



ILL: 5725059

Call bound
Number:
Location: ifm
Maxcost: \$50ifm

DateReq: 5/7/2004 Yes
Date Rec: 5/11/2004 No
Borrower: TXA Conditional
LenderString: IUA.LHL.*DGW,LIN,LIN

Title: Inverse problems in engineering.

Author:

Edition:

Imprint: New York, N.Y. : Gordon & Breach Science Publishers

Article: A. Chakrabarti, P. Daripa, and S. Roy: An Inverse Problem for the Helmholtz Equation Involving Two Semi-infinite Fluids

Vol: 10

No.: 3

Pages: 203-214

Date: 2002

NOTICE
THIS MATERIAL MAY
BE PROTECTED BY
COPYRIGHT LAW
(TITLE 17, US CODE)

Borrowing

Notes:

Fax:

ILL: 5725059 :Borrower: TXA :ReqDate: 20040507 :NeedBefore: 20040606
:Status: IN PROCESS 20040510 :RecDate: :RenewalReq:
:OCLC: 27568317 :Source: OCLCILL :DueDate: :NewDueDate:
:Lender: IUA, LHL, *DGW, LIN, LIN
:CALLNO: :TITLE: Inverse problems in engineering. :IMPRINT: New York, N.Y. :
Gordon & Breach Science Publishers :ARTICLE: A. Chakrabarti, P. Daripa, and S.
Roy: An Inverse Problem for the Helmholtz Equation Involving Two Semi-infinite
Fluids :VOL: 10 :NO: 3 :DATE: 2002 :PAGES: 203-
214 :VERIFIED: <TN:431368>OCLC ISSN: 1068-2767 [Format: Serial]
:PATRON: Daripa, Prabir :SHIP TO: Interlibrary Services/TAMU Libraries/5000
TAMUS/College Station, TX 77843-5000 :BILL TO: SAME...FEIN 746000531 :SHIP
VIA: Ariel 165.91.220.14 :MAXCOST: \$50ifm :COPYRT COMPLIANCE: CCG :FAX:
(979)862-4759; ARIEL 165.91.220.14 :E-MAIL: ilsshare@lib-gw.tamu.edu :BILLING
NOTES: BRI USER CODE 51-1281;CAI # DD000806; LHL D10225 :AFFILIATION:
ARL, BTP, HARLiC, TEXSHARE, TExpress 41 HOU, TEL: (979) 845-5641 :LENDING
CHARGES: :SHIPPED: :SHIP INSURANCE: :LENDING
RESTRICTIONS: :LENDING NOTES: :RETURN TO: :RETURN VIA:

ShipVia: Ariel 165.91.220



NeedBy: 6/6/2004

Return To:

Gelman Library ILL Unit, Rm. 103
The George Washington University
2130 H Street, NW
Washington, DC 20052

ILL: 5725059 Borrower: TXA
Req Date: 5/7/2004 OCLC #: 27568317

Patron: Daripa, Prabir

Author:

Title: Inverse problems in engineering.

Article: A. Chakrabarti, P. Daripa, and S. Roy: An Inverse Problem for the Helmholtz Equation Involving Two Semi-infinite Fluids

Vol.: 10 No.: 3

Date: 2002 Pages: 203-214

Verified: <TN:431368>OCLC ISSN: 1068-2767 [For

Maxcost: \$50ifm Due Date:

Lending Notes:

Bor Notes:

Ship To:

Interlibrary Services

TAMU Libraries

5000 TAMUS

College Station, TX 77843-5000

AN INVERSE PROBLEM FOR THE HELMHOLTZ EQUATION INVOLVING TWO SEMI-INFINITE FLUIDS

A. CHAKRABARTI^{a,*}, PRABIR DARIPA^{b,†} and S. ROY^c

^a*Department of Mathematics, NJIT, University Heights, Newark, NJ-07102;*

^b*Department of Mathematics, Texas A&M University, College Station, TX-77843-3368, USA;*

^c*Department of Mathematics, IIT-Madras, Chennai, 600036, India*

(Received 30 July 2001; In final form 14 November 2001)

An analytical method is developed for solving an inverse problem for Helmholtz's equation associated with two semi-infinite incompressible fluids of different variable refractive indices, separated by a plane interface. The unknowns of the inverse problem are: (i) the refractive indices of the two fluids, (ii) the ratio of the densities of the two fluids, and (iii) the strength of an acoustic source assumed to be situated at the interface of the two fluids. These are determined from the pressure on the interface produced by the acoustic source. The effect of the surface tension force at the interface is taken into account in this paper.

