

Math 412-501

Theory of Partial Differential Equations

**Lecture 2: Diffusion equation.  
Wave equation. Boundary conditions.**

heat equation:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$$

wave equation:

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

Laplace's equation:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

heat equation:

$$\frac{\partial u}{\partial t} = k \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

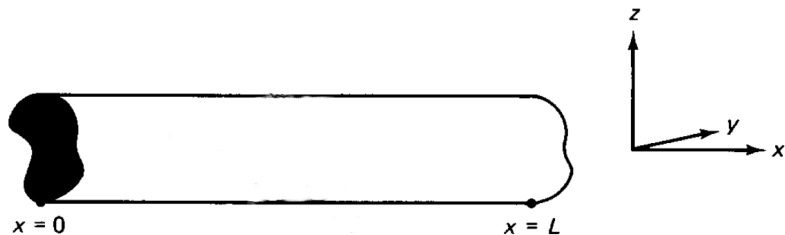
wave equation:

$$\frac{\partial^2 u}{\partial t^2} = c^2 \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

Laplace's equation:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} = 0$$

## Heat conduction in a rod



$u(x, t)$  = temperature

**Heat equation:**

$$c\rho\frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left( K_0 \frac{\partial u}{\partial x} \right) + Q$$

$K_0 = K_0(x)$ ,  $c = c(x)$ ,  $\rho = \rho(x)$ ,  $Q = Q(x, t)$ .

Assuming  $K_0$ ,  $c$ ,  $\rho$  are constant (uniform rod) and  $Q = 0$  (no heat sources), we obtain

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2},$$

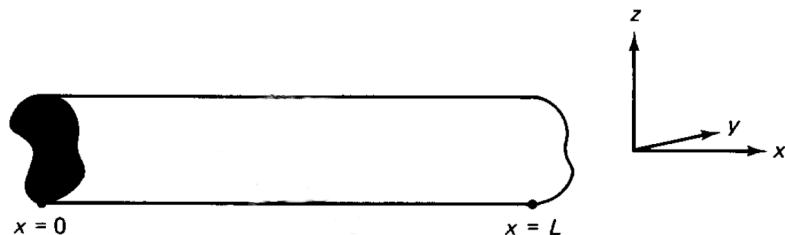
where  $k = K_0(c\rho)^{-1}$  is called the *thermal diffusivity*.

Heat equation is derived from two physical laws:

- conservation of heat energy,
- Fourier's law of heat conduction.

The heat equation is also called the **diffusion equation**.

## Pollutant diffusion in a tube



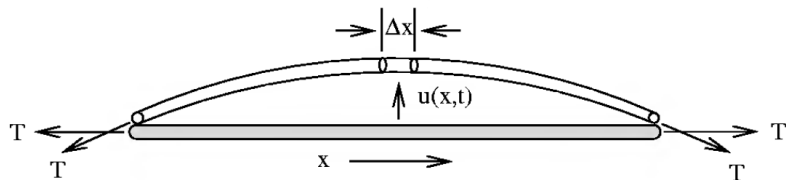
$u(x, t)$  = concentration of the chemical

- conservation of mass
- Fick's law of diffusion

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$$

$k$  = chemical diffusivity

## Vibration of a stretched string



$u(x, t)$  = vertical displacement

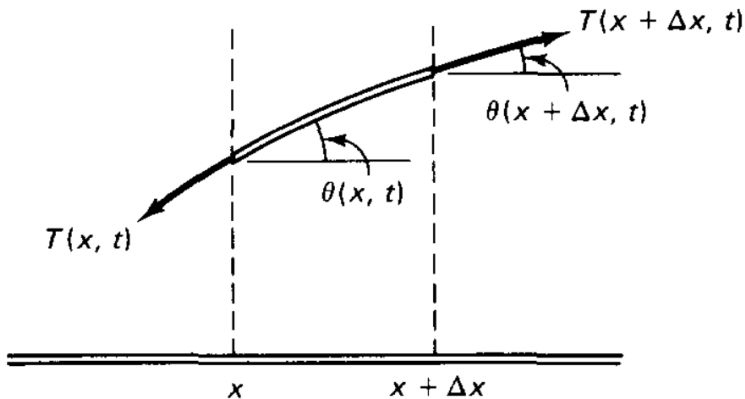
**Newton's law:** mass  $\times$  acceleration = force

