# Second Order Parabolic PDE: Weak Solutions and Galerkin Approximations 

MATH 612, Texas A\&M University

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## Second Order Parabolic PDE

Let $U \subset \mathbb{R}^{n}$ be open and bounded, and for $T>0$ set

$$
U_{T}:=U \times(0, T] .
$$

We'll consider equations of the form

$$
\begin{align*}
& u_{t}+L u=f ;  \tag{P}\\
& \text { in } U_{T} \\
& u=0 ; \\
& \text { on } \partial U \times[0, T] \\
& u=g ; \\
& \text { on } U \times\{t=0\},
\end{align*}
$$

where $L$ denotes a partial differential operator either in divergence form

$$
L u:=-\sum_{i, j=1}^{n}\left(a^{i j}(\vec{x}, t) u_{x_{i}}\right)_{x_{j}}+\sum_{i=1}^{n} b^{i}(\vec{x}, t) u_{x_{i}}+c(\vec{x}, t) u
$$

or non-divergence form

$$
L u:=-\sum_{i, j=1}^{n} a^{i j}(\vec{x}, t) u_{x_{i} x_{j}}+\sum_{i=1}^{n} b^{i}(\vec{x}, t) u_{x_{i}}+c(\vec{x}, t) u .
$$

## Uniform Parabolicity

Definition. For either form of the operator $L$, we say that the partial differential operator $\partial_{t}+L$ is uniformly parabolic in $U_{T}$ if there exists a constant $\theta>0$ so that

$$
\sum_{i, j=1}^{n} a^{i j}(\vec{x}, t) \xi_{i} \xi_{j} \geq \theta|\vec{\xi}|^{2}
$$

for a.e. $(\vec{x}, t) \in U_{T}$ and all $\vec{\xi} \in \mathbb{R}^{n}$. I.e., if we denote $A(\vec{x}, t)=\left(a^{i j}(\vec{x}, t)\right)$, then

$$
\vec{\xi}^{T} A(\vec{x}, t) \vec{\xi} \geq \theta|\vec{\xi}|^{2}
$$

for a.e. $(\vec{x}, t) \in U_{T}$ and all $\vec{\xi} \in \mathbb{R}^{n}$.
Notes. 1. For the heat equation, $L=-\Delta$, so that $A(\vec{x}, t)=I$, and we have

$$
\vec{\xi}^{T} A(\vec{x}, t) \vec{\xi}=|\vec{\xi}|^{2} .
$$

I.e., $\theta=1$.

## Uniform Parabolicity

2. More generally, if $A(\vec{x}, t)$ is symmetric, then by the min-max principle, this condition holds if and only if the eigenvalues of $A(\vec{x}, t)$ are all bounded below by $\theta$. Recall from last semester that a partial differential operator $\partial_{t}+L$ is parabolic if the matrix associated with its second order terms (including $t$ ) has 0 as one of its eigenvalues and its other eigenvalues all have the same sign. So an operator that is uniformly parabolic in $U_{T}$ is certainly parabolic in $U_{T}$.

## Basic Assumptions

In order to avoid repeated statements of assumptions, we'll collect our basic assumptions for this section here. These will be:

1. $U \subset \mathbb{R}^{n}$ is open and bounded, and $T>0$;
2. $a^{i j}=a^{i i}$ for all $i, j \in\{1,2, \ldots, n\}$;
3. $a^{i j}, b^{i}, c \in L^{\infty}\left(U_{T}\right)$ for all $i, j \in\{1,2, \ldots, n\}$;
4. The operator $L$ is in divergence form and uniformly parabolic in $U_{T}$;
5. $f \in L^{2}\left(U_{T}\right), g \in L^{2}(U)$.

We'll refer to this collection of assumptions as Assumptions (A).
Note. In this section, unless explicitly stated otherwise, $(\cdot, \cdot)$ will denote $L^{2}(U)$ inner product, and $\langle\cdot, \cdot\rangle$ will denote the action of an element of $H^{-1}(U)$ on an element of $H_{0}^{1}(U)$.