The application of the proposed analytical method to solve the inverse problem is also illustrated with several examples. In particular, exact solutions of two direct problems are first derived using standard classical methods which are then used in our proposed inverse method to recover the unknowns of the corresponding inverse problems. The results are found to be in excellent agreement.

1 INTRODUCTION

The inverse problem concerning the determination of the triple $\{n(z), \alpha, h_0\}$, arising in acoustics, is considered where the three unknowns (the triple) of the problem are as described below:

- (i) $n(z)$, the variable refraction coefficient of the medium (see Section 2) consisting of two semi-infinite fluids of constant densities ρ_{0+} and ρ_{0-} occupying the regions $z > 0$ and $z < 0$ respectively, where the z -axis is taken normal to the interface and a rectangular coordinate system (x, y, z) is employed.
- (ii) $\alpha = \rho_{0-}/\rho_{0+} > 1$.
- (iii) h_0 , a constant, determining the strength of a time-harmonic ($e^{-i\omega t}$) sound-source of creating fluid assumed to be situated at $z = 0^+$, very near the origin of coordinates on the interface $z = 0$.

*Permanent Address: Indian Institute of Science, Bangalore, 560012, India.

†Corresponding author. E-mail: prabir.daripa@math.tamu.edu

The known data of the problem under consideration are the acoustic pressure functions $p(x, y, 0^\pm)$ at the interface where $p(x, y, 0^\pm) = \lim_{z \rightarrow 0^\pm} p(x, y, z)$. These are discussed in detail in Section 2.

Similar inverse problems have been considered for solution by Ramm [1] among other authors using a Green's function approach when the effect of surface tension at the interface of the two fluids is neglected. The present paper deals with the above inverse problem when the effect of the surface tension at the interface of the two fluids is also taken into account. Below we briefly summarize the method of solving this inverse problem which is given in more detail in the main body of the paper.

In the present context, the corresponding direct problem is concerned with the determination of the pressure functions in the fluid regions. It is assumed in our approach below that this direct problem possesses a unique solution for the class of functions $n(z)$, having at most an algebraic growth at infinity. Under this assumption, the direct problem is cast into an integral equation of the second kind (see Mikhlin [2]) for the determination of the pressure functions. Using a technique for obtaining a series solution of this integral equation for the pressure functions which is valid for sufficiently small values of the wave number $k = \omega/c_0$ (c_0 representing the constant sound speed in free-space), the inverse problem of determination of $n(z)$ is cast into two independent Laplace-inversion problems, one for each of the two regions: $z > 0$ and $z < 0$. The other two unknowns α and h_0 (in the triple $\{n(z), \alpha, h_0\}$) are determined by using a straight-forward limiting procedure.

The paper is organized as follows. In Section 2, we have presented the basic equations and boundary conditions of the direct problem. Section 3 is devoted to the detailed analysis to determine the solution of the inverse problem. For validation purposes of our inverse method, two special cases of the direct problem are considered which are as follows: (i) $n(z) = n_0$ for $z > 0$, and $n(z) = n_1$ for $z < 0$ where n_0 and n_1 are two different known constants; and (ii) $n(z) = n_0$ (a known constant) for $z > 0$, and $n(z) = \exp(2z)$ for $z < 0$. These two direct problems are solved in a straight-forward manner to obtain the input data for the inverse problem. With these as the input data, the results of the corresponding inverse problems for the unknowns are found to be in excellent agreement with their exact values when the various formulae derived in Section 3 for the solution of the inverse problems are utilized.

