

MATH 311

Topics in Applied Mathematics

**Lecture 25:**  
**Bessel functions (continued).**

**Bessel's differential equation** of order  $m \geq 0$ :

$$z^2 \frac{d^2 f}{dz^2} + z \frac{df}{dz} + (z^2 - m^2)f = 0$$

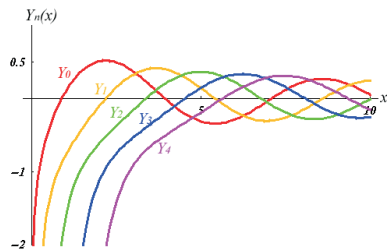
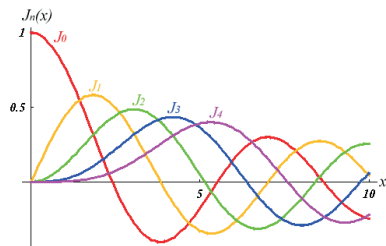
The equation is considered on the interval  $(0, \infty)$ .  
Solutions are called **Bessel functions** of order  $m$ .

$J_m(z)$ : **Bessel function of the first kind,**

$Y_m(z)$ : **Bessel function of the second kind.**

The general Bessel function of order  $m$  is  
 $f(z) = c_1 J_m(z) + c_2 Y_m(z)$ , where  $c_1, c_2$  are  
constants.

# Bessel functions of the 1st and 2nd kind



## Asymptotics at the origin

$J_m(z)$  is regular while  $Y_m(z)$  has a singularity at 0.

As  $z \rightarrow 0$ , we have for any integer  $m > 0$

$$J_m(z) \sim \frac{1}{2^m m!} z^m, \quad Y_m(z) \sim -\frac{2^m (m-1)!}{\pi} z^{-m}.$$

Also,  $J_0(z) \sim 1$ ,  $Y_0(z) \sim \frac{2}{\pi} \log z$ .

To get the asymptotics for a noninteger  $m$ , we replace  $m!$  by  $\Gamma(m+1)$  and  $(m-1)!$  by  $\Gamma(m)$ .

$J_m(z)$  is uniquely determined by this asymptotics while  $Y_m(z)$  is not.

## Asymptotics at infinity

As  $z \rightarrow \infty$ , we have

$$J_m(z) = \sqrt{\frac{2}{\pi z}} \cos \left( z - \frac{\pi}{4} - \frac{m\pi}{2} \right) + O(z^{-1}),$$

$$Y_m(z) = \sqrt{\frac{2}{\pi z}} \sin \left( z - \frac{\pi}{4} - \frac{m\pi}{2} \right) + O(z^{-1}).$$

Both  $J_m(z)$  and  $Y_m(z)$  are uniquely determined by this asymptotics.

For  $m = 1/2$ , these are exact formulas.

Original definition by Bessel (only for integer  $m$ ):

$$J_m(z) = \frac{1}{\pi} \int_0^\pi \cos(z \sin \tau - m\tau) d\tau$$
$$= \frac{1}{\pi} \int_0^\pi \cos(z \sin \tau) \cos(m\tau) d\tau + \frac{1}{\pi} \int_0^\pi \sin(z \sin \tau) \sin(m\tau) d\tau.$$

The first integral is 0 for any odd  $m$  while the second integral is 0 for any even  $m$ . It follows that

$$\cos(z \sin \tau) = J_0(z) + 2 \sum_{n=1}^{\infty} J_{2n}(z) \cos(2n\tau),$$

$$\sin(z \sin \tau) = 2 \sum_{n=1}^{\infty} J_{2n-1}(z) \sin((2n-1)\tau).$$

## Zeros of Bessel functions

Let  $0 < j_{m,1} < j_{m,2} < \dots$  be zeros of  $J_m(z)$  and  $0 < y_{m,1} < y_{m,2} < \dots$  be zeros of  $Y_m(z)$ .

Let  $0 \leq j'_{m,1} < j'_{m,2} < \dots$  be zeros of  $J'_m(z)$  and  $0 < y'_{m,1} < y'_{m,2} < \dots$  be zeros of  $Y'_m(z)$ .

(We let  $j'_{0,1} = 0$  while  $j'_{m,1} > 0$  if  $m > 0$ .)