## The Time-Dependent Bilinear Form

Formally, if $u$ is a smooth solution of

$$
u_{t}+L u=f
$$

with $L$ in divergence form, we can multiply by a test function $\phi \in C_{c}^{\infty}(U)$ and proceed as in our section on elliptic operators to obtain the relation

$$
\left(u_{t}, \phi\right)+B[u, \phi ; t]=(f, \phi),
$$

for each $t \in[0, T]$, where
$B[u, \phi ; t]:=\int_{U}\left\{\sum_{i, j=1}^{n} a^{i j}(\vec{x}, t) u_{x_{i}} \phi_{x_{j}}+\sum_{i=1}^{n} b^{i}(\vec{x}, t) u_{x_{i}} \phi+c(\vec{x}, t) u \phi\right\} d \vec{x}$.
We'll use this to develop our weak formulation for $(\mathcal{P})$.

## Notation

If $u(\cdot, t) \in H_{0}^{1}(U)$ for a.e. $t \in(0, T)$, then we'll regard $u$ as a map $t \mapsto u(\cdot, t)$. In this section, following the convention that Evans adopts, we'll denote such maps with a bold $\mathbf{u}$. I.e.,

$$
\mathbf{u}(t)(\vec{x})=u(\vec{x}, t)
$$

and similarly for $\mathbf{f}(t)$. This allows us to express the strong form of $(\mathcal{P})$ as

$$
\mathbf{u}^{\prime}+L \mathbf{u}=\mathbf{f}, \quad \text { for a.e. } t \in(0, T)
$$

Here, we notice that since $f \in L^{2}\left(U_{T}\right)$, we have
$\|\mathbf{f}\|_{L^{2}\left(0, T ; L^{2}(U)\right)}^{2}=\int_{0}^{T}\|\mathbf{f}(t)\|_{L^{2}(U)}^{2} d t=\int_{0}^{T} \int_{U}|f(\vec{x}, t)|^{2} d \vec{x} d t<\infty$.
I.e., $f \in L^{2}\left(0, T ; L^{2}(U)\right)$.

Motivating the Weak Formulation
The associated weak form of this equation can be expressed as

$$
\left(\mathbf{u}^{\prime}, v\right)+B[\mathbf{u}, v ; t]=(\mathbf{f}, v),
$$

for all $v \in H_{0}^{1}(U)$ and a.e. $t \in(0, T)$. We see that

$$
\begin{aligned}
\left(\mathbf{u}^{\prime},\right. & v) \\
= & =-B[\mathbf{u}, v ; t]+(\mathbf{f}, v) \\
= & -\int_{U}\left\{\sum_{i, j=1}^{n} a^{i j} \mathbf{u}_{x_{i}} v_{x_{j}}+\sum_{i=1}^{n} b^{i} \mathbf{u}_{x_{i}} v+c \mathbf{u} v\right\}+(\mathbf{f}, v) \\
= & \int_{U}\left\{\left(\mathbf{f}-\sum_{i=1}^{n} b^{i} \mathbf{u}_{x_{i}}-c \mathbf{u}\right) v-\sum_{j=1}^{n}\left(\sum_{i=1}^{n} a^{i j} \mathbf{u}_{x_{i}}\right) v_{x_{j}}\right\} d \vec{x} \\
= & \int_{U}\left\{\mathbf{g}^{0} v+\sum_{j=1}^{n} \mathbf{g}^{j} v_{x_{j}}\right\} d \vec{x} .
\end{aligned}
$$