2 MATHEMATICAL FORMULATION

We employ the notations used by Jones [3] with p_0 and ρ_0 as constant pressure and density respectively of a single fluid medium and with $p_0(1+p)$ and $\rho_0(1+\rho)$ as the corresponding perturbed quantities, caused due to a sound source of creating fluid and situated at a point (x_0, y_0, z_0) . Then the basic equations of acoustics can be written down in the following linearized form (neglecting gravity, viscosity and thermal conductivity, and assuming that p and the perturbed fluid velocity \mathbf{v} are small quantities):

$$\rho_0 \frac{\partial \rho}{\partial t} + \rho_0 \operatorname{div} \mathbf{v} = -\frac{\partial m}{\partial t} \delta(x - x_0) \delta(y - y_0) \delta(z - z_0), \quad (1)$$

where 'c' is the speed of sound, $\delta(x)$ denoting the Dirac delta function.

If the medium is homogeneous, the densities ρ_{0+} and ρ_{0-} are constant and time-independent. In the present case, the above equations are modified to include the time-dependence of the sound speed $c(z)$ in the z -direction only. The boundary conditions at $z = 0$ are $p(x, y, 0) = m_0 \exp(-i\omega t)$, $\rho(x, y, 0) = 0$, $\alpha(x, y, 0) = 0$, $h_0(x, y, 0) = 0$. Eqs. (1), (2), and (3) are solved for the pressure function $p(x, y, z)$.

where $n(z) = c_0^2/c^2(z)$ is the refractive index in free-space (or the refractive index $c_0/c(z)$ of the medium). The coefficient of the perturbation in Eq. (4), $\alpha(x, y, z)$, is zero for $z < 0$ and that for $z > 0$ is given by

The kinematic boundary conditions at $z = 0$ are respectively the

- (a) The particle velocity is continuous across the interface.
- (b) The jump in the normal component of the surface-tension is zero.

The linearized boundary conditions at $z = 0$ are (for details) in the

where $\alpha = \rho_{0+}/\rho_{0-}$. In the sequel, we have used the regions $z > 0$ and $z < 0$.

$$\rho_0 \frac{\partial \mathbf{v}}{\partial t} = -p_0 \text{grad } p, \quad (2)$$

$$p = c^2 \frac{\rho_0 \rho}{\rho_0}, \quad (3)$$

where 'c' is the speed of sound and $m(t)$ represents the mass created at the source at time t , $\delta(x)$ denoting Dirac's delta function (considering isentropic flows only).

If the medium under consideration consists of two semi-infinite fluids of constant densities ρ_{0+} and ρ_{0-} occupying the regions $z > 0$ and $z < 0$ respectively, and if the time-dependence is assumed to be of the type $\exp(-i\omega t)$ (i.e. time-harmonic case), the above equations are to be used with appropriate changes in the various quantities involved. In addition to this, we assume that the sound speed $c(z)$ varies in the z -direction only. The source of sound is taken to be of the form: $m(t) = m_0 \exp(-i\omega t)$, (m_0 denoting the strength of the source) located at $(x_0, y_0, z_0) = (0, 0, a)$, $a(> 0)$. Then, by means of a standard elimination procedure adapted to the Eqs. (1), (2), and (3), we obtain the following partial differential equation for the pressure-function $p(x, y, z)$.

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} + k^2 n(z) \right) p = -h_0 k^2 \delta(x) \delta(y) \delta(z - a), \quad (4)$$

where $n(z) = c_0^2/c^2(z)$, $k = \omega/c_0$, and $h_0 = m_0 c_0^2/\rho_{0+}$, with p_{0+} representing the constant unperturbed pressure in the region $z > 0$, and c_0 representing the constant sound speed in free-space (or any reference medium). The function $n(z)$, related to the refractive index $c_0/c(z)$ of the fluid medium under consideration, will be called the refraction coefficient of the medium. It is worth pointing out that the functions $n(z)$ and p , appearing in Eq. (4), actually represent two different functions in the two regions $z > 0$ and $z < 0$ and that there are two equations built-in in Eq. (4).