Then the zeros are interlaced:

$$\begin{aligned} m \leq j'_{m,1} < y_{m,1} < y'_{m,1} < j_{m,1} < \\ < j'_{m,2} < y_{m,2} < y'_{m,2} < j_{m,2} < \dots \end{aligned}$$

Asymptotics of the  $n$ th zeros as  $n \rightarrow \infty$ :

$$j'_{m,n} \approx y_{m,n} \sim \left(n + \frac{1}{2}m - \frac{3}{4}\right)\pi,$$

$$y'_{m,n} \approx j_{m,n} \sim \left(n + \frac{1}{2}m - \frac{1}{4}\right)\pi.$$

## Dirichlet Laplacian in a circle

Eigenvalue problem:

$$\nabla^2 \phi + \lambda \phi = 0 \quad \text{in } D = \{(x, y) : x^2 + y^2 \leq R^2\},$$
$$\phi|_{\partial D} = 0.$$

Separation of variables in polar coordinates:

$\phi(r, \theta) = f(r)h(\theta)$ . Reduces the problem to two one-dimensional eigenvalue problems:

$$r^2 f'' + r f' + (\lambda r^2 - \mu) f = 0, \quad f(R) = 0, \quad |f(0)| < \infty;$$

$$h'' = -\mu h, \quad h(-\pi) = h(\pi), \quad h'(-\pi) = h'(\pi).$$

The latter problem has eigenvalues  $\mu_m = m^2$ ,

$m = 0, 1, 2, \dots$ , and eigenfunctions  $h_0 = 1$ ,

$h_m(\theta) = \cos m\theta$ ,  $\tilde{h}_m(\theta) = \sin m\theta$ ,  $m \geq 1$ .



The 1st intermediate eigenvalue problem:

$$r^2 f'' + r f' + (\lambda r^2 - m^2) f = 0, \quad f(R) = 0, \quad |f(0)| < \infty.$$

New variable  $z = \sqrt{\lambda} \cdot r$  reduces the equation to Bessel's equation of order  $m$ . Hence the general solution is  $f(r) = c_1 J_m(\sqrt{\lambda} r) + c_2 Y_m(\sqrt{\lambda} r)$ , where  $c_1, c_2$  are constants.

Singular condition  $|f(0)| < \infty$  holds if  $c_2 = 0$ .

Nonzero solution exists if  $J_m(\sqrt{\lambda} R) = 0$ .

Thus there are infinitely many eigenvalues  $\lambda_{m,1}, \lambda_{m,2}, \dots$ , where  $\sqrt{\lambda_{m,n}} R = j_{m,n}$ , i.e.,  $\lambda_{m,n} = (j_{m,n}/R)^2$ .  
Associated eigenfunctions:  $f_{m,n}(r) = J_m(j_{m,n} r/R)$ .

The eigenfunctions  $f_{m,n}(r) = J_m(j_{m,n} r/R)$  are orthogonal relative to the inner product

$$\langle f, g \rangle_r = \int_0^R f(r) g(r) r dr.$$

Any piecewise continuous function  $g$  on  $[0, R]$  is expanded into a **Fourier-Bessel series**

$$g(r) = \sum_{n=1}^{\infty} c_n J_m\left(j_{m,n} \frac{r}{R}\right), \quad c_n = \frac{\langle g, f_{m,n} \rangle_r}{\langle f_{m,n}, f_{m,n} \rangle_r},$$

that converges in the mean (with weight  $r$ ).

If  $g$  is piecewise smooth, then the series converges at its points of continuity.

## Eigenvalue problem:

$$\nabla^2 \phi + \lambda \phi = 0 \quad \text{in } D = \{(x, y) : x^2 + y^2 \leq R^2\},$$
$$\phi|_{\partial D} = 0.$$

**Eigenvalues:**  $\lambda_{m,n} = (j_{m,n}/R)^2$ , where  $m = 0, 1, 2, \dots$ ,  $n = 1, 2, \dots$ , and  $j_{m,n}$  is the  $n$ th positive zero of the Bessel function  $J_m$ .

**Eigenfunctions:**  $\phi_{0,n}(r, \theta) = J_0(j_{0,n} r/R)$ .

For  $m \geq 1$ ,  $\phi_{m,n}(r, \theta) = J_m(j_{m,n} r/R) \cos m\theta$  and  $\tilde{\phi}_{m,n}(r, \theta) = J_m(j_{m,n} r/R) \sin m\theta$ .

## Neumann Laplacian in a circle

Eigenvalue problem:

$$\nabla^2 \phi + \lambda \phi = 0 \quad \text{in } D = \{(x, y) : x^2 + y^2 \leq R^2\},$$
$$\frac{\partial \phi}{\partial n} \Big|_{\partial D} = 0.$$

Again, separation of variables in polar coordinates,  $\phi(r, \theta) = f(r)h(\theta)$ , reduces the problem to two one-dimensional eigenvalue problems:

$$r^2 f'' + r f' + (\lambda r^2 - \mu) f = 0, \quad f'(R) = 0, \quad |f(0)| < \infty;$$

$$h'' = -\mu h, \quad h(-\pi) = h(\pi), \quad h'(-\pi) = h'(\pi).$$

The 2nd problem has eigenvalues  $\mu_m = m^2$ ,  $m = 0, 1, 2, \dots$ , and eigenfunctions  $h_0 = 1$ ,  $h_m(\theta) = \cos m\theta$ ,  $\tilde{h}_m(\theta) = \sin m\theta$ ,  $m \geq 1$ .

