

Construction of suitable weak solutions for the 3D incompressible NSEs

Jean-Luc Guermond

Department of Mathematics
Texas A&M University, College Station, TX
and LIMSI, Orsay, FR

Rencontres Niçoises de Mécanique des Fluides
Laboratoire J.A.Dieudonné
Nice University
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Outline

1 BASIC FACTS ABOUT THE 3D NSEs



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- 2 SUITABLE WEAK SOLUTIONS



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BASIC FACTS ABOUT THE NSE



Claude Louis Marie Henri
Navier



George Gabriel Stokes

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INCOMPRESSIBLE 3D NAVIER–STOKES EQUATIONS

- u : velocity, p : pressure



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$$\left\{ \begin{array}{l} \partial_t u + u \cdot \nabla u + \nabla p - \nu \nabla^2 u = f \quad \text{in } \Omega \\ \nabla \cdot u = 0 \quad \text{in } \Omega, \\ u|_{\Gamma} = 0 \quad \text{or } u \text{ is periodic,} \\ u|_{t=0} = u_0, \end{array} \right.$$



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- u_0 is the initial data.
- f a source term.
- ρ is chosen equal to unity.
- ν is viscosity (inverse of Reynolds number).



EXISTENCE

- J. Leray (1934): introduces the notion of **turbulent solution**.
A turbulent solution is a **weak solution** in
 $u \in L^2(0, T; \mathbf{H}^1(\Omega)) \cap L^\infty(0, T; \mathbf{L}^2(\Omega))$
+ global energy inequality:

$$\frac{1}{2} \|u\|_{\mathbf{L}^2}^2 + \nu \int_0^T \|\nabla u\|_{L^2}^2 dt \leq \frac{1}{2} \|u_0\|_{\mathbf{L}^2}^2 + \int_0^T \int_{\Omega} f \cdot u dx dt$$



EXISTENCE

- J. Leray uses **mollification** to prove existence:

$$\psi \in D(\mathbb{R}^3), \psi \geq 0, \int_{\mathbb{R}^3} \psi = 1, \psi_\epsilon(x) = \frac{1}{\epsilon} \psi\left(\frac{x}{\epsilon}\right).$$

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- E. Hopf (1951) *et al.* uses the **Galerkin technique** to prove existence.

Idea: Discretize and pass to the limit.



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- Are weak solutions **unique** in the large?



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- Uniqueness for small data (essentially in $\mathbf{H}^{1/2}(\Omega)$, Fujita-Kato (1964)) or large highly oscillatory data (become small fast by diffusion).
- Uniqueness in more regular classes (with short time existence, Prodi-Serrin (1959-1963))



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 \Rightarrow evolution under NSE: $B_{\infty,\infty}^{-1} \rightarrow C([0, T]; B_{\infty,\infty}^{-1})$ is not continuous.



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- Cheskidov, Shvydkoy (Arxiv, 2009): The evolution under NSE $H^s \cap B_{\infty,\infty}^{-1} \rightarrow C([0, T]; B_{\infty,\infty}^{-1})$ is not continuous, $\forall s < \frac{1}{2}$.



SUITABLE WEAK SOLUTION

Definition (V. Scheffer (1976))

A NS weak solution is said to be **suitable weak solutions** iff (u, p) is a weak solution and

$$\partial_t\left(\frac{1}{2}u^2\right) + \nabla \cdot \left(u\left(\frac{1}{2}u^2 + p\right)\right) - \nu \nabla^2\left(\frac{1}{2}u^2\right) + \nu(\nabla u)^2 - f \cdot u \leq 0.$$

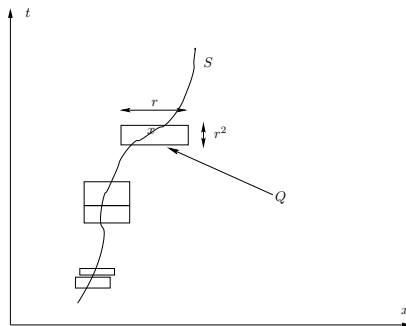
in $\mathcal{D}'((0, T) \times \Omega)$



SUITABLE WEAK SOLUTION

- Singular set: S

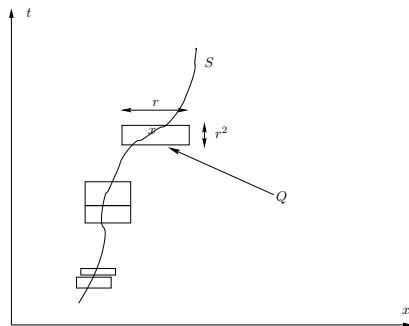
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- $\mathcal{P}^1(S) = \lim_{\delta \rightarrow 0^+} \inf \{ \sum r_i^1, : S \subset \cup Q(M_i, r_i), r_i < \delta \}$



