# Calculation of Highly Oscillatory Integrals by Quadrature Methods 

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## Outline

(9) Motivation

- Study of Vacuum Energy
- Oscillatory Integrals
- Earlier Literature
(2) Our Results
- Main Results
- Implementation


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## A model for Vacuum Energy

Our model of quantum vacuum energy density near the boundary has the form $\lambda z^{\alpha}$.


Figure: Steeply rising potential near the boundary

## Why are we interested?

Spectral analysis of the rising potential gives Energy momentum tensor:

$$
\begin{equation*}
\bar{T}(z)=\frac{1}{\pi^{3}} \int_{0}^{\infty} d \rho \int_{0}^{1} d u \sqrt{1-u^{2}} \cos (2 z \rho u-2 \delta(\rho u)) \tag{1}
\end{equation*}
$$

where,

$$
\delta(u)=\operatorname{ArcTan}\left(-u\left(\frac{\operatorname{AiryAi}\left(-u^{2}\right)}{\operatorname{Airy} \mathrm{Ai}^{\prime}\left(-u^{2}\right)}\right)\right)
$$

## What does it look like?



Figure: The oscillatory cosine function
left:u goes from 100 to 100.0002
right: u goes from 100 to 100.004

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## $\bar{T}(z)$ is highly oscillatory.

- Takes hours, if not days to calculate.
- Only for $\alpha=1$
- We need to check for higher values of $\alpha$.
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## Newton-Cotes Rule

## - Trapezoidal rule




## Figure: Plot showing integration by trapezoidal rule



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- Trapezoidal rule
- Simpson's rule



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$$
\int_{a}^{b} f(x) d x \approx c_{1} f(a)+c_{2} f(b)=\frac{(b-a)}{2}(f(a)+f(b))
$$

## Gauss-Quadrature

- $\int_{a}^{b} f(x) d x \approx c_{1} f\left(x_{1}\right)+c_{2} f\left(x_{2}\right)$.
- Here, $c_{1}, c_{2}, x_{1}$, and $x_{2}$ are all unknowns.
- In this case, these four constants are found by integrating third order polynomials and equating the coefficients.



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$$
\begin{gathered}
x_{1}=\frac{b-a-1}{2} \frac{b+a}{\sqrt{3}}+\frac{b+a}{2} \\
x_{2}=\frac{b-a}{2} \frac{1}{\sqrt{3}}+\frac{b+a}{2} \\
c_{1}=\frac{b-a}{2}, \text { and } c_{2}=\frac{b-a}{2} .
\end{gathered}
$$

## Filon's method



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$$
\int_{a}^{b} f(x) \sin \omega x d x \text { and } \int_{0}^{\infty} \frac{f(x)}{x} \sin \omega x
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m_{\mu}\left(x_{\mu}\right)=f\left(x_{\mu}\right), m_{\mu+1}\left(x_{\mu+1}\right)=f\left(x_{\mu+1}\right), \text { and } m_{\mu+2}\left(x_{\mu+2}\right)=f\left(x_{\mu+2}\right)
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## ClenshawCurtis Method

## C.W.Clenshaw and A.R. Curtis in 1960. Expand $f(x)$ in

 Chebyshev polynomials.$f(x)=F(t)=\frac{1}{2} a_{0}+a_{1} T_{1}(t)+a_{2} T_{2}(t)+\ldots+\frac{1}{2} a_{n} T_{n}(t),(a \leq x \leq b)$ (2)
where,


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\end{equation*}
$$

where,

$$
\begin{equation*}
T_{n}(t)=\cos \left(\mathrm{n} \cos ^{-1}(t), t=\frac{2 x-(b+a)}{b-a}\right. \tag{3}
\end{equation*}
$$

and this eventually reduces to

$$
\begin{equation*}
f(x)=\frac{a_{0}}{2} T_{0}(x)+\sum_{n=1}^{\infty} a_{n} T_{n}(x), x_{n}=\cos \left(\frac{n \pi}{N}\right) \tag{4}
\end{equation*}
$$

## Levin-Iserles' Method

## Improvement over Filon's method




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\begin{align*}
Q_{2}^{F}[f] & =\left(-\frac{1}{i \omega}-6 \frac{1+e^{i \omega}}{i \omega^{3}}+12 \frac{1-e^{i \omega}}{\omega^{4}}\right) f(0)  \tag{5}\\
& +\left(\frac{e^{i \omega}}{i \omega}+6 \frac{1+e^{i \omega}}{i \omega^{3}}-12 \frac{1-e^{i \omega}}{\omega^{4}}\right) f(1) \\
& +\left(-\frac{1}{\omega^{2}}-2 \frac{2+e^{i \omega}}{i \omega^{3}}+6 \frac{1-e^{i \omega}}{\omega^{4}}\right) f^{\prime}(0) \\
& +\left(\frac{e^{i \omega}}{\omega^{2}}-2 \frac{1+e^{i \omega}}{i \omega^{3}}+6 \frac{1-e^{i \omega}}{\omega^{4}}\right) f^{\prime}(1)
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## Was it worth the time?

- Yes, and No.
- Iserles' method did not work for our $\bar{T}(z)$ integral.
- Were able to calculate integrals much faster.
- Not very consistent.


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## What quadrature method to choose?

- Levin-Iserles' method seems more promising.
- Clenshawcurtis' qudrature method also works well.


Figure: $\bar{T}(z)$ using Levin and Clenshawcurtis method

## Summary

- Highly oscillatory integrals can be calculatedmuch faster than by conventional methods.
- choose methods judiciously.
- Reduce error for integrands with large frequency.
- Outlook
- Not enough data for conclusion.
- Check for higher values of $\alpha$.


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## For Further Reading I

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