

Solutions to Chapter 3

1. Summing the two trig identities that are given in the problem yields

$$\cos mt \cos \lambda t = (\cos(\lambda + m)t + \cos((\lambda - m)t) / 2$$

Integrating the expression on the right gives

$$\int_{-\pi}^{\pi} \cos mt \cos \lambda t dt = \frac{1}{2} \left(\frac{\sin(\lambda + m)t}{\lambda + m} + \frac{\sin(\lambda - m)t}{\lambda - m} \right) \Big|_{-\pi}^{\pi}$$

Inserting the limits and simplifying yields

$$\frac{-2(-1)^m \lambda \sin(\lambda \pi)}{(m^2 - \lambda^2)}$$

Setting $m = 3$ shows that the Fourier transform of $\cos(3t)$ (from $-\pi$ to π) is

$$-\frac{\sqrt{2} \lambda \sin(\pi \lambda)}{\sqrt{\pi} (-9 + \lambda^2)}$$

as claimed.

2. Since \sin is an odd function, we only need to compute the sine part of the Fourier transform. That is, we must compute

$$\frac{-i}{\sqrt{2\pi}} \int_{-\pi}^{\pi} \sin 3t \sin \lambda t dt$$

Subtracting the two trig identities given in problem 1, gives

$$\frac{-i}{\sqrt{2\pi}} \sin 3t \sin \lambda t = \frac{-i}{\sqrt{2\pi}} (\cos(\lambda - 3)t - \cos(\lambda + 3)t) / 2$$

which when integrated from $-\pi$ to π gives

$$\frac{i}{2\sqrt{2\pi}} \left(\frac{\sin(\lambda - 3)t}{\lambda - 3} - \frac{\sin(\lambda + 3)t}{\lambda + 3} \right) \Big|_{-\pi}^{\pi} = 3 \frac{i \sqrt{2} \sin(\pi \lambda)}{\sqrt{\pi} (-3 + \lambda) (3 + \lambda)}$$

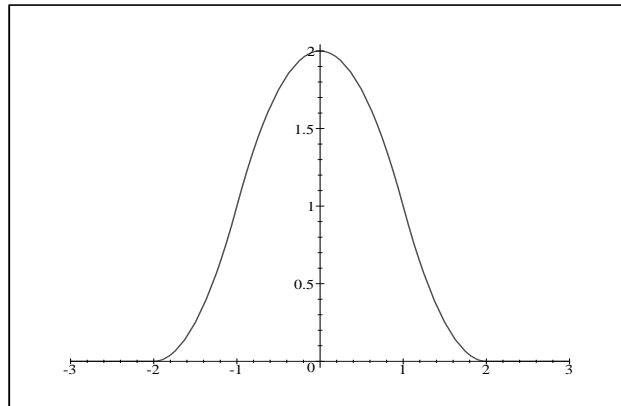
3. The Maple code for displaying f and computing its Fourier transform is given below.

```
> f:=t-> piecewise(t<-2,0,t<-1,t^2+4*t+4,t<1,2-t^2,t<2,t^2-4*t+4,0);
f := t -> piecewise(t < -2, 0, t < -1, t^2+4t+4, t < 1, 2-t^2, t < 2, t^2-4t+4, 0)
```

```
> f(t);
```

$$\begin{cases} 0 & t < -2 \\ t^2 + 4t + 4 & -2 < t < -1 \\ 2 - t^2 & -1 < t < 1 \\ t^2 - 4t + 4 & 1 < t < 2 \\ 0 & \text{otherwise} \end{cases}$$

```
> plot(f,-3..3);
```



Note that f is continuous. Each piece matches up at the boundary points. For example the limit of $t^2 + 4t + 4$ as t approaches -2 is zero, which is the same as the value of f on the interval $t < -2$. The other pieces match up similarly.

The derivative is given by

```
> D(f)(t);
```

$$\begin{cases} 0 & t < -2 \\ 2t + 4 & -2 < t < -1 \\ -2t & -1 < t < 1 \\ 2t - 4 & 1 < t < 2 \\ 0 & \text{otherwise} \end{cases}$$

The derivative is also continuous since each of its pieces match up at the boundary defining points. For example, the limit of $2t + 4$ as t approaches -2 is 0 which agrees with f' for $t < -2$.