Motivating the Weak Formulation
I.e.,

$$
\left(\mathbf{u}^{\prime}, v\right)=\int_{U}\left\{\mathbf{g}^{0} v+\sum_{j=1}^{n} \mathbf{g}^{j} v_{x_{j}}\right\} d \vec{x}
$$

where for $\mathbf{u}(t) \in H_{0}^{1}(U)$, we have

$$
\begin{aligned}
& \mathbf{g}^{0}(t)=\mathbf{f}(t)-\sum_{i=1}^{n} b^{i}(\vec{x}, t) \mathbf{u}(t)_{x_{i}}-c(\vec{x}, t) \mathbf{u}(t) \in L^{2}(U) \\
& \mathbf{g}^{j}(t)=-\sum_{i=1}^{n} a^{i j}(\vec{x}, t) \mathbf{u}_{x_{i}} \in L^{2}(U)
\end{aligned}
$$

Motivating the Weak Formulation
If we compare with Theorem 5.9.1, we see that this suggests that we should have $\mathbf{u}^{\prime}(t) \in H^{-1}(U)$ for a.e. $t \in(0, T)$, with

$$
\begin{aligned}
\left\|\mathbf{u}^{\prime}(t)\right\|_{H^{-1}(U)} & \leq\left(\int_{U} \sum_{j=0}^{n}\left|\mathbf{g}^{j}(t)\right|^{2} d \vec{x}\right)^{1 / 2} \\
& \leq C\left(\|\mathbf{u}(t)\|_{H^{1}(U)}+\|\mathbf{f}(t)\|_{L^{2}(U)}\right)
\end{aligned}
$$

This motivates our weak formulation of $(\mathcal{P})$.

The Parabolic Weak Formulation
Definition. We say that a function $\mathbf{u} \in L^{2}\left(0, T ; H_{0}^{1}(U)\right)$, with $\mathbf{u}^{\prime} \in L^{2}\left(0, T ; H^{-1}(U)\right)$, is a weak solution of $(\mathcal{P})$ if
(i) $\left\langle\mathbf{u}^{\prime}, v\right\rangle+B[\mathbf{u}, v ; t]=(\mathbf{f}, v)$ for all $v \in H_{0}^{1}(U)$ and a.e. $t \in(0, T)$, and
(ii) $\mathbf{u}(0)=g$.

Note. We know from Theorem 5.9.3 that under these assumptions, we have $\mathbf{u} \in C\left([0, T] ; L^{2}(U)\right)$, so the pointwise evaluation $\mathbf{u}(0)=g$ is justified.

## Galerkin Approximations

We'll approach existence by first constructing solutions to certain finite-dimensional approximations of the weak formulation of $(\mathcal{P})$, and then taking an appropriate limit. This method is named after the Russian mathematician Boris Galerkin (1871-1945).

To begin, let $\left\{w_{k}\right\}_{k=1}^{\infty}$ denote an orthogonal basis of $H_{0}^{1}(U)$ that is also an orthonormal basis of $L^{2}(U)$. For example, such a basis is constructed in the proof of Theorem 6.5.2 in Evans as eigenfunctions for the Laplacian operator $L=-\Delta$ on $U$.

We fix any $m \in\{1,2, \ldots\}$, and look for solutions of the weak formulation of $(\mathcal{P})$ of the form

$$
\mathbf{u}_{m}(t)=\sum_{k=1}^{m} d_{m}^{k}(t) w_{k},
$$

where the coefficient functions $\left\{d_{m}^{k}(t)\right\}_{k=1}^{m}$ are to be determined.

## Galerkin Approximations

If we substitute this ansatz into the weak formulation of $(\mathcal{P})$, we obtain the relation
$\left\langle\sum_{k=1}^{m} d_{m}^{k \prime}(t) w_{k}, v\right\rangle+B\left[\sum_{k=1}^{m} d_{m}^{k}(t) w_{k}, v ; t\right]=(\mathbf{f}(t), v), \quad \forall v \in H_{0}^{1}(U)$,
and using linearity
$\sum_{k=1}^{m} d_{m}^{k \prime}(t)\left\langle w_{k}, v\right\rangle+\sum_{k=1}^{m} d_{m}^{k}(t) B\left[w_{k}, v ; t\right]=(f(t), v), \quad \forall v \in H_{0}^{1}(U)$.
Here, we have $w_{k} \in H_{0}^{1}(U)$ for each fixed $k \in\{1,2, \ldots, m\}$, so we can replace $\left\langle w_{k}, v\right\rangle$ with $\left(w_{k}, v\right)_{L^{2}(U)}$.

