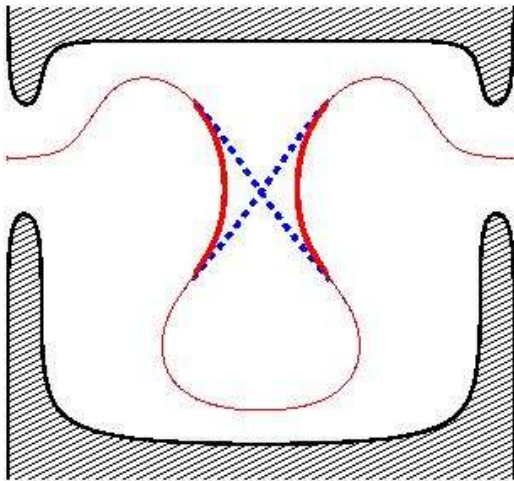


# The classical limit of quantum transport

Saar Rahav

Dept. of Chemistry and Biochemistry  
University of Maryland



With **Piet Brouwer**

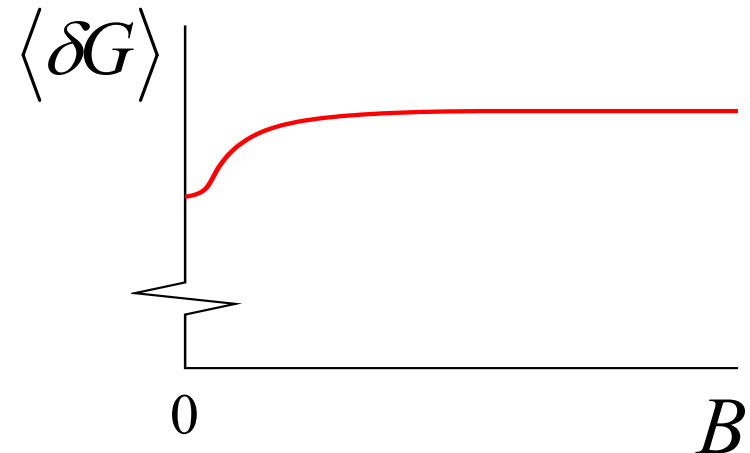
Cornell University

Support: NSF, Packard  
foundation

# Interference effects in transport

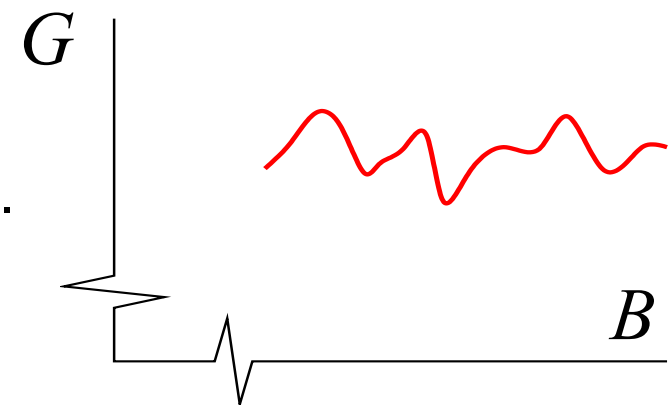
## Weak Localization:

Small negative correction for the average conductance at  $B = 0$



## Conductance fluctuations:

Reproducible fluctuation of the conductance when a parameter is varied.



First studied for disordered samples.

Anderson, Abrahams, Ramakrishnan (1979)

Gorkov, Larkin, Khmel'nitskii (1979)

Altshuler (1985)

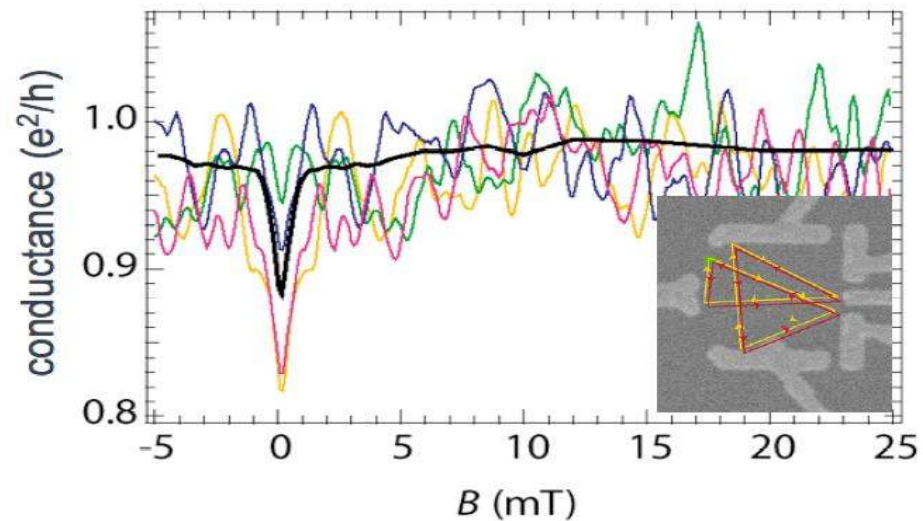
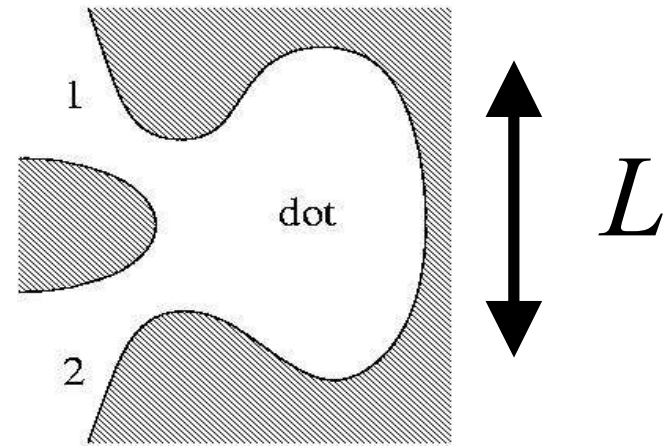
Lee and Stone (1985)

# Ballistic quantum dots

No impurities, smooth boundaries:

Chaotic classical dynamics

Exhibit interference effects:



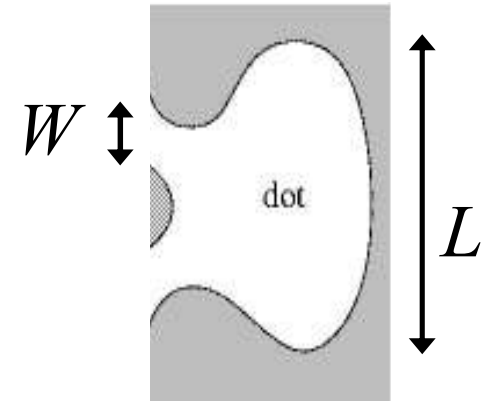
Marcus group

# RMT, Universality

Random matrix theory  $N = k_F W \gg 1$

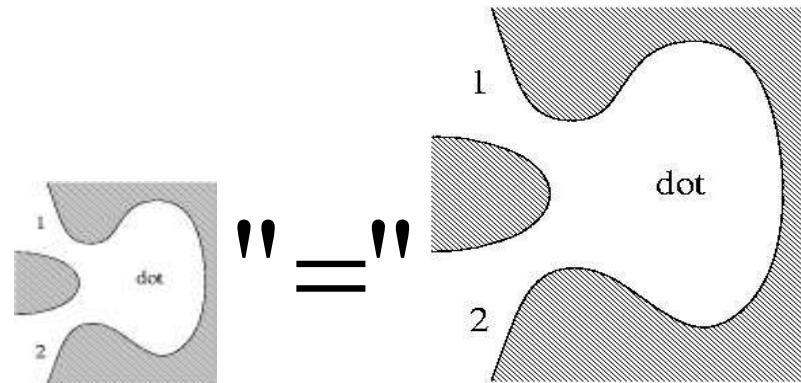
$$\delta G_{\text{RMT}} = \frac{1}{4} \left( \frac{2e^2}{h} \right)$$

$$\text{var } G_{\text{RMT}} = \frac{1}{16} \left( \frac{2e^2}{h} \right)^2$$



Jalabert, Pichard, Beenakker (1994)  
Baranger and Mello (1994)

**RMT:  $\delta G$  and  $\text{var } G$  do not depend on dot size  $L$**



# RMT, Universality ???

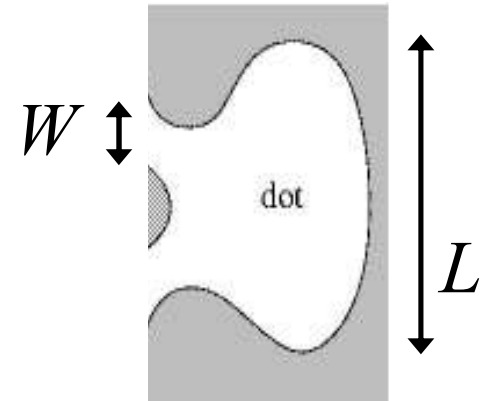
Semiclassical theory

$$\delta G = \delta G_{\text{RMT}} (k_F L)^{-2/\lambda \tau_D}$$

large

small

Aleiner, Larkin  
(1996)



$$\lambda \sim \frac{v_F}{L} \text{ Lyapunov exponent}$$

$$\tau_D \sim \frac{L^2}{W v_F} \text{ dwell time}$$

$$\frac{1}{\lambda \tau_D} \sim \frac{W}{L}$$

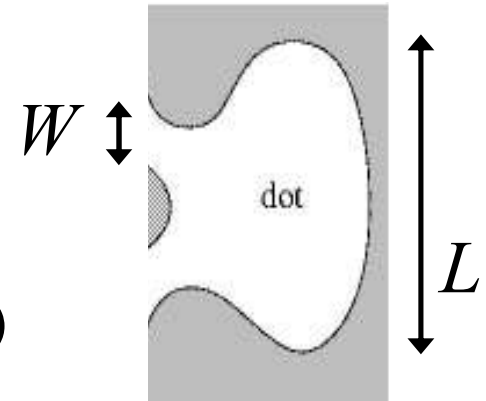
# RMT, Universality ???