The kinematic and dynamic interface conditions (see Drazin [4], Lamb [5]) are respectively the following.

- The particle-displacement at the interface $z=0$ must be continuous, and
- The jump in normal stress (pressure) at the interface must be balanced with the surface-tension force T at the interface.

The linearized versions of these interface conditions can be expressed (See Lamb [5] for details) in the following forms:

$$\alpha \frac{\partial p_+}{\partial z} = \frac{\partial p_-}{\partial z}, \quad \text{on } z = 0, \quad (5)$$

$$p_+ - p_- = \frac{-\varepsilon}{k^2} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \frac{\partial p_+}{\partial z}, \quad \text{on } z = 0, \quad (6)$$

where $\alpha = \rho_{0-}/\rho_{0+}$, and $\varepsilon = T/c_0^2 \rho_{0+}$ are two parameters of the problem. Here, and in the sequel, we have used the symbols p_+ and p_- for the pressure function p in the two regions $z > 0$ and $z < 0$ respectively.

In addition to the above interface conditions, the functions p_+ and p_- must also satisfy the radiation conditions at infinity in order to ensure the uniqueness of the solution of the boundary value problem posed by the Eqs. (4)–(6). We assume that the unknown function $n(z)$ tends to a constant limit as $|z| \rightarrow \infty$ so that the radiation conditions are applicable in the two regions $z > 0$ and $z < 0$.

This concludes the formulation of the direct problem of acoustics where $n(z)$, α , h_0 , ε , a and k are known and the pressure function is unknown. But, what we wish to consider here is the inverse problem of determining the triple $\{n(z), \alpha, h_0\}$ from a knowledge of the data p_+ and p_- at the interface $z=0$ when ε , a and k are given. In the next section, we have presented an analytical method of solution of this inverse problem under the assumption that $a \rightarrow 0^+$, i.e., the source is placed at the interface.

3 METHOD OF SOLUTION

We first non-dimensionalize the partial differential Eq. (4) and the conditions (5) and (6) by using the non-dimensional coordinates \bar{x} , \bar{y} , \bar{z} and the non-dimensional parameters \bar{k} , $\bar{\varepsilon}$ and \bar{h}_0 as defined by relations

$$\bar{x} = x/L, \quad \bar{y} = y/L, \quad \bar{z} = z/L, \quad \bar{k} = kL, \quad \bar{\varepsilon} = \varepsilon/L, \quad \bar{h}_0 = h_0/L^3 \quad \text{and} \quad \bar{a} = a/L,$$

L representing a typical length-scale. We then drop the bars everywhere and proceed to solve the problem as described below.

Using a double Fourier transform as defined by

$$\bar{p}(z) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} p(x, y, z) \exp[i(\lambda_1 x + \lambda_2 y)] dx dy, \quad (\lambda^2 = \lambda_1^2 + \lambda_2^2) \quad (7)$$

the boundary value problem posed by the Eqs. (4)–(6) can be reduced to the problem of solving the following integral equation (see [6,7]): (using the assumption that $n(z)$ possesses at most an algebraic growth as $|z| \rightarrow \infty$)

$$\bar{p}(z') = k^2 h_0 G(a, z') + k^2 \int_{-\infty}^{\infty} n(z) G(z, z') \bar{p}(z) dz, \quad (-\infty < z' < \infty) \quad (8)$$

where the function $G(z, z')$ satisfies the following equation and conditions:

$$\left(\frac{\partial^2}{\partial z^2} - \lambda^2\right)G = -\delta(z - z'), \quad (9)$$

$$\frac{\partial G_+}{\partial z} = \frac{\partial G_-}{\partial z}, \quad G_+ = \alpha G_- + \beta \frac{\partial G_-}{\partial z}, \quad \text{on } z = 0, \quad (10)$$

$$G(z, z') \rightarrow 0 \quad \text{as } |z| \rightarrow \infty, \quad (11)$$

with $\beta = \varepsilon \lambda^2 / k^2$. Here $G_+ = G(z, z')$ for $z > 0$, and $G_- = G(z, z')$ for $z < 0$.