Here is the computation of the Fourier transform. Since f is even, we need only compute the cosine part of the transform.

```
> (1/sqrt(2*Pi))*Int(f(t)*cos(lambda*t),t=-2..2); value(");
```

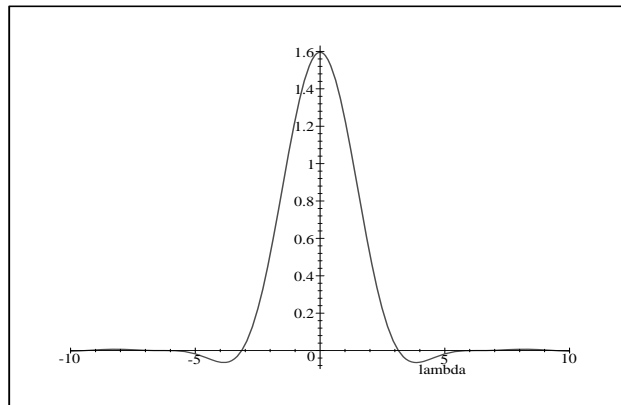
$$\frac{1}{2} \frac{\sqrt{2} \int_{-2}^2 \begin{cases} 0 & t < -2 \\ t^2 + 4t + 4 & -2 < t < -1 \\ 2 - t^2 & -1 < t < 1 \\ t^2 - 4t + 4 & 1 < t < 2 \\ 0 & \text{otherwise} \end{cases} \cos(\lambda t) dt}{\sqrt{\pi}} = -4 \frac{\sqrt{2} \sin(\lambda) (\cos(\lambda) - 1)}{\sqrt{\pi} \lambda^3}$$

> trans:=";

$$trans := -4 \frac{\sqrt{2} \sin(\lambda) (\cos(\lambda) - 1)}{\sqrt{\pi} \lambda^3}$$

Here is the plot of the transform over the interval $-10 \leq \lambda \leq 10$.

> plot(trans, lambda=-10..10);



Note that the Fourier transform decays like $1/\lambda^3$. There is a pattern here. A function which has finite support and bounded leads to a Fourier transform that decays like $1/\lambda$ (see example 21 in the text, which was also done in class). If the function is also continuous then its transform decays like $1/\lambda^2$ (such as in problem 1). If the function also has a continuous derivative, then the transform decays like $1/\lambda^3$ as in this exercise.

5. To prove the \mathcal{F}^{-1} is linear, we have

$$\begin{aligned} \mathcal{F}^{-1}[f + g](x) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [f(\lambda) + g(\lambda)] e^{i\lambda x} d\lambda \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [f(\lambda)] e^{i\lambda x} d\lambda + \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [g(\lambda)] e^{i\lambda x} d\lambda \\ &= \mathcal{F}^{-1}[f](x) + \mathcal{F}^{-1}[g](x) \end{aligned}$$

To prove

$$\mathcal{F}^{-1}[\lambda^n f(\lambda)](t) = (-i)^n \frac{d^n}{dt^n} \{ \mathcal{F}^{-1}[f](t) \}$$

We start with the right side:

$$\begin{aligned} (-i)^n \frac{d^n}{dt^n} \{ \mathcal{F}^{-1}[f](t) \} &= (-i)^n \frac{d^n}{dt^n} \left\{ \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [f(\lambda)] e^{i\lambda t} d\lambda \right\} \\ &= (-i)^n \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [f(\lambda)] \frac{d^n}{dt^n} \{ e^{i\lambda t} \} d\lambda \\ &= (-i)^n \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [f(\lambda)] (i\lambda)^n e^{i\lambda t} d\lambda \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [f(\lambda)] (\lambda)^n e^{i\lambda t} d\lambda \\ &= \mathcal{F}^{-1}[\lambda^n f(\lambda)](t) \end{aligned}$$

To prove

$$\mathcal{F}^{-1}[f^{(n)}(\lambda)](t) = (-it)^n \mathcal{F}^{-1}[f](t)$$

we start with the left side with $n = 1$:

$$\mathcal{F}^{-1}[f'(\lambda)](t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f'(\lambda) e^{i\lambda t} d\lambda$$

We integrate by parts with $dv = f'$ and $u = e^{i\lambda t}$ to obtain

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f'(\lambda) e^{i\lambda t} d\lambda = \frac{1}{\sqrt{2\pi}} \left(f(\lambda) e^{i\lambda t} \Big|_{\lambda=-\infty}^{\infty} - \int_{-\infty}^{\infty} f(\lambda) (ite^{i\lambda t}) d\lambda \right)$$