## Galerkin Approximations

In fact, since $\mathbf{u}_{m}(t)$ is a finite-dimensional approximation of the solution, we expect that this is too much to ask, but we can think of replacing the requirement that this be true for all $v \in H_{0}^{1}(U)$ with the requirement that it be true for all $v \in \operatorname{Span}\left\{w_{j}\right\}_{j=1}^{m}$. I.e., we require

$$
\sum_{k=1}^{m} d_{m}^{k \prime}(t)\left(w_{k}, w_{j}\right)+\sum_{k=1}^{m} d_{m}^{k}(t) B\left[w_{k}, w_{j} ; t\right]=\left(\mathbf{f}(t), w_{j}\right)
$$

for all $j \in\{1,2, \ldots, m\}$. We set

$$
e^{j k}(t):=B\left[w_{k}, w_{j} ; t\right] \quad \text { and } \quad f^{j}(t):=\left(\mathbf{f}(t), w_{j}\right),
$$

and use orthonormality of $\left\{w_{k}\right\}_{k=1}^{n}$ to obtain the first-order system of ODE

$$
d_{m}^{j \prime}(t)=-\sum_{k=1}^{m} e^{j k}(t) d_{m}^{k}(t)+f^{j}(t), \quad j=1,2, \ldots m
$$

## Galerkin Approximations

For initial values, we would like to set

$$
\mathbf{u}_{m}(0)=g \in L^{2}(U)
$$

but again this is too much to ask. Instead, we formally determine the coefficients $\left\{d_{m}^{k}(0)\right\}_{k=1}^{m}$ that we would need in order to have the relation

$$
g=\sum_{k=1}^{\infty} d_{m}^{k}(0) w_{k}
$$

Again using orthonormality of $\left\{w_{k}\right\}_{k=1}^{\infty}$, we see that

$$
d_{m}^{j}(0)=\left(g, w_{j}\right),
$$

for each $j \in\{1,2, \ldots, m\}$.

## Galerkin Approximations

We now want to assert something about solvability for the ODE system

$$
\begin{aligned}
& d_{m}^{j \prime}(t)=-\sum_{k=1}^{m} e^{j k}(t) d_{m}^{k}(t)+f^{j}(t), \quad j=1,2, \ldots m \\
& d_{m}^{j}(0)=\left(g, w_{j}\right)
\end{aligned}
$$

and for this we need to better understand the nature of the coefficents $\left\{f^{j}(t)\right\}_{j=1}^{m}$ and $\left\{e^{j k}(t)\right\}_{j=1}^{m}$. First,

$$
\begin{aligned}
\left\|f^{j}\right\|_{L^{2}(0, T)}^{2} & =\int_{0}^{T}\left|\left(\mathbf{f}(t), w_{j}\right)\right|^{2} d t \stackrel{\text { c.s. }}{\leq} \int_{0}^{T}\|\mathbf{f}(t)\|_{L^{2}(U)}^{2}\left\|w_{j}\right\|_{L^{2}(U)}^{2} d t \\
& =\|\mathbf{f}\|_{L^{2}\left(0, T ; L^{2}(U)\right)}<\infty
\end{aligned}
$$

so $f^{j} \in L^{2}(0, T ; \mathbb{R})$ for all $j \in\{1,2, \ldots, m\}$.