The exact d
forward fashio

I. For $z' > 0$:

$$G(z, z') =$$

II. For $z' < 0$:

$$G(z, z')$$

Using these tw

I. For $z' > 0$,

$$\bar{p}(z') = \frac{k^2 h_0}{2\lambda} \left[\varepsilon \right.$$

$$+ \frac{1}{\lambda[\varepsilon \lambda^3]}$$

$$+ \frac{k^2}{2\lambda} \int_0^{\infty}$$

II. For $z' < 0$,

$$\bar{p}(z') = \frac{\alpha}{\lambda[\varepsilon \lambda^3 +}$$

$$+ \frac{k^2}{2\lambda} \int_{-\infty}^0$$

$$+ \frac{1}{\lambda[\varepsilon \lambda^3]}$$

We find tha
growth at infi
possesses a uni
lishes the existe
We write the s

which is the eq

The exact determination of the function $G(z, z')$ can be carried out in a straightforward fashion and we find that

I. For $z' > 0$:

$$G(z, z') = \begin{cases} \frac{1}{2\lambda} \left[\exp(-\lambda|z - z'|) + \frac{(\alpha + \beta\lambda - 1)}{(\alpha + \beta\lambda + 1)} \exp[-\lambda(z + z')] \right], & z > 0 \\ \frac{1}{2\lambda} \left[\frac{2}{(1 + \alpha + \beta\lambda)} \exp[\lambda(z - z')] \right], & z < 0 \end{cases} \quad (12)$$

II. For $z' < 0$:

$$G(z, z') = \begin{cases} \frac{1}{2\lambda} \left[\frac{2\alpha}{(1 + \alpha + \beta\lambda)} \exp[\lambda(z' - z)] \right], & z > 0 \\ \frac{1}{2\lambda} \left[\exp(-\lambda|z - z'|) + \frac{(1 - \alpha + \beta\lambda)}{(1 + \alpha + \beta\lambda)} \exp[\lambda(z + z')] \right], & z < 0. \end{cases} \quad (13)$$

Using these two forms of the function $G(z, z')$ in the above integral Eq. (8), we obtain

I. For $z' > 0$,

$$\begin{aligned} \tilde{p}(z') = & \frac{k^2 h_0}{2\lambda} \left[\exp(-\lambda|a - z'|) + \left\{ \frac{\epsilon\lambda^3 + (\alpha - 1)k^2}{\epsilon\lambda^3 + (\alpha + 1)k^2} \right\} \exp[-\lambda(a + z')] \right] \\ & + \frac{k^4}{\lambda[\epsilon\lambda^3 + (\alpha + 1)k^2]} \int_{-\infty}^0 n(z) \tilde{p}(z) \exp[\lambda(z - z')] dz \\ & + \frac{k^2}{2\lambda} \int_0^{\infty} n(z) \tilde{p}(z) \left\{ \exp(-\lambda|z - z'|) + \left(\frac{\epsilon\lambda^3 + (\alpha - 1)k^2}{\epsilon\lambda^3 + (\alpha + 1)k^2} \right) \exp[-\lambda(z + z')] \right\} dz, \end{aligned} \quad (14)$$

II. For $z' < 0$,

$$\begin{aligned} \tilde{p}(z') = & \frac{\alpha k^4 h_0}{\lambda[\epsilon\lambda^3 + (\alpha + 1)k^2]} \exp[\lambda(z' - a)] \\ & + \frac{k^2}{2\lambda} \int_{-\infty}^0 n(z) \tilde{p}(z) \left\{ \exp(-\lambda|z - z'|) + \left(\frac{\epsilon\lambda^3 + (1 - \alpha)k^2}{\epsilon\lambda^3 + (1 + \alpha)k^2} \right) \exp[\lambda(z + z')] \right\} dz \\ & + \frac{\alpha k^4}{\lambda[\epsilon\lambda^3 + (\alpha + 1)k^2]} \int_0^{\infty} n(z) \tilde{p}(z) \exp[\lambda(z' - z)] dz. \end{aligned} \quad (15)$$