SUITABLE WEAK SOLUTION

Theorem (Caffarelli-Kohn-Nirenberg (1982))

If (u, p) is a *suitable weak solutions*, then $\mathcal{P}^1(S) = 0$



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Best partial regularity theorem to date.



CONSTRUCTION OF SUITABLE SOLUTIONS

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EX 1: CONSTRUCTION BY MOLLIFICATION

- Leray's mollification

$$\partial_t u_\epsilon + (\psi_\epsilon * u_\epsilon) \cdot \nabla u_\epsilon + \nabla p_\epsilon - \nu \nabla^2 u_\epsilon = f$$



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Theorem (Leray (1934), Duchon-Robert (2000))

Unique weak solution for all $t > 0$, and $u_\epsilon \rightharpoonup u$ (up to subsequences) and u is *suitable*.



EX 2: CONSTRUCTION BY MOLLIFICATION/NLGM

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- Set $\mathbf{X}_\varepsilon = \dot{\mathbf{P}}_{N_\varepsilon}$ (velocity space).
- Let $P_\varepsilon : \dot{\mathbf{L}}^2(\Omega) \rightarrow \mathbf{X}_\varepsilon$ be L^2 -projection.

$$\dot{\mathbf{L}}^2(\Omega) = \mathbf{X}_\varepsilon \oplus (\mathbf{X}_\varepsilon)^\perp$$



EX 2: CONSTRUCTION BY MOLLIFICATION/NLGM

- Solve for u_ϵ, p_ϵ s.t.

$$\begin{cases} \partial_t(P_\epsilon u_\epsilon) + P_\epsilon u_\epsilon \cdot \nabla u_\epsilon + \nabla p_\epsilon - \nu \nabla^2 u_\epsilon = \mathbf{f} & \text{in } \Omega \\ \nabla \cdot u_\epsilon = 0 & \text{in } \Omega, \\ u_\epsilon \text{ is periodic, } u_\epsilon|_{t=0} = u_0, \end{cases}$$



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Theorem

For all $\epsilon > 0$, problem is well-posed (existence + uniqueness).
 $u_\epsilon \rightarrow u, p_\epsilon \rightarrow p$ as $\epsilon \rightarrow 0$ (in appropriate spaces, up to subsequences), u and p are **suitable** weak solution to N.S.



EX 3: CONSTRUCTION BY HYPERVISCOSITY

- Add **vanishing hyperviscosity** (Lions (1959)). Ω is the d -torus, d is the space dimension.

$$\left\{ \begin{array}{l} \partial_t u_\epsilon + u_\epsilon \cdot \nabla u_\epsilon + \nabla p_\epsilon - \nu \nabla^2 u_\epsilon + \epsilon^{2\alpha} (-\nabla^2)^\alpha u_\epsilon = f, \\ \nabla \cdot u_\epsilon = 0 \\ u_\epsilon \text{ is periodic,} \\ u_\epsilon|_{t=0} = u_0. \end{array} \right.$$



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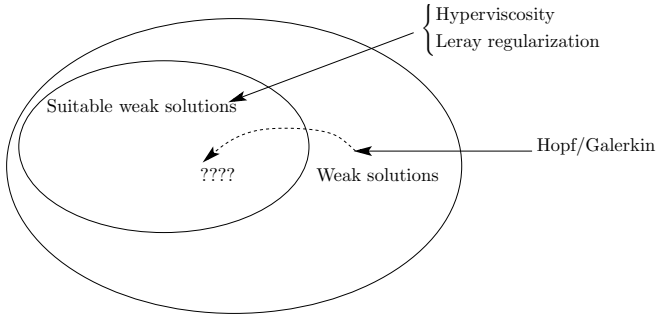
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Theorem (Lions (1959), Beirão da Veiga (1985))