## Semiclassical theory

$$\delta G = \delta G_{\text{RMT}} (k_F L)^{-2/\lambda_{\text{TD}}}$$

$$\delta G = \delta G_{\text{RMT}} (k_F L)^{-1/\lambda_{\text{TD}}}$$

Aleiner, Larkin  
(1996)  
Adagideli (2003)



## Numerical simulations

$$\text{var } G = \text{var } G_{\text{RMT}}$$

$$\delta G = \delta G_{\text{RMT}}$$

Tworzydło, Tajic, Beenakker (2004)  
Jacquod and Sukhorukov (2004)

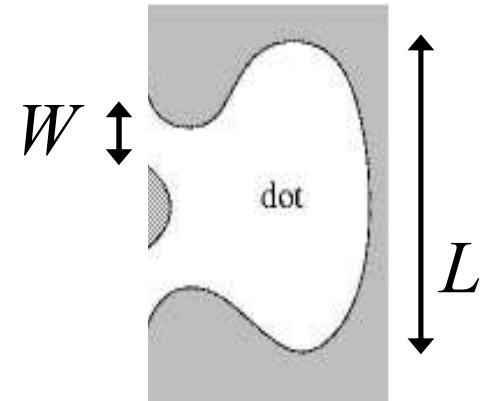
Theory and numerics disagree!

# RMT, Universality ???

## Semiclassical theory

$$\delta G = \delta G_{\text{RMT}} (k_F L)^{-2/\lambda\tau_D}$$

$$\delta G = \delta G_{\text{RMT}} (k_F L)^{-1/\lambda\tau_D}$$



## Numerical simulations

$$\text{var } G = \text{var } G_{\text{RMT}}$$

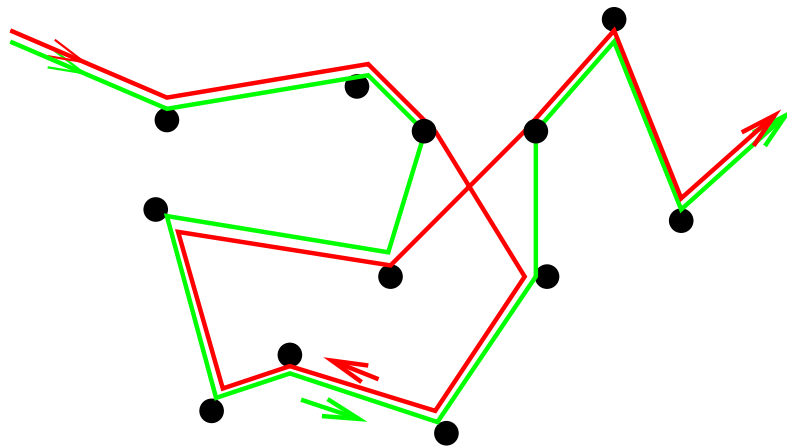
$$\delta G = \delta G_{\text{RMT}}$$

**This talk: Classical correlations  
for**

- Weak localization -  $\delta G$
  - Conductance fluctuation -  $\text{var } G$
- Remark: Parametric correlations**

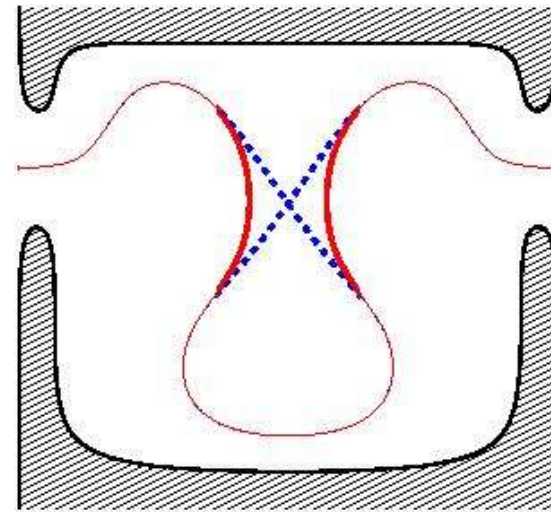
# Semiclassical picture of weak localization

**Disordered systems:**



Anderson, Abrahams, Ramakrishnan (1979), Gorkov, Larkin, Khmelnitskii (1979)

**Ballistic systems:**



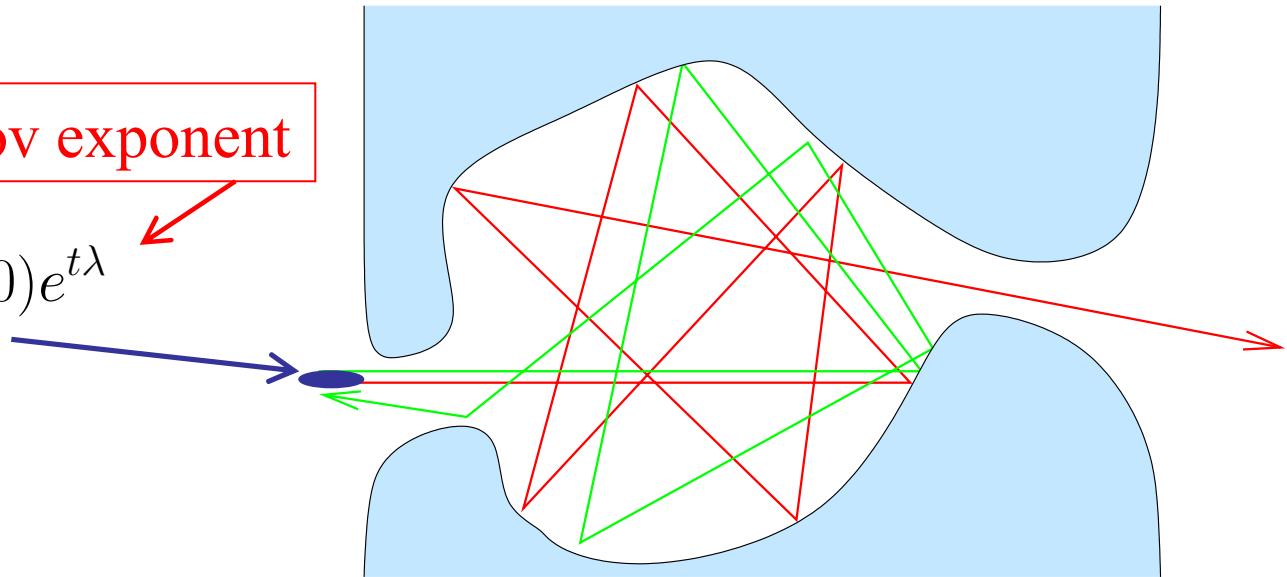
Aleiner and Larkin (1996), Richter and Sieber (2001).

A 'small region' responsible for phase difference



$\lambda$ : Lyapunov exponent

$$\delta \mathbf{r}(t) = \delta \mathbf{r}(0) e^{t\lambda}$$



Quantum uncertainty in position or direction of incoming wavepacket is magnified by chaotic boundary scattering.

