

UNDERGRADUATE RESEARCH Scholar

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The Question

- Are numerical solutions to pattern-forming partial differential equations sensitive to time stepping methods?
- The Kurimoto-Sivashinsky Equation is a good model equation to study for this question

Has analogs in atmospheric science



Has analogs in atmospheric science



What makes it so interesting?

- Assume solution is periodic on interval L (common assumption in atmospheric models)
- Given initial condition $u(x, 0) = u_0(x)$
- Divide L in to N parts (so N is the spatial resolution)



- For large values of L (> 12π), equation produces chaotic solutions
- For smaller values of L solutions have a wide array of structure

• Define
$$\hat{L} = \frac{L}{2\pi}$$

Solutions are very sensitive to \hat{L}

Initial condition a randomized wave with small (10^{-5}) amplitude, $\hat{L} = 3.6398, N = 128$



Solutions are very sensitive to \hat{L}

• Increase \hat{L} by 0.0001 with exact same initial condition



Numerical Methods

Separate the spatial (x) and temporal (t) derivatives so it looks like

$$\frac{\partial u}{\partial t} = -u\frac{\partial u}{\partial x} - \frac{\partial^2 u}{\partial x^2} - \frac{\partial^4 u}{\partial x^4}$$

- Use "pseudo-spectral" method for the spatial derivatives
- Time stepping method for time derivative

Solving the spatial derivatives

• If $u(x, t_n)$ is known, then we can use the Discrete Fourier Transform to approximate u(x) as,

$$u(x,t_n) \approx \sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} \tilde{u}(k)e^{i\hat{k}x}$$

Where
$$\hat{k} = \frac{2\pi k}{L}$$

So differentiation becomes simple multiplication

$$\frac{\partial^2 u}{\partial x^2} \approx -\sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} \hat{k}^2 \tilde{u}(k) e^{i\hat{k}x} \qquad \qquad \frac{\partial^4 u}{\partial x^4} \approx \sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} \hat{k}^4 \tilde{u}(k) e^{i\hat{k}x}$$

Solving the spatial derivatives

- For the nonlinear term, since $u(x, t_n)$ is known and $u\frac{\partial u}{\partial x} = \frac{1}{2}\frac{\partial u^2}{\partial x}$,
- Can calculate $\frac{u^2}{2}$ and then proceed as with the linear terms.

• Let
$$v = \frac{u^2}{2}$$

So from the original equation,

$$\frac{\partial u}{\partial t} = -\frac{\partial v}{\partial x} - \frac{\partial^2 u}{\partial x^2} - \frac{\partial^4 u}{\partial x^4}$$

Solving the spatial derivatives

- For the nonlinear terms, since $u(x, t_n)$ is known and $u\frac{\partial u}{\partial x} = \frac{1}{2}\frac{\partial u^2}{\partial x}$,
- Can calculate $\frac{u^2}{2}$ and then proceed as with the linear terms.
- Let $v = \frac{u^2}{2}$
- We have $\frac{\partial \tilde{u}}{\partial t} \approx -\sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} i\hat{k}\tilde{v}(k)e^{i\hat{k}x} + \sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} (\hat{k}^2 - \hat{k}^4)\tilde{u}(k)e^{i\hat{k}x}$

Spectral view of the equation

From the spectral view of the equation,

$$\frac{\partial \widetilde{u}}{\partial t} = -\sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} i\widehat{k}\widetilde{v}(k)e^{i\widehat{k}x} + \sum_{k=-\frac{N}{2}}^{\frac{N}{2}-1} (\widehat{k}^2 - \widehat{k}^4)\widetilde{u}(k)e^{i\widehat{k}x}$$

- The 2nd derivative is a forcing term for low wave numbers $(|\hat{k}| > 1)$
- The 4th derivative is a source of dissipation in the high wave numbers
- The nonlinear term transfers energy from low to high wave numbers



Spectral view of the equation



Time-Stepping methods

- Basic method:
 - Leapfrog
- Two modifications
 - 1. Leapfrog + periodic Predictor-Corrector,
 - 2. Robert-Asselin-Williams (RAW) filter
- How do these two methods compare in the formation of structure in this equation?

Leapfrog time-stepping scheme

The leapfrog scheme (centered difference) is an approximation for the time derivative

$$\frac{u(x,t_n+\Delta t)-u(x,t_n-\Delta t)}{2\Delta t} \approx \frac{\partial u}{\partial t}\Big|_{t_n} = f(u(x,t_n))$$

Where the right hand side is treated as a function *f*This method is unstable



Predictor Corrector method

- Stability can be improved even more by restarting every 25 steps in time.
- When starting from a single initial condition $u(t_0)$, need $u(t_0 + \Delta t)$ to use leapfrog again.
- Use Forward Euler method* to calculate $u\left(t_0 + \frac{\Delta t}{2}\right)$
- Use leapfrog on $u(t_0)$ and $u\left(t_0 + \frac{\Delta t}{2}\right)$ to calculate $u(t_0 + \Delta t)$ then continue as before

Forward Euler:
$$\left. \frac{\partial u}{\partial t} \right|_{t_n} \approx \frac{u(t_n + \Delta t) - u(t_n)}{\Delta t}$$

Robert-Asselin-Williams Filter

A separate improvement on the basic leapfrog scheme is the RAW filter

$$\frac{u(x,t_n+\Delta t)-\overline{u(x,t_n-\Delta t)}}{2\Delta t} = f(\overline{u(x,t_n)})$$

Where,

$$\overline{u(x,t_n)} = u(x,t_n) + \frac{v(1-\alpha)}{2} \left[\overline{u(x,t_n-2\Delta t)} - 2\overline{u(x,t_n-\Delta t)} + u(x,t_n)\right]$$

And,

$$\overline{\overline{u(x,t_n-\Delta t)}} = \overline{u(x,t_n-\Delta t)} - \frac{\nu\alpha}{2} \left[\overline{\overline{u(x,t_n-2\Delta t)}} - 2\overline{u(x,t_n-\Delta t)} + u(x,t_n)\right]$$

Robert-Asselin-Williams Filter

- What this is saying is we compute the next time step
- Then push $u(t_n)$ and $u(t_n + \Delta t)$ towards the



What is the problem?



U EVOLUTION, L*= 5.9, N = 128

What is the problem?



Compared to more accurate method



Compared to more accurate method



Conclusion: Time-stepping method matters

- 4th order Runge-Kutta method takes time and memory
- Can also be difficult to implement in existing code
- In general RAW gives a better idea of the behavior of the solution than the Predictor Corrector method
- It is also very simple to update existing code
- Williams has produced a more general filter to give up to 7th order accuracy
- When looking at the development of structure timestepping matters!