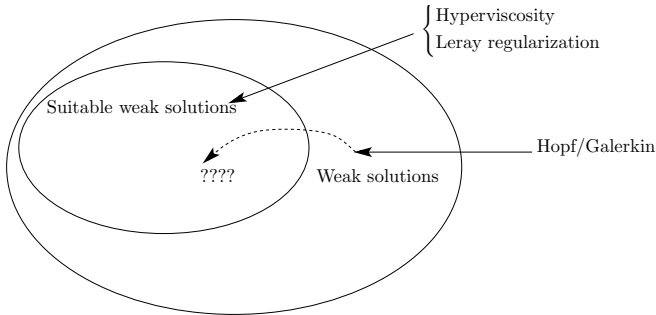
*Unique weak solution for all $t > 0$ if $\alpha > \frac{d+2}{4}$, and $u_\epsilon \xrightarrow{\epsilon \rightarrow 0} u$ (up to subsequences) and u is **suitable**.*



MORE TRACTABLE QUESTIONS

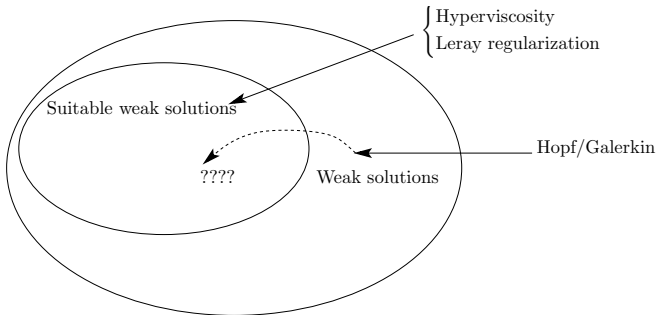


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Q1: Is the set of suitable solutions a **proper** subset of weak solutions?

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Q2: Do the Galerkin solutions end up to be **suitable** after all?



GALERKIN APPROX IN TORUS



Boris Grigorievich
Galerkin

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- Assume there is $c > 0$ independent of h such that

$$\forall q_h \in M_h, \quad \sup_{0 \neq v_h \in X_h} \frac{(\nabla q_h, v_h)}{\|v_h\|_{\mathbf{L}^2}} \geq c \|\nabla q_h\|_{\mathbf{L}^2}.$$



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- Modify the nonlinear term as follows:

$$b_h(u, v, v) = \begin{cases} (u \cdot \nabla u + \frac{1}{2} u \nabla \cdot u, v) & \text{(Temam, 1967)} \\ ((\nabla \times u) \times u + \frac{1}{2} \nabla(\mathcal{K}_h(u^2)), v) \end{cases}$$

where $\mathcal{K}_h : L^2(\Omega) \longrightarrow M_h$, linear L^2 -stable approximation operator.



HYPOTHESES/DEFINITIONS (ctd.)

Definition (Discrete commutator property)

There is an operator $P_h \in \mathcal{L}(H_{\#}^1(\Omega); X_h)$ such that for all ϕ in $W_{\#}^{2,\infty}(\Omega)$ and all $v_h \in X_h$

$$\|\phi v_h - P_h(\phi v_h)\|_{H^l} \leq c h^{1+m-l} \|v_h\|_{H^m} \|\phi\|_{W^{m+1,\infty}}, \quad 0 \leq l \leq m \leq 1.$$



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- Similar hypothesis on pressure space M_h .
- Define $\Phi : \psi \mapsto \phi\psi$, then

$$P_h(\phi v_h) - \phi v_h = P_h(\Phi(v_h)) - \Phi(P_h(v_h)) := [P_h, \Phi](v_h)$$



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- For instance we want: $\| [P_h, \Phi] v_h \|_{L^2} \leq c h \| v_h \|_{L^2} \| \phi \|_{W^{1,\infty}}$,
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i.e. $\| [P_h, \Phi] \|_{\mathcal{L}(\mathbb{P}_N \cap L^2, L^2)} \leq c h \| \phi \|_{W^{1,\infty}}$.
- FE and wavelet-based approximation spaces have the discrete commutator property



HYPOTHESES/DEFINITIONS (ctd.)