Aleiner and Larkin (1996)

# Ehrenfest time

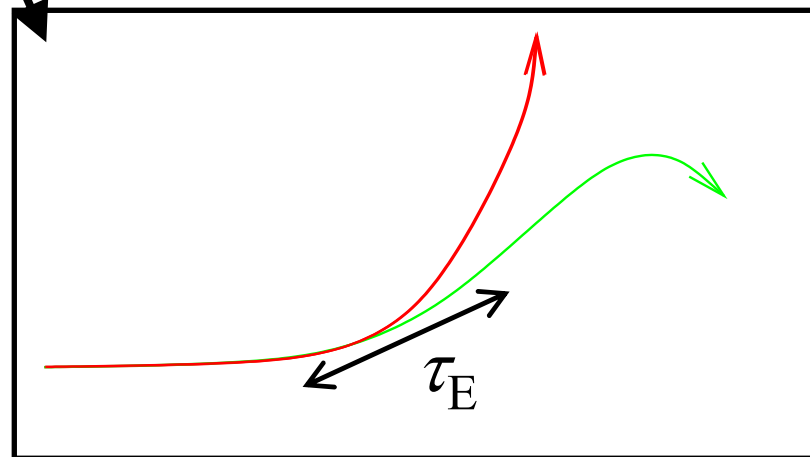
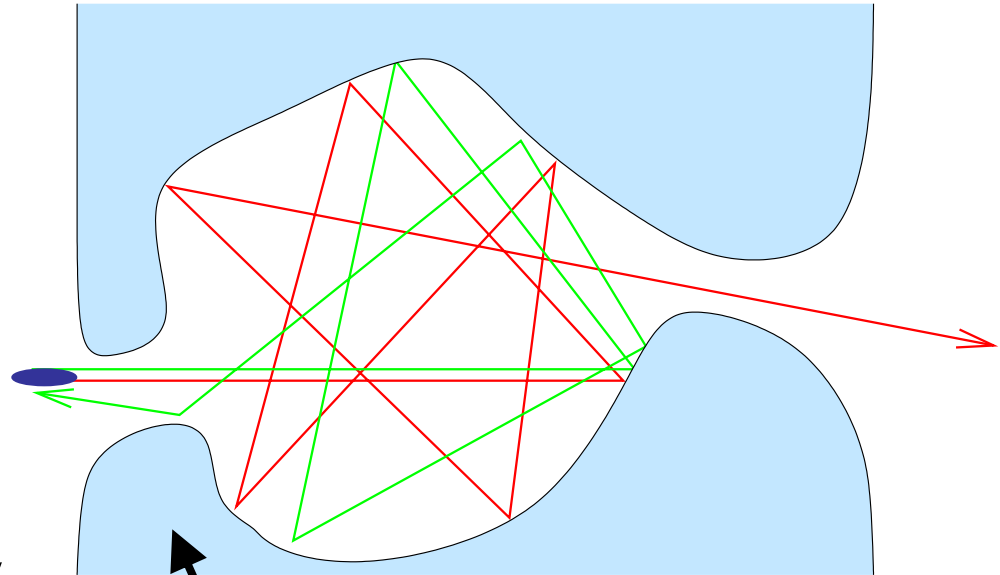
'Ehrenfest time'  $\tau_E$ :  
Time until initial  
uncertainty  $1/k_F$  has  
reached dot size  $L$ :

$$L = \frac{1}{k_F} e^{\lambda \tau_E}$$

$$\tau_E = \frac{1}{\lambda} \ln k_F L$$

$\lambda$ : Lyapunov exponent

Random Matrix Theory  
valid if  $\tau_E \ll \tau_D$



# Theory for $\tau_E \ll \tau_D$

$$G = \frac{2e^2}{h} \sum_{m,n} |t_{mn}|^2 \quad \text{Landauer formula}$$

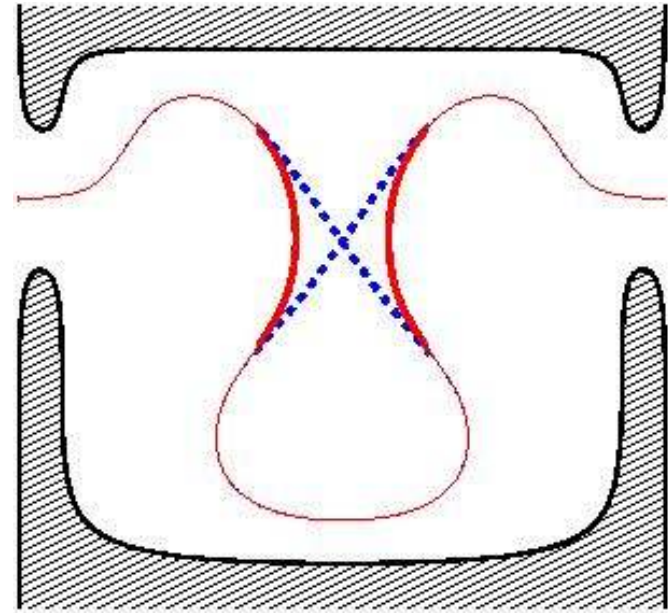
$$|t_{mn}|^2 \sim \sum_{\alpha,\beta} A_\alpha A_\beta e^{i(S_\alpha - S_\beta)/\hbar},$$

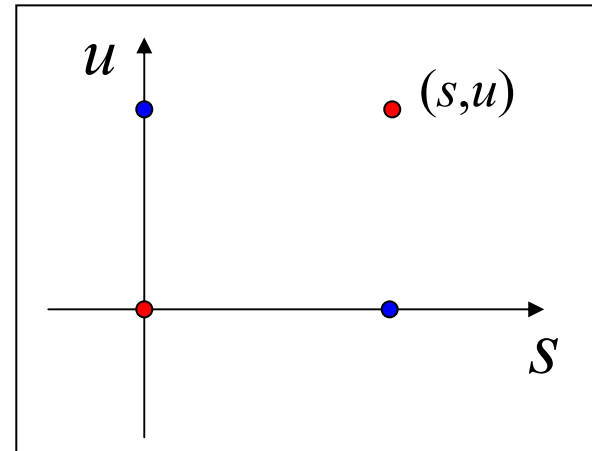
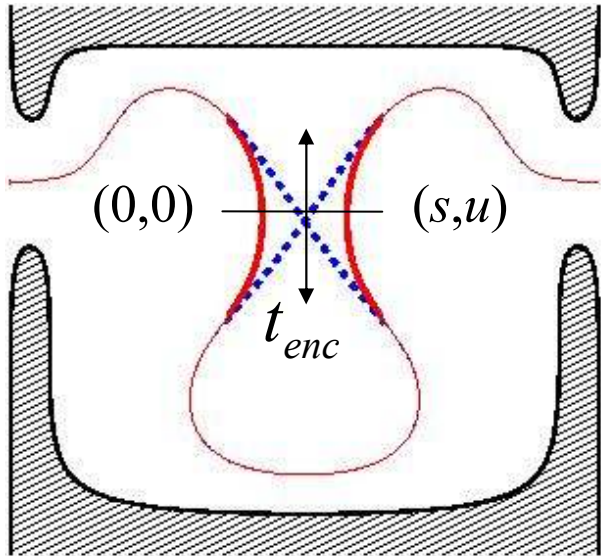
Jalabert, Baranger, Stone (1990)

- $S_\alpha, S_\beta$ : classical action
- angles of  $\alpha, \beta$  consistent with transverse momentum in lead,

$$p_\perp(m) = \pm \pi \hbar m / W_j, \quad m = 1, \dots, N_j,$$

- $A_\alpha, A_\beta$ : stability amplitudes





$s, u$ : distances along stable, unstable phase space directions

encounter region:  $|s|, |u| < c$

$c$ : classical cut-off scale

$$t_{\text{enc}} = \frac{1}{\lambda} \left( \ln \frac{c}{|s|} + \ln \frac{c}{|u|} \right)$$

Action difference  $S_\alpha - S_\beta = su$

$$\begin{array}{l} s \propto e^{-\lambda t} \\ u \propto e^{\lambda t} \end{array}$$

Richter and Sieber (2003)

Spehner (2003)

Turek and Richter (2003)

Müller *et al.* (2004)

Heusler *et al.* (2006)

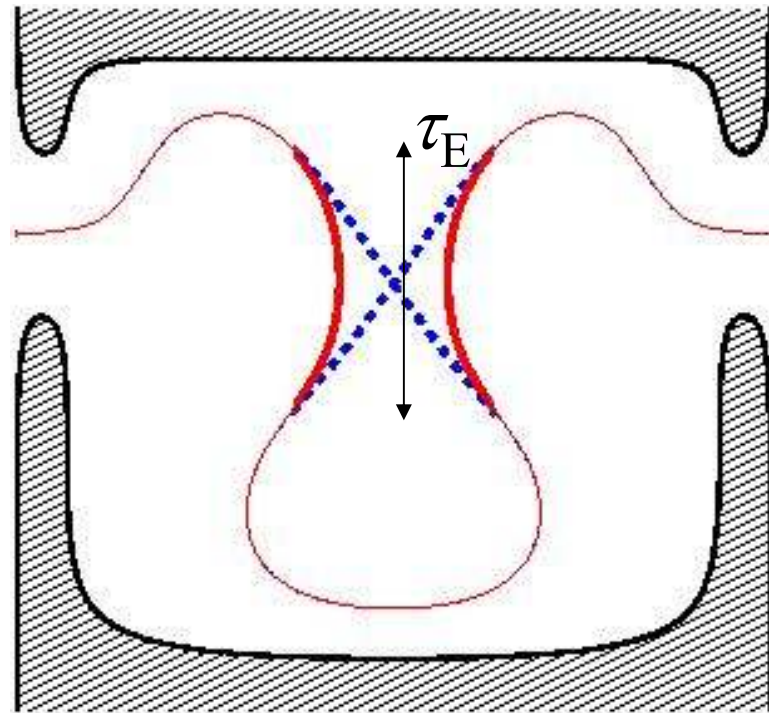
# Theory for $\tau_E \sim \tau_D$

Encounters have 'size':