The 1st one-dimensional eigenvalue problem:

$$r^2 f'' + r f' + (\lambda r^2 - m^2) f = 0, \quad f'(R) = 0, \quad |f(0)| < \infty.$$

For  $\lambda > 0$ , the general solution of the equation is  $f(r) = c_1 J_m(\sqrt{\lambda} r) + c_2 Y_m(\sqrt{\lambda} r)$ , where  $c_1, c_2$  are constants.

Singular condition  $|f(0)| < \infty$  holds if  $c_2 = 0$ .

Nonzero solution exists if  $J'_m(\sqrt{\lambda} R) = 0$ .

Thus there are infinitely many eigenvalues  $\lambda_{m,1}, \lambda_{m,2}, \dots$ ,

where  $\sqrt{\lambda_{m,n}} R = j'_{m,n}$ , i.e.,  $\lambda_{m,n} = (j'_{m,n}/R)^2$ .

Associated eigenfunctions:  $f_{m,n}(r) = J_m(j'_{m,n} r/R)$ .

$\lambda = 0$  is an eigenvalue only for  $m = 0$ .

## Eigenvalue problem:

$$\nabla^2 \phi + \lambda \phi = 0 \quad \text{in } D = \{(x, y) : x^2 + y^2 \leq R^2\},$$
$$\left. \frac{\partial \phi}{\partial n} \right|_{\partial D} = 0.$$

**Eigenvalues:**  $\lambda_{m,n} = (j'_{m,n}/R)^2$ , where  $m = 0, 1, 2, \dots$ ,  $n = 1, 2, \dots$ , and  $j'_{m,n}$  is the  $n$ th positive zero of  $J'_m$  (exception:  $j'_{0,1} = 0$ ).

**Eigenfunctions:**  $\phi_{0,n}(r, \theta) = J_0(j'_{0,n} r/R)$ .

In particular,  $\phi_{0,1} = 1$ .

For  $m \geq 1$ ,  $\phi_{m,n}(r, \theta) = J_m(j'_{m,n} r/R) \cos m\theta$  and  $\tilde{\phi}_{m,n}(r, \theta) = J_m(j'_{m,n} r/R) \sin m\theta$ .

## Laplacian in a circular sector

Eigenvalue problem:

$$\nabla^2 \phi + \lambda \phi = 0 \quad \text{in } D = \{(r, \theta) : r < R, 0 < \theta < L\},$$
$$\phi|_{\partial D} = 0.$$

Again, separation of variables in polar coordinates,  $\phi(r, \theta) = f(r)h(\theta)$ , reduces the problem to two one-dimensional eigenvalue problems:

$$r^2 f'' + r f' + (\lambda r^2 - \mu) f = 0, \quad f(0) = f(R) = 0;$$
$$h'' = -\mu h, \quad h(0) = h(L) = 0.$$

The 2nd problem has eigenvalues  $\mu_m = \left(\frac{m\pi}{L}\right)^2$ ,  $m = 1, 2, \dots$ , and eigenfunctions  $h_m(\theta) = \sin \frac{m\pi\theta}{L}$ .

The 1st one-dimensional eigenvalue problem:

$$r^2 f'' + r f' + (\lambda r^2 - \nu^2) f = 0, \quad f(0) = f(R) = 0.$$

Here  $\nu^2 = \mu_m$ . We may assume that  $\lambda > 0$ .

The general solution of the equation is

$f(r) = c_1 J_\nu(\sqrt{\lambda} r) + c_2 Y_\nu(\sqrt{\lambda} r)$ , where  $c_1, c_2$  are constants.

Boundary condition  $f(0) = 0$  holds if  $c_2 = 0$ .

Nonzero solution exists if  $J_\nu(\sqrt{\lambda} R) = 0$ .

Thus there are infinitely many eigenvalues  $\lambda_{m,1}, \lambda_{m,2}, \dots$ , where  $\sqrt{\lambda_{m,n}} R = j_{\nu,n}$ , i.e.,  $\lambda_{m,n} = (j_{\nu,n}/R)^2$ .

Associated eigenfunctions:  $f_{m,n}(r) = J_\nu(j_{\nu,n} r/R)$ .

Note that  $\nu = m\pi/L$ .



## Eigenvalue problem:

$$\nabla^2 \phi + \lambda \phi = 0 \quad \text{in } D = \{(r, \theta) : r < R, 0 < \theta < L\},$$
$$\phi|_{\partial D} = 0.$$

**Eigenvalues:**  $\lambda_{m,n} = (j_{\frac{m\pi}{L},n}/R)^2$ , where  $m = 1, 2, \dots$ ,  $n = 1, 2, \dots$ , and  $j_{\frac{m\pi}{L},n}$  is the  $n$ th positive zero of the Bessel function  $J_{\frac{m\pi}{L}}$ .

## Eigenfunctions:

$$\phi_{m,n}(r, \theta) = J_{\frac{m\pi}{L}}(j_{\frac{m\pi}{L},n} \cdot r/R) \sin \frac{m\pi\theta}{L}.$$