# Solutions of Einstein's equations with cylindrical symmetry 

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- Relevant to describing the metric outside a cosmic string
- Perform calculations analogous to the textbook treatment of the spherically symmetric case

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- $\Phi, \Lambda$, and $\Psi$ are unknown functions of $r$ only
- Range of $\phi$ runs from 0 to $\phi_{*}$ (not necessarily $2 \pi$ ); can be made to be 0 to $2 \pi$ by rescaling $\phi$


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\begin{aligned}
& \Gamma_{t r}^{t}=\Gamma_{r t}^{t}=\Phi^{\prime} \\
& \Gamma^{r}{ }_{t t}=\Phi^{\prime} e^{2(\phi-\Lambda)} \\
& \Gamma_{r r}^{r}=\Lambda^{\prime} \\
& \Gamma^{r}{ }_{\phi \phi}=-r e^{-2 \Lambda} \\
& \Gamma^{r} \\
& { }_{z z}=-\Psi^{\prime} e^{2(\Psi-\Lambda)} \\
& \Gamma_{r \phi}^{\phi}=\Gamma^{\phi}{ }_{\phi r}=\frac{1}{r} \\
& \Gamma_{r z}^{z}=\Gamma_{z r}^{z}=\Psi^{\prime}
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- $R^{\alpha}{ }_{\beta \mu \nu}=\Gamma^{\alpha}{ }_{\beta \nu, \mu}-\Gamma^{\alpha}{ }_{\beta \mu, \nu}+\Gamma^{\alpha}{ }_{\sigma \mu} \Gamma^{\sigma}{ }_{\beta \nu}-\Gamma^{\alpha}{ }_{\sigma \nu} \Gamma^{\sigma}{ }_{\beta \mu}$


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\begin{aligned}
& R_{\phi \phi t}^{t}=r \Phi^{\prime} e^{-2 \Lambda} \\
& R_{\phi \phi r}^{r}=-r \Lambda^{\prime} e^{-2 \Lambda} \\
& R_{z z r}^{r}=\left(\Psi^{\prime \prime}+\Psi^{\prime 2}-\Psi^{\prime} \Lambda^{\prime}\right) e^{2(\Psi-\Lambda)} \\
& R_{t t r}^{r}=-\left(\Phi^{\prime \prime}+\Phi^{\prime 2}-\Phi^{\prime} \Lambda^{\prime}\right) e^{2(\Phi-\Lambda)} \\
& R_{t t z}^{z}=-\Psi^{\prime} \phi^{\prime} e^{2(\Phi-\Lambda)} \\
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& R_{t t}=\left(\Phi^{\prime \prime}+\Phi^{\prime 2}-\Phi^{\prime} \Lambda^{\prime}+\frac{1}{r} \phi^{\prime}+\Psi^{\prime} \Phi^{\prime}\right) e^{2(\Phi-\Lambda)} \\
& R_{r r}=-\Phi^{\prime \prime}-\Phi^{\prime 2}+\Phi^{\prime} \Lambda^{\prime}+\frac{1}{r} \Lambda^{\prime}-\Psi^{\prime \prime}-\psi^{\prime 2}+\Lambda^{\prime} \Psi^{\prime} \\
& R_{\phi \phi}=r\left(\Lambda^{\prime}-\Phi^{\prime}-\Psi^{\prime}\right) e^{-2 \Lambda} \\
& R_{z z}=-\left(\Psi^{\prime \prime}+\Psi^{\prime 2}-\Psi^{\prime} \Lambda^{\prime}+\Psi^{\prime} \Phi^{\prime}+\frac{1}{r} \Psi^{\prime}\right) e^{2(\Psi-\Lambda)}
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& \Lambda^{\prime}=\Omega^{\prime}=\Phi^{\prime}+\Psi^{\prime} \\
& \Phi^{\prime \prime}+\frac{1}{r} \Phi^{\prime}=0 \\
& \Psi^{\prime \prime}+\frac{1}{r} \Psi^{\prime}=0 \\
& \Phi^{\prime} \Psi^{\prime}+\frac{1}{r} \phi^{\prime}+\frac{1}{r} \Psi^{\prime}=0
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& \Phi=\ln \left(r^{a_{1}}\right)+\ln \left(a_{2}\right) \\
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- After rescaling $t, z$, and $r$ to absorb constants (and with $\left.a:=a_{1}, b:=b_{1}\right)$ : $d s^{2}=-r^{2 a} d t^{2}+r^{2(a+b)} d r^{2}+K^{2} r^{2} d \phi^{2}+r^{2 b} d z^{2}$ - leads to two conventions for $\phi$


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## Transforming Between Metric Forms

- When does $r=0$ in our metric form correspond to $\bar{r}=\infty$ in the alternate metric form?
- From our form to Garfinkle's: $\bar{r}=\frac{c}{a+b+1} r^{a+b+1}$
- Exponent of $r$ is negative whenever $a+b<-1$, which occurs whenever $b<-1$ (or, equivalently, $a<-1$ )
- From our form to Weyl's: $\bar{r}=\frac{c}{a+1} r^{a+1}$
- Exponent of $r$ is negative whenever $a<-1(b<-1)$
- From our form to Rosen's: $\bar{r}=\frac{c}{b+1} r^{b+1}$
- Exponent of $r$ is negative whenever $b<-1(a<-1)$
- Thus in all three cases, $r=0$ in our metric corresponds to $\bar{r}=\infty$ in the alternate metric whenever $a<-1 \Leftrightarrow b<-1 \Leftrightarrow a+b<-1$


## Future Work

- Look at $R_{\alpha \beta \gamma \delta} R^{\alpha \beta \gamma \delta}$


## Future Work

- Look at $R_{\alpha \beta \gamma \delta} R^{\alpha \beta \gamma \delta}$
- Try connecting exterior solution to string of finite radius


## Thank You

- Questions?

AND OVER THERE WE HAVE THE LABYRINTH GUARDS. ONE ALWAYS LIES, ONE ALWAYS TELLS THE TRUTH, AND ONE STABS PEOPLE WHO ASK TRICKY QUESTIONS.