$$S_\alpha - S_\beta \simeq \hbar$$



$$t_{\text{enc}} \simeq \frac{1}{\lambda} \ln \frac{c^2}{\hbar} = \tau_E$$

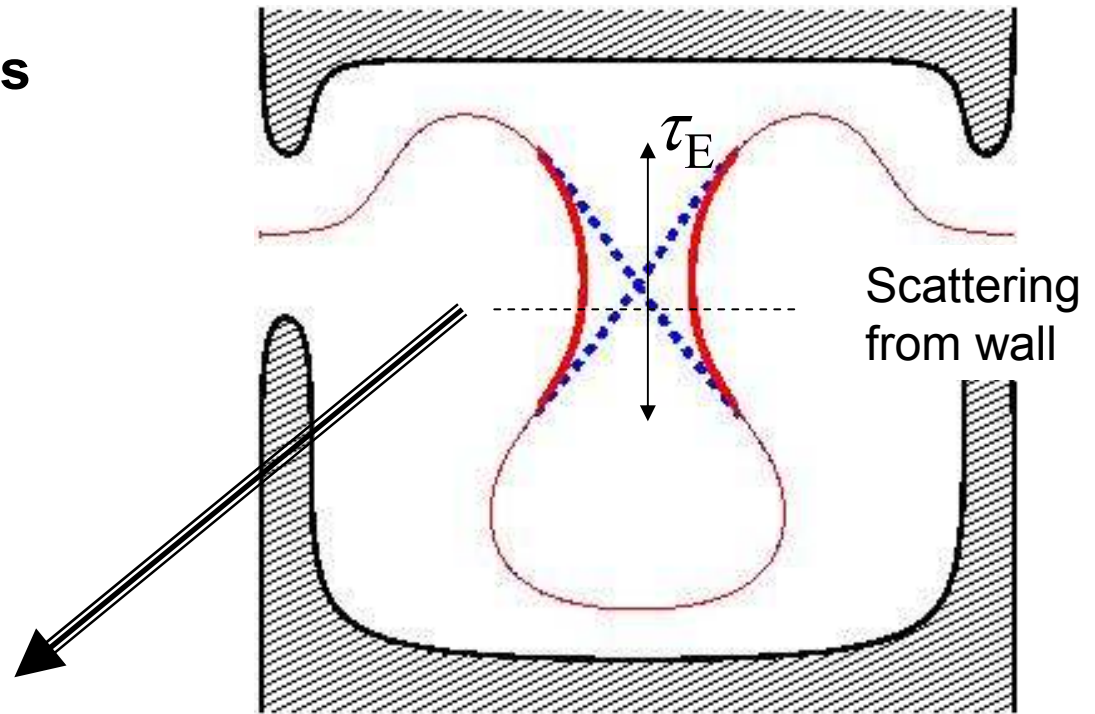
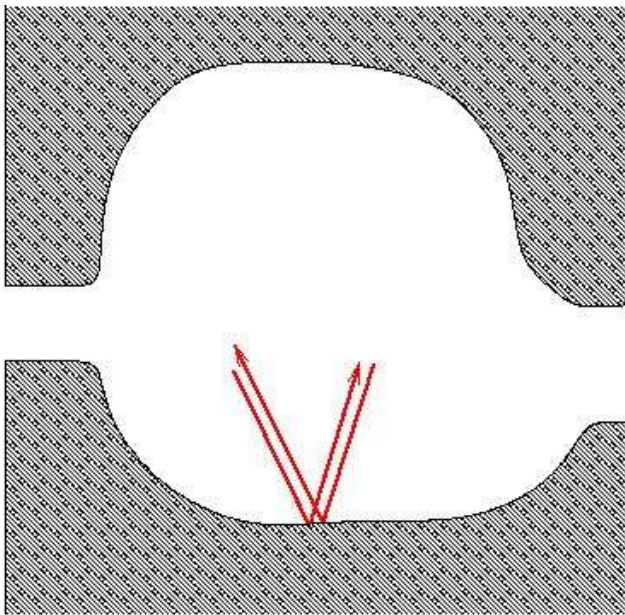


Uncorrelated escape:  $P(t) = e^{-t/\tau_D}$

$$P_{\text{stay}} = ??$$

# Theory for $\tau_E \sim \tau_D$

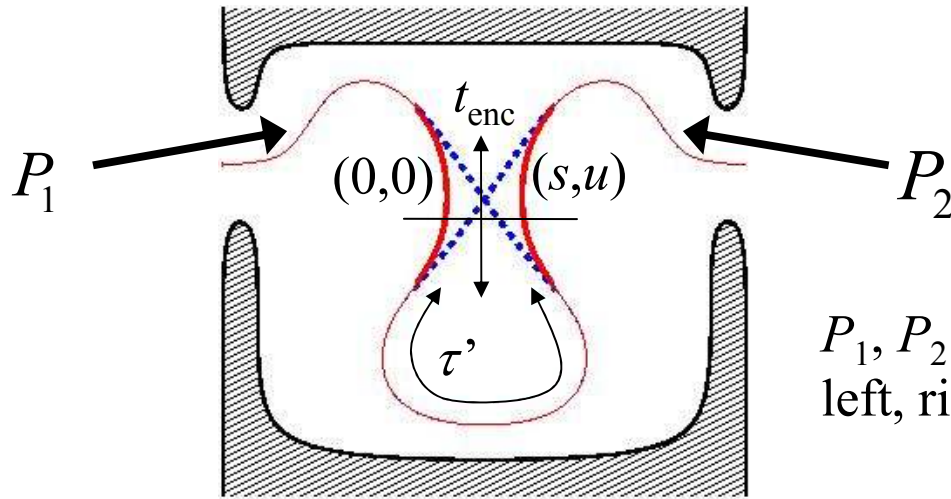
Correlated motion means correlated escape !



$$P_{stay} = e^{-t_{enc}/\tau_D} \simeq e^{-\tau_E/\tau_D}$$

Not  $e^{-2\tau_E/\tau_D}$

# Theory for $\tau_E \sim \tau_D$

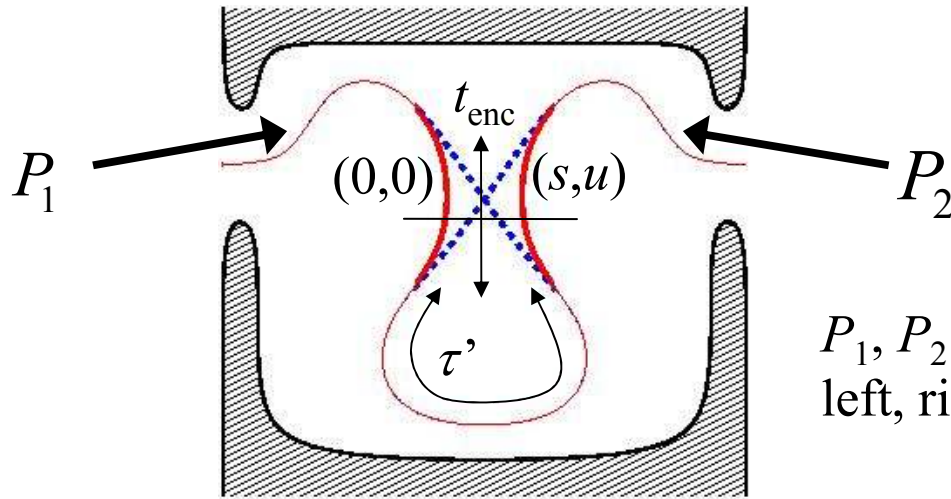


$P_1, P_2$ : probabilities to enter, exit through left, right contacts

$$\begin{aligned} \delta G &= -\frac{2e^2}{h} P_1 P_2 \int d\tau' \int_{-c}^c ds du \frac{e^{isu/\hbar - (t_{\text{enc}} + \tau')/\tau_D}}{2\pi\hbar t_{\text{enc}}} \\ &= \frac{2e^2}{h} P_1 P_2 e^{-\tau_E/\tau_D} \end{aligned}$$

SR and Brouwer (2005), Brouwer and SR (2006), Jacquod and Whitney (2006).

# Theory for $\tau_E \sim \tau_D$



$P_1, P_2$ : probabilities to enter, exit through left, right contacts

$$\begin{aligned}\delta G &= \delta G_{\text{RMT}} e^{-\tau_E/\tau_D} \\ &= \delta G_{\text{RMT}} (k_F L)^{-1/\lambda\tau_D}\end{aligned}$$

SR and Brouwer (2005), Brouwer and SR (2006), Jacquod and Whitney (2006).