- Localization is the key property:

$$\|\psi - P_h\psi\|_{L^2(K)} \leq ch_K \|\nabla\psi\|_{L^2(\Delta_K)}$$

K is a cell, $h_K := \text{diam}(K)$, Δ_K is a patch containing K ,
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- Second key: P_h is a projection

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- Fourier-based approximation spaces **do not have the discrete commutator property** (No localization for interpolation properties).



GALERKIN FORMULATION

- Seek $u_h \in \mathcal{C}^1([0, T]; X_h)$ and $p_h \in \mathcal{C}^0([0, T]; M_h)$ such that for all $v_h \in X_h$, all $q_h \in M_h$, and all $t \in [0, T]$

$$\begin{cases} (\partial_t u_h, v) + b_h(u_h, u_h, v) - (p_h, \nabla \cdot v) + \nu(\nabla u_h, \nabla v) = \langle f, v \rangle, \\ (\nabla \cdot u_h, q) = 0, \\ u_h|_{t=0} = \mathcal{I}_h u_0, \end{cases}$$

where $\mathcal{I}_h : L^2(\Omega) \longrightarrow V_h$, L^2 -stable interpolation operator.



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Theorem (Guermond (2006))

*Under the above hypotheses, if X_h and M_h have **the discrete commutator property**, the couple (u_h, p_h) convergences to a **suitable** solution to NS.*



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- Question was **open** since Scheffer (1977).



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Question: Does the result hold for Dirichlet BCs?



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$$\begin{aligned}(u_h \cdot \nabla u_h, P_h(\phi u_h)) &= (u_h \cdot \nabla u_h, \phi u_h) + (u_h \cdot \nabla u_h, P_h(\phi u_h) - \phi u_h) \\ &= (u_h \cdot \nabla (\frac{1}{2} u_h^2), \phi) + (u_h \cdot \nabla u_h, [P_h, \Phi] u_h) \\ &\approx (\frac{1}{2} \nabla \cdot (u_h^2 u_h), \phi) + (u_h \cdot \nabla u_h, [P_h, \Phi] u_h) \\ &= -(\frac{1}{2} u_h^2 u_h, \nabla \phi) + (u_h \cdot \nabla u_h, [P_h, \Phi] u_h).\end{aligned}$$



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- Use the **discrete commutator** property to pass to the limit.



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- Use the **discrete commutator** property to pass to the limit.
- Idem form the term $u_h p_h$



OUTLINE



Johann Peter Gustav
Lejeune Dirichlet

- 1 BASIC FACTS ABOUT THE 3D NSEs
- 2 SUITABLE WEAK SOLUTIONS
- 3 GALERKIN IN TORUS
- 4 GALERKIN + DIRICHLET
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HYPOTHESES/DEFINITIONS

- Finite element spaces, $X_h \subset \mathbf{H}_{\#}^1(\Omega)$ for **velocity** and $M_h \subset H_{\#}^1(\Omega)$ for **pressure**.



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- Modify the nonlinear term as follows:

$$b_h(u, v, v) = \begin{cases} (u \cdot \nabla u + \frac{1}{2} u \nabla \cdot u, v) & \text{(Temam, 1967)} \\ ((\nabla \times u) \times u + \frac{1}{2} \nabla(\mathcal{K}_h(u^2)), v) \end{cases}$$

where $\mathcal{K}_h : L^2(\Omega) \longrightarrow M_h$, linear L^2 -stable approximation operator.



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PRELIMINARY RESULTS

Lemma

$\exists c_l > 0$ (*non-increasing function*) $\exists c_u > 0$ (*non-decreasing function*), independent of h :

$$c_l(|s|) \|v_h\|_{\tilde{\mathbf{H}}_0^s} \leq \|v_h\|_{\mathbf{V}_h^s} \leq c_u(|s|) \|v_h\|_{\tilde{\mathbf{H}}_0^s}, \quad \begin{cases} -\frac{1}{2} < s < \frac{3}{2}, & \text{lower,} \\ -\frac{3}{2} < s < \frac{3}{2}, & \text{upper} \end{cases}$$



