in the sets), so for some constant C we have the claimed inequality

$$||u||_{H^2(U)} \leq C(||f||_{L^2(U)} + ||u||_{L^2(U)}).$$

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Higher Boundary Regularity

We'll conclude by summarizing two results on higher boundary regularity.

Theorem 6.3.5. Suppose $U \subset \mathbb{R}^n$ is open and bounded with C^{m+2} boundary, a^{ij} , b^i , $c \in C^{m+1}(\overline{U})$ ($\forall i, j \in \{1, 2, ..., n\}$) for some $m \in \{0, 1, 2, ...\}$, and L is uniformly elliptic. Also, suppose $f \in H^m(U)$, and $u \in H^1_0(U)$ is a weak solution of

 $Lu = f \quad \text{in } U$ $u = 0, \quad \text{on } \partial U.$

Then $u \in H^{m+2}(U)$, and there exists a constant *C*, depending only on *m*, *U* and the coefficients of *L*, so that

$$||u||_{H^{m+2}(U)} \leq C(||f||_{H^m(U)} + ||u||_{L^2(U)}).$$

Note. If Reg $(H^{m+2}) > 0$, then *u* is continuous up to the boundary.

Higher Boundary Regularity

Theorem 6.3.6. Suppose $U \subset \mathbb{R}^n$ is open and bounded with C^{∞} boundary, a^{ij} , b^i , $c \in C^{\infty}(\overline{U})$ ($\forall i, j \in \{1, 2, ..., n\}$), and L is uniformly elliptic. Also, suppose $f \in C^{\infty}(\overline{U})$, and $u \in H_0^1(U)$ is a weak solution of

 $Lu = f \quad \text{in } U$ $u = 0, \quad \text{on } \partial U.$

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Then $u \in C^{\infty}(\overline{U})$.