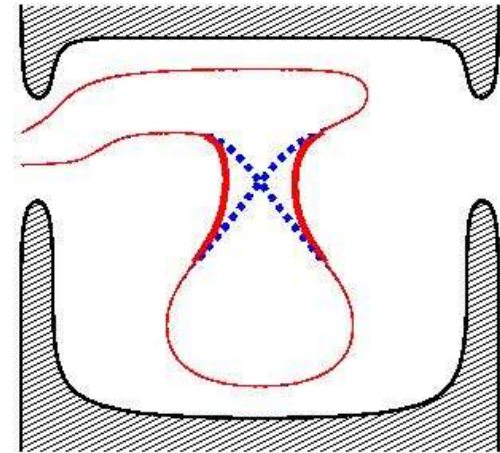


# Unitarity

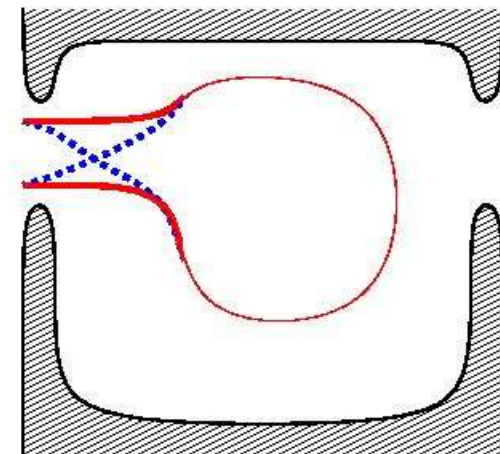
$\delta G$  can also be calculated through the dot's reflection

$$\begin{aligned}\delta G &= -2\frac{e^2}{h}P_1^2e^{-\tau_E/\tau_D} + 2\frac{e^2}{h}P_1e^{-\tau_E/\tau_D} \\ &= 2\frac{e^2}{h}P_1(1 - P_1)e^{-\tau_E/\tau_D} \\ &= \delta G_{\text{RMT}}e^{-\tau_E/\tau_D}\end{aligned}$$

$P_1, P_2=1-P_1$ : probabilities to enter, exit through left, right contacts



Weak localization

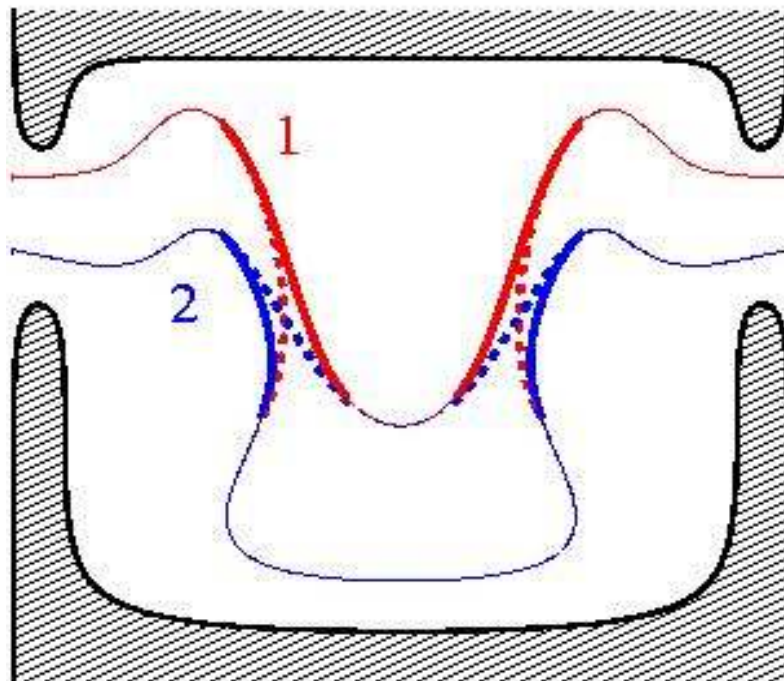


Coherent backscattering

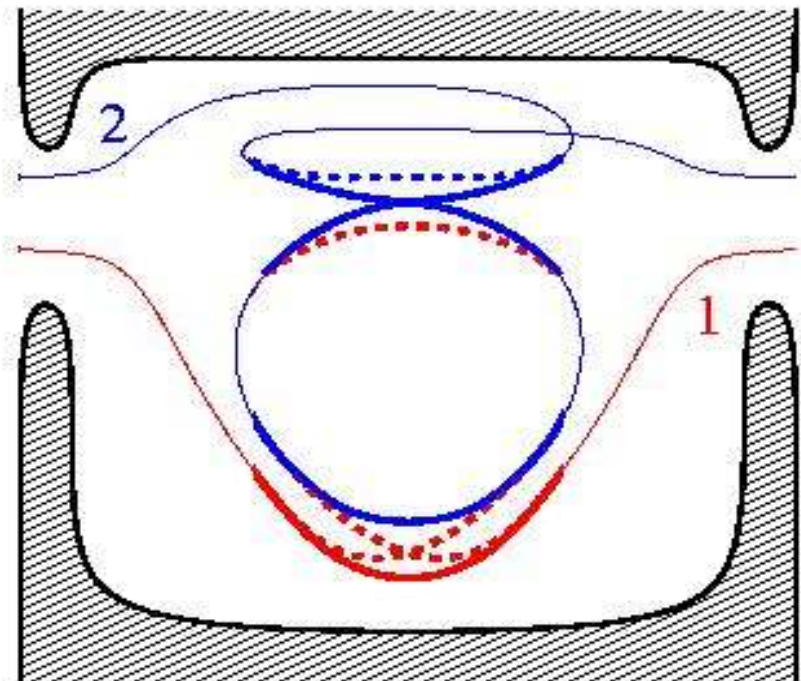
# Conductance fluctuations

Two classes of trajectories:

“Diffusion constant”



“Density of states”



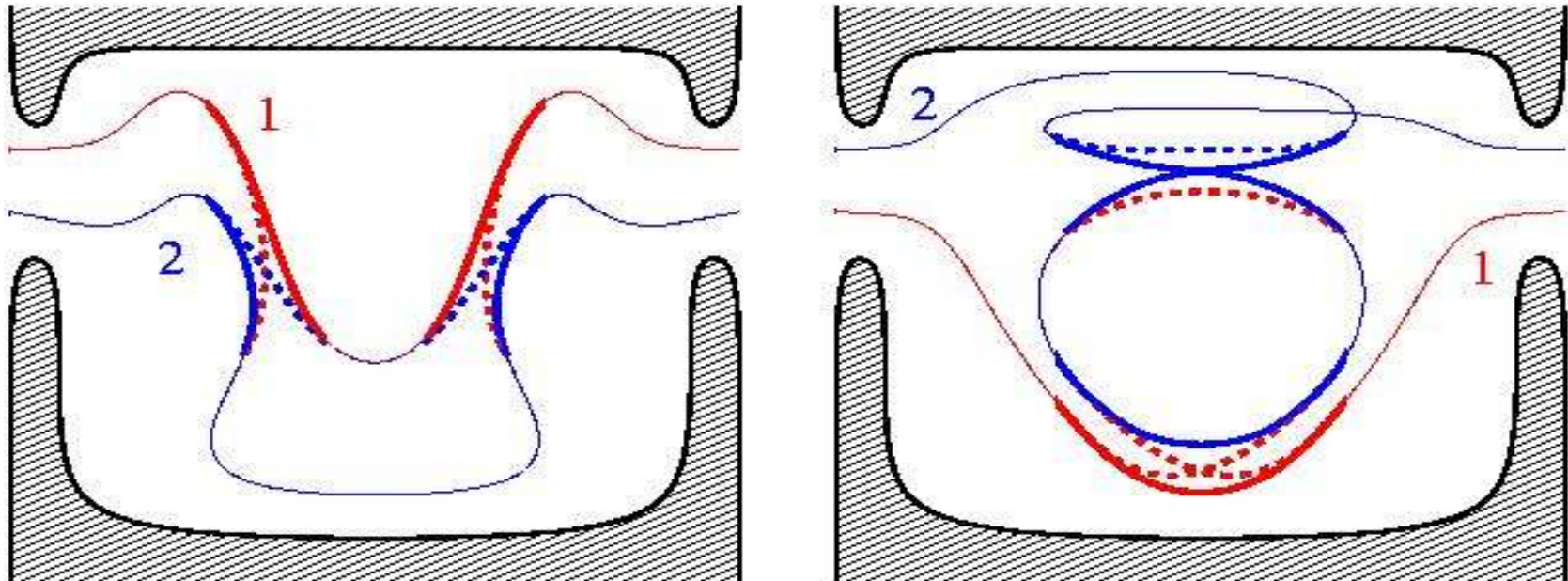
RMT:

1

:

0

# Conductance fluctuations



Semiclassical theory:

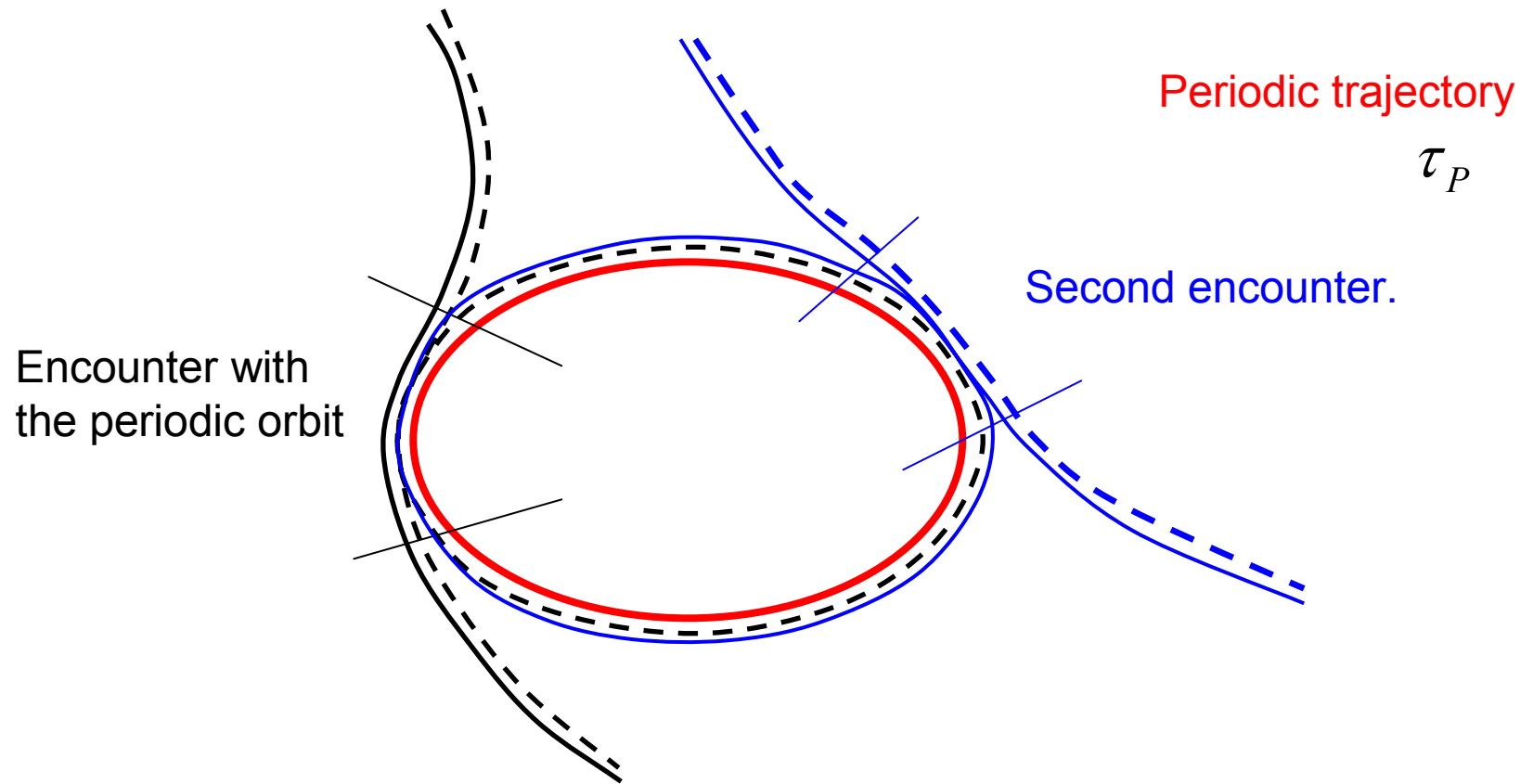
$$\text{var } G_{\text{RMT}} \times e^{-2\tau_E/\tau_D} \quad + \quad \text{var } G_{\text{RMT}} \times (1 - e^{-2\tau_E/\tau_D})$$

$$\text{var } G = \text{var } G_{\text{RMT}}$$

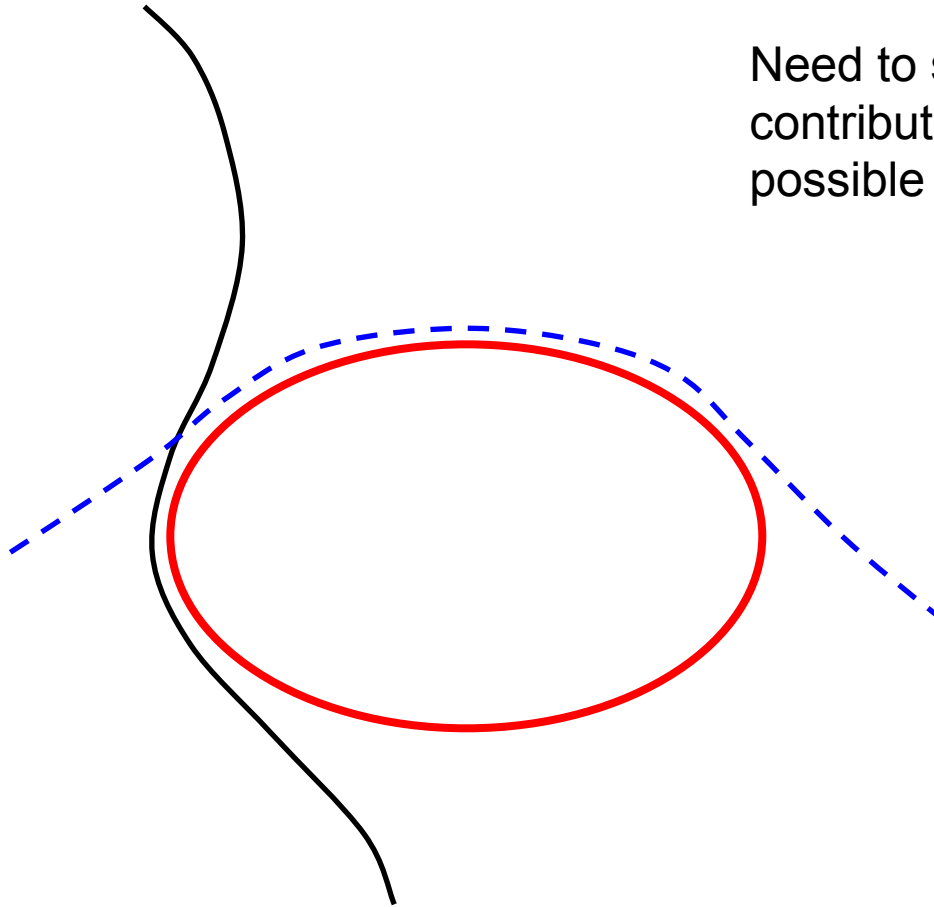
Brouwer and SR (2006). Brouwer (2007).

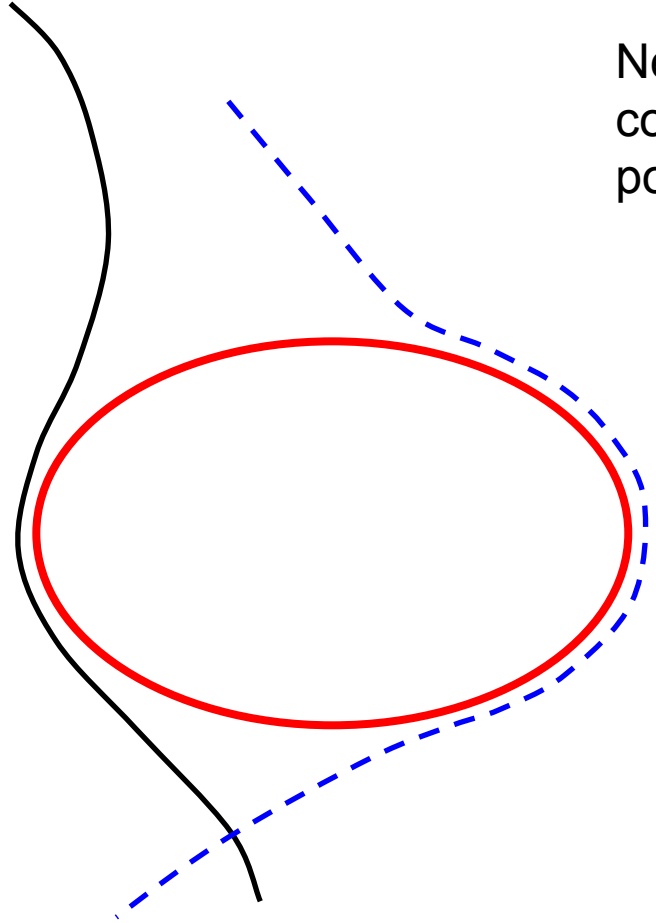
**Conductance fluctuations independent of Ehrenfest time !**