We find that for the class of functions $n(z)$ under consideration (i.e., of algebraic growth at infinity), the integral equation for $\tilde{p}(z)$ as given by the Eqs. (14) and (15) possesses a unique solution for sufficiently small values of the parameter k . This establishes the existence of a unique solution of the direct problem for such small values of k . We write the solution of the Eqs. (14) and (15) in the form

$$\tilde{p}(z) = k^2 \tilde{p}_2(z) + k^4 \tilde{p}_4(z) + k^6 \tilde{p}_6(z) + \dots, \quad (16)$$

which is the equivalent to the Neumann series solution of the above integral equations.

In the special case, when $\varepsilon = 0$, i.e., when the effect of surface tension is negligible, we find, from the relations (14) and (15), that all the terms of the expansion (16), except the first term, which contains the function \tilde{p}_2 , are zero. In the case when $\varepsilon \neq 0$, we proceed to determine these functions in a manner as described below.

We multiply both sides of Eqs. (14) and (15) by the quantity $\lambda[\varepsilon\lambda^3 + (\alpha + 1)k^2]$ and substitute the relation (16). Then equating the terms involving k^2, k^4 and k^6 , we arrive at the following results.

(A) For $z' > 0$,

$$\lambda\tilde{p}_2(z') = (h_0/2)[\exp(-\lambda|a - z'|) + \exp[-\lambda(a + z')]], \tag{17a}$$

$$\begin{aligned} \varepsilon\lambda^4\tilde{p}_4(z') + \lambda(1 + \alpha)\tilde{p}_2(z') &= (h_0/2)[(\alpha + 1)\exp(-\lambda|a - z'|) + (\alpha - 1)\exp[-\lambda(a + z')]] \\ &+ (\varepsilon\lambda^3/2) \int_0^\infty n(z)\tilde{p}_2(z)\{\exp(-\lambda|z - z'|) + \exp[-\lambda(z + z')]\}dz. \end{aligned} \tag{17b}$$

(B) For $z' < 0$,

$$\tilde{p}_2(z') = 0, \tag{18a}$$

$$\varepsilon\lambda^4\tilde{p}_4(z') = \alpha h_0 \exp[\lambda(z' - a)], \tag{18b}$$

$$\begin{aligned} \varepsilon\lambda^4\tilde{p}_6(z') + \lambda(\alpha + 1)\tilde{p}_4(z') &= (\varepsilon\lambda^3/2) \int_{-\infty}^0 n(z)\tilde{p}_4(z)\exp\{-\lambda|z - z'|\} + \exp[\lambda(z + z')]\} dz \\ &+ \alpha \int_0^\infty n(z)\tilde{p}_2(z)\exp[\lambda(z' - z)] dz. \end{aligned} \tag{18c}$$

If we now take limits $z' \rightarrow 0^+$ and $z' \rightarrow 0^-$ of Eqs. (17b) and (18c) respectively, we obtain

$$\int_0^\infty n(z)\exp(-\lambda z)\{\exp(-\lambda|a - z|) + \exp[-\lambda(a + z)]\}dz = \phi(\lambda), \tag{19}$$

$$\int_{-\infty}^0 n(z)\exp(2\lambda z)dz = \psi(\lambda), \tag{20}$$

where

$$\phi(\lambda) = \frac{2}{\varepsilon\lambda^2 h_0} [\varepsilon\lambda^4\tilde{p}_4(0^+) + h_0 \exp(-\lambda a)], \tag{21}$$

$$\psi(\lambda) = \frac{\lambda \exp(\lambda a)}{\alpha h_0} \left[\varepsilon\lambda^4\tilde{p}_6(0^-) + \frac{\alpha h_0}{\lambda} \left\{ \frac{(\alpha + 1)\exp(-\lambda a)}{\varepsilon\lambda^2} - \frac{1}{2}\phi(\lambda) \right\} \right]. \tag{22}$$

It is easy to see that the boundary conditions are satisfied by the functions substituted in (16).