The boundary terms are zero since $f(\lambda) = 0$ for large λ . Therefore

$$\mathcal{F}^{-1}[f'(\lambda)](t) = - \int_{-\infty}^{\infty} f(\lambda) (ite^{i\lambda t}) d\lambda = -it \mathcal{F}^{-1}[f(\lambda)](t)$$

To show the statement regarding the Laplace transform, we have

$$\mathcal{F}[f](\lambda) = \frac{1}{\sqrt{2\pi}} \int_0^{\infty} f(t) e^{-i\lambda t} dt$$

Since $f(t) = 0$ for $t < 0$, the lower limit of integration starts at 0 rather than $-\infty$. On the other hand

$$\frac{1}{\sqrt{2\pi}} \mathcal{L}f(i\lambda) = \frac{1}{\sqrt{2\pi}} \int_0^{\infty} f(t) e^{-ti\lambda} dt$$

which is the same as $\mathcal{F}[f](\lambda)$.

8. Here is the proof of Theorem 3.12

Shifts. We have

$$\begin{aligned}\widehat{z}_j &= \sum_{k=0}^{n-1} z_k \bar{w}^{kj} \quad \text{where } w = e^{i2\pi/n} \\ &= \sum_{k=0}^{n-1} y_{k+1} \bar{w}^{kj}\end{aligned}$$

We change indices by letting $\ell = k + 1$ to obtain

$$\begin{aligned}\widehat{z}_j &= \sum_{\ell=1}^n y_\ell \bar{w}^{(\ell-1)j} \\ &= \bar{w}^{-j} \sum_{\ell=1}^n y_\ell \bar{w}^{\ell j} \\ &= w^j \sum_{\ell=1}^n y_\ell \bar{w}^{\ell j}\end{aligned}$$

where the last equality holds since $\bar{w}^{-1} = w$ (recall that $w = e^{2\pi i/n}$). The last line is $w^j \widehat{y}_j$ as desired.

Convolutions. We must show

$$(x * y)_{k+n} = (x * y)_k$$

The left side is

$$(x * y)_{k+n} = \sum_{j=0}^{n-1} x_j y_{k+n-j}$$

Since $y_{k+n-j} = y_{k-j}$, we have

$$(x * y)_{k+n} = \sum_{j=0}^{n-1} x_j y_{k+n-j} = \sum_{j=0}^{n-1} x_j y_{k-j} = (x * y)_k$$

Convolution Theorem. We have

$$\begin{aligned}\widehat{(x * y)}_\ell &= \sum_{k=0}^{n-1} (x * y)_k \bar{w}^{\ell k} \\ &= \sum_{k=0}^{n-1} \sum_{j=0}^{n-1} x_{k-j} y_j \bar{w}^{\ell k} \\ &= \sum_{k=0}^{n-1} \sum_{j=0}^{n-1} x_{k-j} \bar{w}^{\ell(k-j)} y_j \bar{w}^{\ell j} \\ &= \sum_{j=0}^{n-1} \sum_{t=-j}^{n-1-j} x_t \bar{w}^{t\ell} y_j \bar{w}^{\ell j} \quad \text{with } t = k - j \\ &= \sum_{j=0}^{n-1} \sum_{t=0}^{n-1} x_t \bar{w}^{t\ell} y_j \bar{w}^{\ell j}\end{aligned}$$

where the last equation follows from the fact that $x_{q-n} = x_q$ and $w^{(q-n)\ell} = w^{q\ell}$ (so the inner sum from $t = -j$ to $t = -1$ is the same as the sum from $t = n - j$ to $n - 1$). The last expression on the right is

$$\sum_{t=0}^{n-1} x_t \bar{w}^{t\ell} \sum_{j=0}^{n-1} y_j \bar{w}^{\ell j} = \hat{x}_\ell \hat{y}_\ell$$

as desired.

Last Property. We have

$$\hat{y}_{n-k} = \sum_{j=0}^{n-1} y_j \bar{w}^{j(n-k)}$$

Now use the fact that $w^n = 1$ and that $\bar{w}^{-jk} = w^{jk}$ (again, recall that $w = e^{2\pi i/n}$). We obtain

$$\hat{y}_{n-k} = \sum_{j=0}^{n-1} y_j w^{jk}$$

Since the y_j are real, the right side is

$$\overline{\sum_{j=0}^{n-1} y_j \bar{w}^{jk}} = \widehat{y}_k$$

as desired.