$\rho(x)$  = mass density

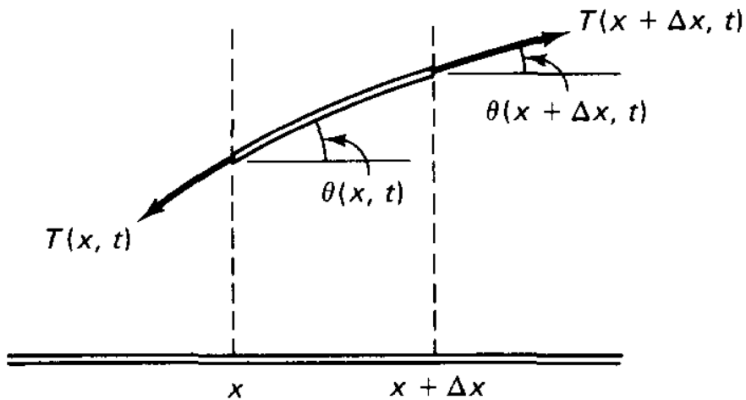
$T(x, t)$  = magnitude of tensile force

$Q(x, t)$  = (vertical) external forces on a unit mass

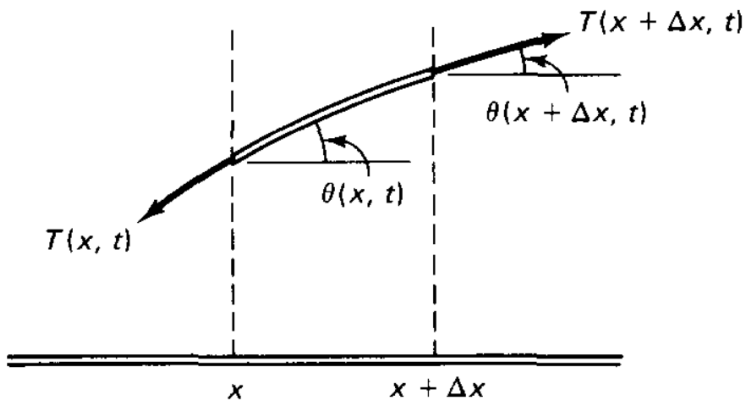




**perfectly flexible string:** no resistance to bending  
 $\theta(x, t) =$  angle between the horizon and the string



$$\tan \theta = \frac{\partial u}{\partial x}$$



vertical component of tensile force =

$$T(x + \Delta x, t) \cdot \sin \theta(x + \Delta x, t) - T(x, t) \cdot \sin \theta(x, t)$$

$$\rho(x) \cdot \Delta x \cdot \frac{\partial^2 u}{\partial t^2} = T(x + \Delta x, t) \cdot \sin \theta(x + \Delta x, t) - T(x, t) \cdot \sin \theta(x, t) + \rho(x) \cdot \Delta x \cdot Q(x, t)$$

$$\rho(x) \cdot \frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial x} \left( T(x, t) \cdot \sin \theta(x, t) \right) + \rho(x) \cdot Q(x, t)$$

We assume that  $\theta \ll 1$ , hence  $\sin \theta \approx \tan \theta$ .

$$\rho(x) \frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial x} \left( T \frac{\partial u}{\partial x} \right) + \rho(x) Q(x, t)$$

**perfectly elastic string:** tension is proportional to stretching (Hooke's law)

Since  $\theta \ll 1$ , we assume  $T(x, t) \approx T_0 = \text{const.}$

$$\rho(x) \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} + \rho(x) Q(x, t)$$

Assuming  $\rho = \text{const}$  and  $Q = 0$ , we obtain

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

where  $c^2 = T_0/\rho$ .