While the results (17a) and (17b) are obtained by the refractive index method, the results (18a) and (18b) are also utilized to determine the values $\tilde{p}_4(0^+)$ and $\tilde{p}_6(0^-)$ as shown below in the next section. We then solve the problem by the usual Laplace transform procedure.

In the results (17a) and (17b),

where

The results (17a) and (17b) for $z > 0$ and $z < 0$ respectively, are these Laplace transforms. It is clear that the functions \tilde{p}_4 and \tilde{p}_6 are constant along lines parallel to the z -axis. Hence the Laplace transform of the proposed problem is obtained by Ramm (see Ramm, 1964).

It is important to note that the relations (25) and (26) are satisfied by $\tilde{p}(0^\pm)$, which are the values of \tilde{p} at $z = 0$. In fact, by using the relations (25) and (26),

$$\tilde{p}_4(0^+)$$

$$\tilde{p}_6(0^-)$$

so that the actual values of \tilde{p}_4 and \tilde{p}_6 are obtained using the known values of $\phi(\lambda)$ and $\psi(\lambda)$.

It is easy to check from the above results that the interface conditions (5) and (6) are satisfied by the various components of the \tilde{p} -functions when series (16) for \tilde{p} is substituted in these interface conditions.

While the results (19) and (20) can be viewed as two identities which must be satisfied by the refraction coefficients $n(z)$ of two fluid media under consideration, we can also utilize these results for the determination of the function $n(z)$ if the limiting values $\tilde{p}_4(0^+)$ and $\tilde{p}_6(0^-)$ are known (or equivalently, if $\tilde{p}(0^+)$ and $\tilde{p}(0^-)$ are known as shown below in Eq. (27)) with the assumption that $a \rightarrow 0^+$. Under such circumstances, we then solve the problem of determination of the function $n(z)$ by a Laplace inversion procedure.

In the results (19)–(22), if we take limit as $a \rightarrow 0^+$, we obtain

$$\int_0^\infty n(z) \exp(-2\lambda z) dz = \frac{\phi_0(\lambda)}{2}, \tag{23}$$

$$\int_{-\infty}^0 n(z) \exp(2\lambda z) dz = \psi_0(\lambda), \tag{24}$$

where

$$\phi_0(\lambda) = \frac{2}{\varepsilon \lambda^2 h_0} [\varepsilon \lambda^4 \tilde{p}_4(0^+) + h_0], \tag{25}$$

$$\psi_0(\lambda) = \frac{\lambda}{\alpha h_0} \left[\varepsilon \lambda^4 \tilde{p}_6(0^-) + \frac{\alpha h_0}{\lambda} \left\{ \frac{(\alpha + 1)}{\varepsilon \lambda^2} - \frac{1}{2} \phi_0(\lambda) \right\} \right]. \tag{26}$$

The results (23) and (24) can be viewed as Laplace transforms of the function $n(z)$ for $z > 0$ and $z < 0$ respectively. Hence the function $n(z)$ can be recovered by inverting these Laplace transforms (see [1]). As is well-known, these Laplace inversions require that the functions $\phi_0(\lambda)$ and $\psi_0(\lambda)$ are known for values of λ lying on certain straight lines parallel to the imaginary axis (or equivalent to it), in the complex λ plane. Hence the Laplace inversion problems given by the relations (25) and (26) are ill-posed problems and these can be resolved numerically by using the ideas proposed by Ramm (see [9]).