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and for all $s \in (-\frac{3}{2}, 0]$

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GALERKIN FORMULATION

- Seek $u_h \in \mathcal{C}^1([0, T]; X_h)$ and $p_h \in \mathcal{C}^0([0, T]; M_h)$ such that for all $v_h \in X_h$, all $q_h \in M_h$, and all $t \in [0, T]$

$$\begin{cases} (\partial_t u_h, v) + b_h(u_h, u_h, v) - (p_h, \nabla \cdot v) + \nu(\nabla u_h, \nabla v) = \langle f, v \rangle, \\ (\nabla \cdot u_h, q) = 0, \\ u_h|_{t=0} = \mathcal{I}_h u_0, \end{cases}$$

where $\mathcal{I}_h : L^2(\Omega) \longrightarrow V_h$, L^2 -stable interpolation operator.



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Theorem (Guermond (2007))

*Under the above hypotheses, if X_h and M_h have **the discrete commutator property**, the couple (u_h, p_h) convergences to a **suitable solution to NS**.*



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⇒ Hopf and Leray solutions are suitable



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- **The key difficulty:** Cannot use $\nabla^2 p = -\partial_j(u_i \partial_i u_j)$ to get pressure (no control on $\partial_n p$ at the boundary!)



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- Alternative: (control on $\partial_t u_h$ and $\nabla_h^2 u_h$) + (Inf-sup) \Rightarrow (control on p_h).
- **Key:** Use Fourier in time to derive estimates on $\partial_t u_h$ in negative norms.



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There is c independent of h so that,

$$\|\partial_t u_h\|_{H^{\tau-1}((0,T); \mathbf{H}^{-\alpha}(\Omega))} + \|u_h\|_{H^{\tau}((0,T); \mathbf{H}^{-\alpha}(\Omega))} \leq c,$$

for all $\alpha \in [\frac{1}{4}, \frac{1}{2})$ and for all $\tau < \bar{\tau} := \frac{2}{5}(1 + \alpha)$.



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$$\text{SW (1986)} \Rightarrow \partial u_t \in H^{-\frac{1}{2}-\epsilon}(0, T; \mathbf{H}^{-\frac{1}{2}-\epsilon}(\Omega))$$

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- Conjecture: $\partial u_t \in L^{\frac{10}{9}}(0, T; \mathbf{L}^{\frac{3}{2}}\Omega)$.



THE MAIN RESULT (Sketch of proof)

Corollary

There is c independent of h such that for $s \in [\frac{3}{10}, \frac{1}{2}]$

$$\|p_h\|_{H^{-r}((0,T);H^s(\Omega))} \leq c,$$

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Enough a priori estimates to pass to the limit (+discrete commutator).



FOURIER APPROX IN TORUS



Joseph Fourier

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- Same problem if one uses special bases: e.g. eigenvectors of Stokes operator.
- **Bottom line:** The commutator property is **not strong enough**.



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Lemma (Negative result)

$\exists c_1, c_2, \forall \phi \in W^{1,\infty}(\Omega)$

$$c_1 N^{-1} \|\phi\|_{W^{1,\infty}} \leq \|[P_N, \Phi]\|_{\mathcal{L}(H^1, L^2)} \leq c_2 N^{-1} \|\phi\|_{W^{1,\infty}}.$$



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- $\Rightarrow \|[P_N, \Phi]u_N\|_{L^\infty} \rightarrow 0$ if $u_N \rightarrow u$ in H^1 strong, (**weak convergence is not enough!**)



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- Does not go to 0, unless $u_N \rightarrow u$ in $L^2(0, T; \mathbf{H}^1)$ **strong** (energy equality).



CONJECTURE

- Result would be true if for all $k_0 > 0$,

$$\lim_{N \rightarrow +\infty} \int_0^T \sup_{N < k < N + k_0} (\widehat{u_N \cdot \nabla u_N})_k^2 dt = 0$$



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- Note that if u Leray solution

$$\lim_{k \rightarrow +\infty} \int_0^T (\widehat{u \cdot \nabla u})_k^2 dt = 0$$

Consequence of Riemann Lebesgue Lemma,
 $\|u \cdot \nabla u\|_{L^1(\Omega)}$ is in $L^2(0, T)$.



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