## Galerkin Approximations

Likewise, for each $e^{j k}$,

$$
\left\|e^{j k}\right\|_{L^{\infty}(0, T)}=\left\|B\left[w_{k}, w_{j} ; t\right]\right\|_{L^{\infty}(0, T)}
$$

and

$$
\begin{aligned}
& \left|B\left[w_{k}, w_{j} ; t\right]\right|=\left|\int_{U}\left\{\sum_{i, l=1}^{n} a^{i l}\left(w_{k}\right)_{x_{j}}\left(w_{j}\right)_{x_{l}}+\sum_{i=1}^{n} b^{i}\left(w_{k}\right)_{x_{i}} w_{j}+c w_{k} w_{j}\right\} d \vec{x}\right| \\
& \leq C \int_{U}\left\{\sum_{i, l=1}^{n}\left|\left(w_{k}\right)_{x_{j}}\right|\left|\left(w_{j}\right)_{x_{l}}\right|+\sum_{i=1}^{n}\left|( w _ { k } ) _ { x _ { i } } \left\|w_{j}\left|+\left|w_{k} \| w_{j}\right|\right\} d \vec{x}\right.\right.\right. \\
& \leq \tilde{C}\left\|w_{k}\right\|_{H^{1}(U)}\left\|w_{j}\right\|_{H^{1}(U)}<\infty,
\end{aligned}
$$

so $e^{j k} \in L^{\infty}(0, T ; \mathbb{R})$ for all $j, k \in\{1,2, \ldots, m\}$.

## Galerkin Approximations

In summary, for the ODE system

$$
\begin{aligned}
& d_{m}^{j \prime}(t)=-\sum_{k=1}^{m} e^{j k}(t) d_{m}^{k}(t)+f^{j}(t), \quad j=1,2, \ldots m \\
& d_{m}^{j}(0)=\left(g, w_{j}\right)
\end{aligned}
$$

we have $f^{j} \in L^{2}(0, T ; \mathbb{R})$ and $e^{j k} \in L^{\infty}(0, T ; \mathbb{R})$ for all
$j, k \in\{1,2, \ldots, m\}$, which is more than we need for what we want to assert. Proceeding similarly as we did first semester, we can show that as long as $f^{j}, e^{j k} \in L^{1}(0, T ; \mathbb{R})$, then there exists a unique absolutely continuous solution

$$
\mathbf{d}_{m}(t)=\left(d_{m}^{1}(t), d_{m}^{2}(t), \ldots, d_{m}^{m}(t)\right)
$$

to this system on $[0, T]$.
For details on this existence, see, e.g., Theorem 2.1 in "Spectral Theory of Ordinary Differential Operators," by Joachim Weidmann, Lecture Notes in Mathematics 1258 (1987).

## Galerkin Approximations

By construction, we see that

$$
\begin{equation*}
\mathbf{u}_{m}(t)=\sum_{k=1}^{m} d_{m}^{k}(t) w_{k} \tag{G}
\end{equation*}
$$

solves the finite-dimensional weak problem,

$$
\begin{equation*}
\left\langle\mathbf{u}_{m}^{\prime}, w_{j}\right\rangle+B\left[\mathbf{u}_{m}, w_{j} ; t\right]=\left(\mathbf{f}, w_{j}\right), \quad \forall j \in\{1,2, \ldots, m\} \tag{FDW}
\end{equation*}
$$

for a.e. $t \in(0, T)$, along with

$$
\begin{equation*}
d_{m}^{k}(0)=\left(g, w_{k}\right), \quad \forall k \in\{1,2, \ldots, m\} \tag{IC}
\end{equation*}
$$

In this way, we have established the following theorem from Evans:
Theorem 7.1.1. Let Assumptions (A) hold. Then for each integer $m \in\{1,2, \ldots\}$, there exists a unique function $\mathbf{u}_{m}$ of the form (G) solving (FDW), (IC).

## Galerkin Approximations

We refer to $\mathbf{u}_{m}(t)$ constructed in this way as the Galerkin approximation of our sought solution $\mathbf{u}(t)$. Our goal will be to show that as $m \rightarrow \infty$, the sequence $\left\{\mathbf{u}_{m}\right\}_{m=1}^{\infty}$ converges to a solution of the weak formulation of $(\mathcal{P})$ in an appropriate sense.