# The 'density-of-states' configurations:

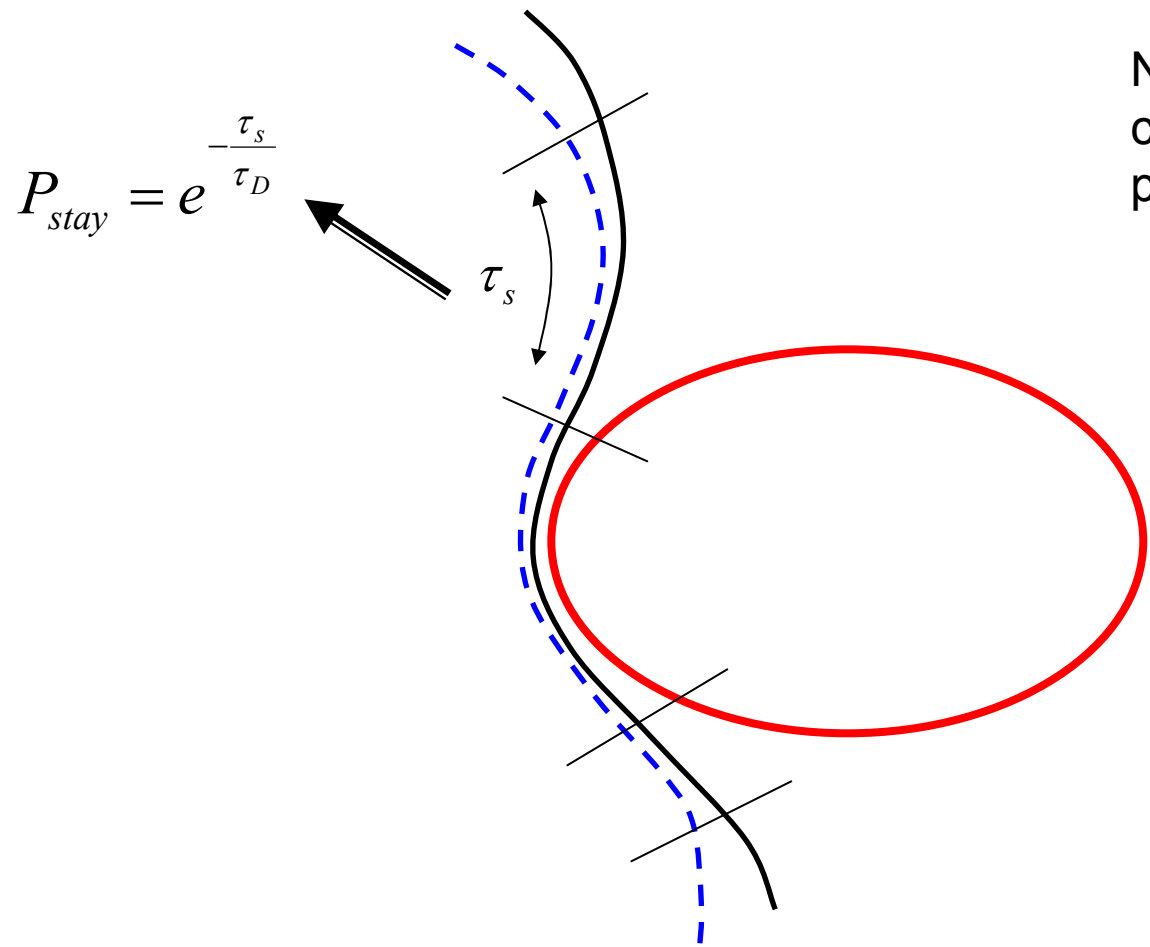


Need to sum  
contributions from all  
possible encounters





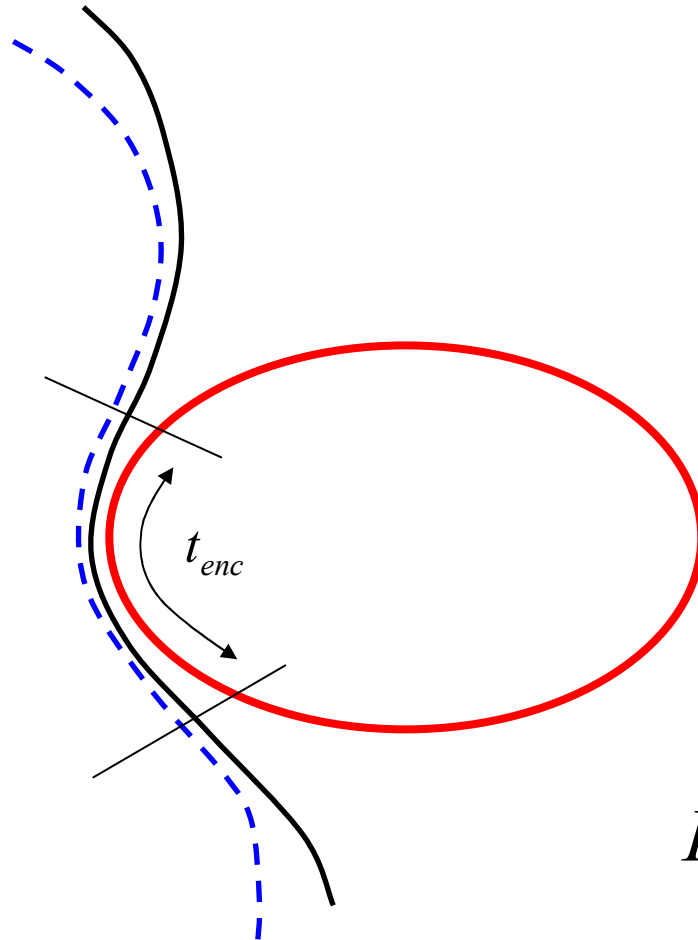
Need to sum  
contributions from all  
possible encounters



Need to sum contributions from all possible encounters

Systematic, non-oscillatory contribution from trajectories correlated away from the periodic orbit

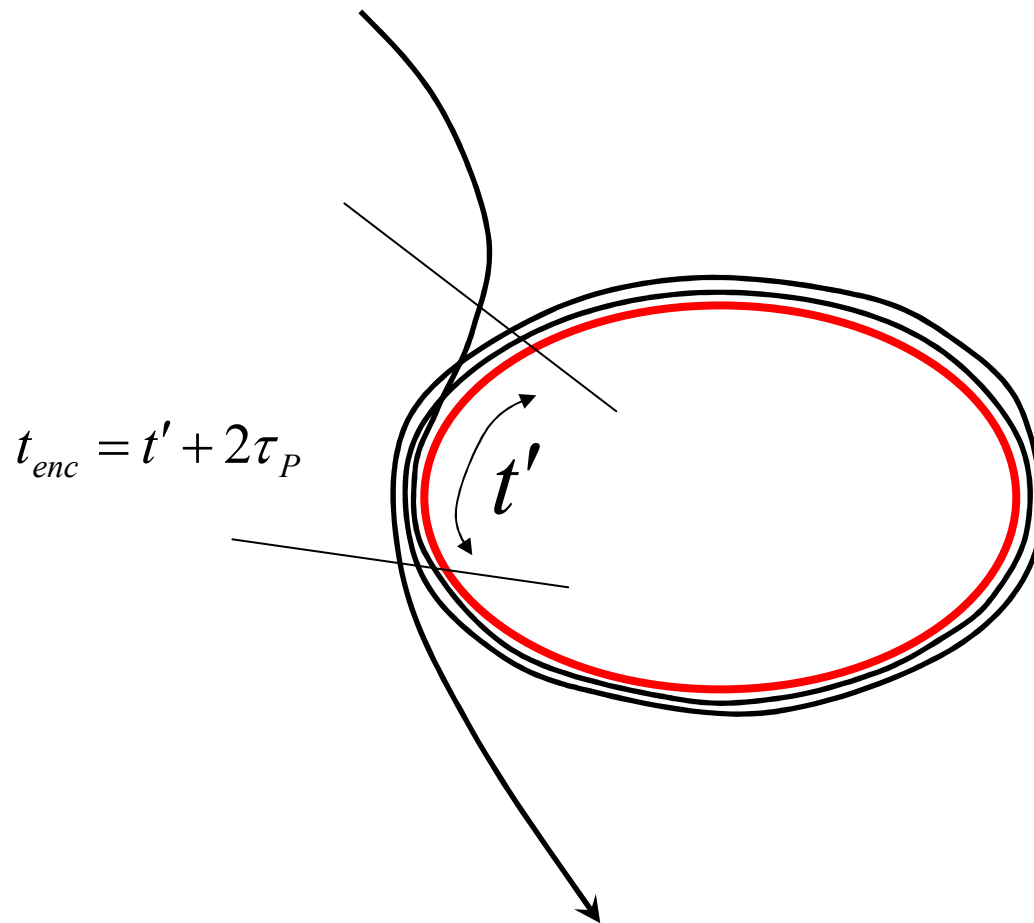
Survival probability  $t_{enc} < \tau_P$



$$P = \exp\left[-\frac{t_{enc}}{\tau_D}\right]$$



Survival probability:  $t_{enc} > \tau_P$



$$P = \exp\left[-\frac{\tau_P}{\tau_D}\right]$$

Can encircle many times without affecting the survival probability!