This is **one-dimensional wave equation**.

## Initial and boundary conditions for ODEs

$$y'(t) = y(t), 0 \leq t \leq L.$$

General solution:  $y(t) = C_1 e^t$ , where  $C_1 = \text{const.}$

To determine a unique solution, we need one **initial condition**.

For example,  $y(0) = 1$ . Then  $y(t) = e^t$  is the unique solution.

$$y''(t) = -y(t), \quad 0 \leq t \leq L.$$

General solution:  $y(t) = C_1 \cos t + C_2 \sin t$ , where  $C_1, C_2$  are constant.

To determine a unique solution, we need **two** initial conditions. For example,  $y(0) = 1, y'(0) = 0$ . Then  $y(t) = \cos t$  is the unique solution.

Alternatively, we may impose boundary conditions. For example,  $y(0) = 0, y(L) = 1$ . In the case  $L = \pi/2$ ,  $y(t) = \sin t$  is the unique solution.

**Initial value problem** = ODE + initial conditions

**Boundary value problem** = ODE + boundary conditions

Initial value problem  $y'' = -y$ ,  $y(0) = a$ ,  $y'(0) = b$  always has a unique solution.

Boundary value problem  $y'' = -y$ ,  $y(0) = a$ ,  $y(L) = b$  may not have a unique solution for some triples  $(a, b, L)$ .

For example, let  $L = \pi$  and  $a = 0$ . Then the boundary value problem has no solution if  $b \neq 0$ . In the case  $b = 0$ , it has infinitely many solutions  $y(t) = C_1 \sin t$ ,  $C_1 = \text{const}$ .



## Heat equation

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}, \quad 0 \leq x \leq L, \quad 0 \leq t \leq T.$$

**Initial condition:**  $u(x, 0) = f(x)$ , where  $f : [0, L] \rightarrow \mathbb{R}$ .

**Boundary conditions:**  $u(0, t) = u_1(t)$ ,  
 $u(L, t) = u_2(t)$ , where  $u_1, u_2 : [0, T] \rightarrow \mathbb{R}$ .

Boundary conditions of the **first kind**: prescribed temperature.

Another boundary conditions:  $\frac{\partial u}{\partial x}(0, t) = \phi_1(t)$ ,  
 $\frac{\partial u}{\partial x}(L, t) = \phi_2(t)$ , where  $\phi_1, \phi_2 : [0, T] \rightarrow \mathbb{R}$ .

Boundary conditions of the **second kind**:  
prescribed heat flux.

A particular case:  $\frac{\partial u}{\partial x}(0, t) = \frac{\partial u}{\partial x}(L, t) = 0$   
(insulated boundary).

## Robin conditions:

$$-\frac{\partial u}{\partial x}(0, t) = -h \cdot \left( u(0, t) - u_1(t) \right),$$

$$-\frac{\partial u}{\partial x}(L, t) = h \cdot \left( u(L, t) - u_2(t) \right),$$

where  $h = \text{const} > 0$  and  $u_1, u_2 : [0, T] \rightarrow \mathbb{R}$ .

Boundary conditions of the **third kind**: Newton's law of cooling.

Also, we may consider **mixed** boundary conditions, for example,  $u(0, t) = u_1(t)$ ,  $\frac{\partial u}{\partial x}(L, t) = \phi_2(t)$ .

## Wave equation

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}, \quad 0 \leq x \leq L, \quad 0 \leq t \leq T.$$

**Two** initial conditions:  $u(x, 0) = f(x)$ ,  
 $\frac{\partial u}{\partial t}(x, 0) = g(x)$ , where  $f, g : [0, L] \rightarrow \mathbb{R}$ .

Some boundary conditions:  $u(0, t) = u(L, t) = 0$ .

**Dirichlet conditions:** fixed ends.

Another boundary conditions:

$$\frac{\partial u}{\partial x}(0, t) = \frac{\partial u}{\partial x}(L, t) = 0.$$

**Neumann conditions:** free ends.