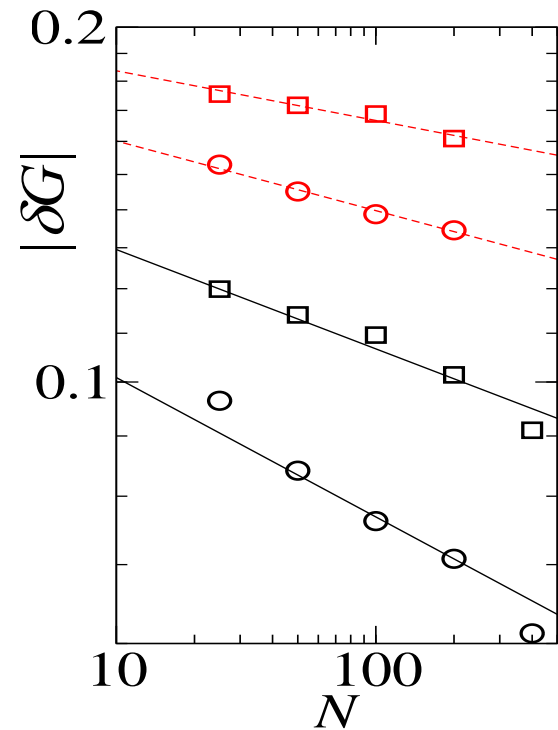
# Numerical results

$$\begin{aligned}\delta G &= \delta G_{\text{RMT}} e^{-\tau_E/\tau_D} \\ &= \delta G_{\text{RMT}} (k_F L)^{-1/\lambda\tau_D}\end{aligned}$$

$$\ln|\delta G| = -\frac{1}{\lambda\tau_D} \ln N + \text{const}$$

$\lambda, \tau_D$  are known

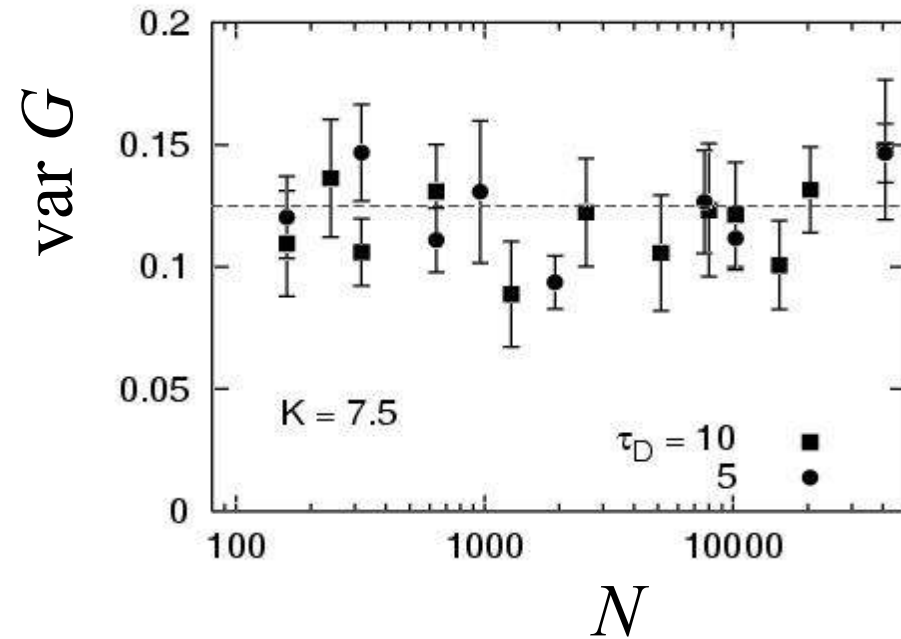
Numerical results consistent with theory !



SR and Brouwer (2005)

# Numerical results

$$\text{var } G = \text{var } G_{\text{RMT}}$$



Tworzydło *et al.* (2004)

Jacquod and Sukhorukov (2004)

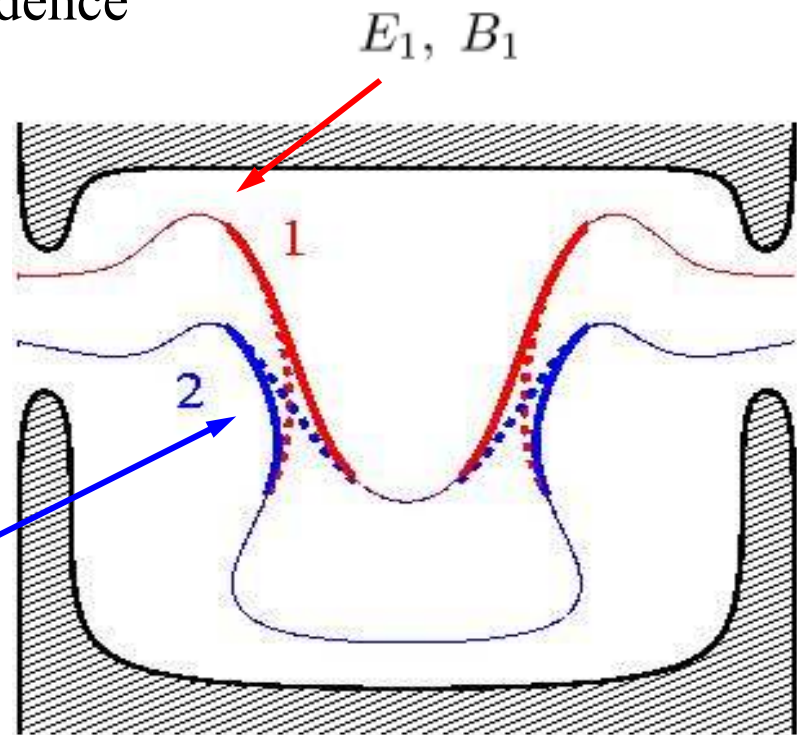
# Parametric correlations

Is there a deep reason for  $\tau_E$  independence of  $\text{Var}(G)$ ?

Study:

$$\langle \delta G(E_1, B_1) \delta G(E_2, B_2) \rangle$$

$E_2, B_2$



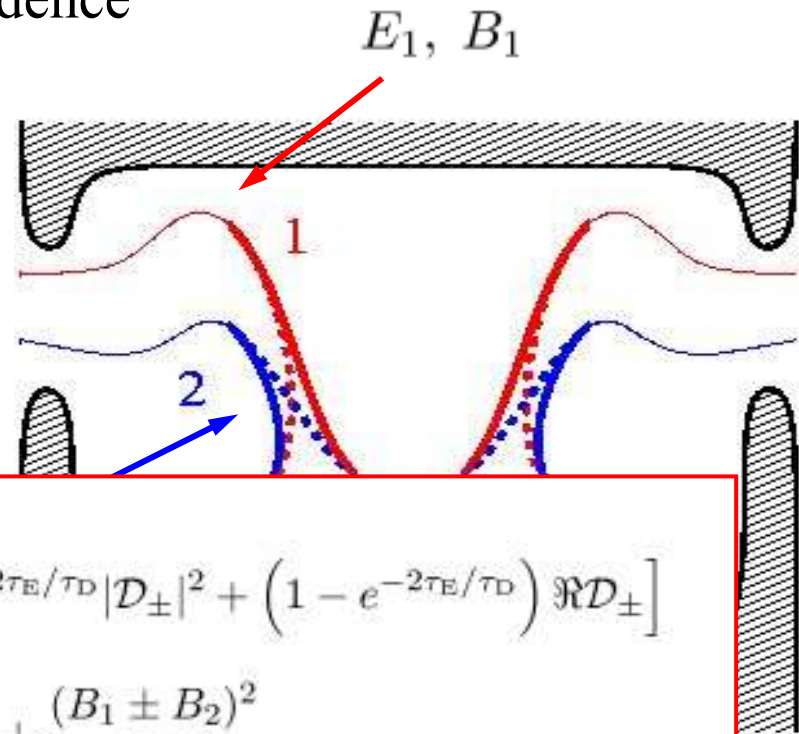
$$\begin{aligned} \langle \delta S \rangle &\propto i(E_1 - E_2)t \\ \langle \delta S^2 \rangle &\propto (B_1 - B_2)^2 t \end{aligned}$$

# Parametric correlations

Is there a deep reason for  $\tau_E$  independence of  $\text{Var}(G)$ ?

Study:

$$\langle \delta G(E_1, B_1) \delta G(E_2, B_2) \rangle$$



$$\langle \delta G(E_1, B_1) \delta G(E_2, B_2) \rangle = P_1^2 P_2^2 \sum_{\pm} \left[ e^{-2\tau_E/\tau_D} |\mathcal{D}_{\pm}|^2 + (1 - e^{-2\tau_E/\tau_D}) \Re \mathcal{D}_{\pm} \right]$$

$$\mathcal{D}_{\pm}^{-1} = 1 - i \frac{E_1 - E_2}{E^*} + \frac{(B_1 \pm B_2)^2}{2(B^*)^2}$$

The parametric correlations depend on  $\tau_E$ . New, non-RMT, universal regime. Brouwer and SR (2007)

# Summary:

**Wave phenomena in ballistic quantum dots appear only after the Ehrenfest time!**

## Summary of results:

- **Weak localization**  $\delta G = \delta G_{\text{RMT}} \times e^{-\tau_E/\tau_D}$
- **Conductance fluctuations**  $\text{var } G = \text{var } G_{\text{RMT}}$
- Many other: shot noise, dephasing, coherent backscattering, quantum pumping ...