It is important to recognize that all the functions appearing in the right of the relations (25) and (26) can be determined, from the knowledge of the given data $\tilde{p}(0^\pm)$, which are the Fourier transforms of the limiting values of the pressure functions. In fact, by using the first results of the Eqs. (17a) and (18a) (when $a \rightarrow 0^+$), we obtain:

$$\begin{aligned} \tilde{p}_4(0^+) &= \lim_{k \rightarrow 0} \{ [\tilde{p}(0^+) - k^2 \tilde{p}_2(0^+)] k^{-4} \} = \lim_{k \rightarrow 0} \{ [\tilde{p}(0^+) - k^2 \lambda^{-1} h_0] k^{-4} \}, \\ \tilde{p}_6(0^-) &= \lim_{k \rightarrow 0} \{ [\tilde{p}(0^-) - k^4 \tilde{p}_4(0^-)] k^{-6} \} = \lim_{k \rightarrow 0} \left\{ [\tilde{p}(0^-) - \frac{\alpha}{\varepsilon} k^4 \lambda^{-4} h_0] k^{-6} \right\}, \end{aligned} \tag{27}$$

so that the actual evaluation of the functions $\phi_0(\lambda)$ and $\psi_0(\lambda)$ can be carried out directly using the known data $\tilde{p}(0^\pm)$.

The above procedure of recovering the unknown function $n(z)$ can be successfully completed, if the other two unknowns of the problem, i.e., the quantities α and h_0 are determined uniquely. To achieve this goal, we use Eqs. (17a) and (18b) and obtain the following formulae (when $a \rightarrow 0^+$).

$$h_0 = \lambda \tilde{p}_2(0^+), \quad \text{and} \quad \frac{\alpha h_0}{\varepsilon} = \lambda^4 \tilde{p}_4(0^-). \tag{28}$$

Thus both the unknowns h_0 and α can be uniquely determined for a given data $\tilde{p}(0^\pm)$ and the relations that

$$\tilde{p}_2(0^+) = \lim_{k \rightarrow 0} \{ \tilde{p}(0^+) k^{-2} \}, \quad \text{and} \quad \tilde{p}_4(0^-) = \lim_{k \rightarrow 0} \{ \tilde{p}(0^-) k^{-4} \}. \tag{29}$$

The inverse problem under consideration thus gets settled and its unique solution can be determined once the Laplace inversions of the two relations (23) and (24) are completed.

The various formulae derived above can be used to construct the following algorithm to recover the values of the three unknowns $\{n(z), \alpha, h_0\}$ from input data $\tilde{p}(0^\pm)$, provided such input data is either known a priori as a solution of the direct problem or obtained by solving a direct problem in a different manner for known values of the triple $\{n(z), \alpha, h_0\}$.

Inverse Algorithm-I: To recover $\{n(z), \alpha, h_0\}$ from $\tilde{p}(0^\pm)$
 [Note: Input data: $\tilde{p}(0^\pm)$ when $a \rightarrow 0^+$.]

- Step 1 Find $\tilde{p}_2(0^+)$ and $\tilde{p}_4(0^-)$ from known input data $\tilde{p}(0^\pm)$ using (29).
- Step 2 Find h_0 and α using (28).
- Step 3 Find $\tilde{p}_4(0^+)$ and $\tilde{p}_6(0^-)$ using (27).
- Step 4 Find $\Phi_0(\lambda)$ and $\Psi_0(\lambda)$ using (25) and (26) respectively.
- Step 5 Finally, find $n(z)$, the refraction coefficient of the medium, by taking inverse Laplace transforms of $\Phi_0(\lambda)$ and $\Psi_0(\lambda)$ according to the relations (23) and (24).

We finally make the following important observations on the effectiveness of the above algorithm.

In Steps 1 and 3 of the algorithm, it is required to calculate certain limits (Eqs. (29) and (27)) using the input data $p(x, y, 0^\mp)$. The problem of calculating these limits is equivalent to finding derivatives of the input data $p(x, y, 0^\mp)$ with respect to wavenumber k . Hence these derivatives must be well defined which in essence requires that the input data and its Fourier transform must be analytic functions of the wavenumber k .

In Step 5 of the above algorithm, the Laplace inversions corresponding to integral Eqs. (23) and (24) are to be performed numerically (which we have not taken up in this paper) by solving two moment problems (see Martin and Ramm [9]), which are the following integral equations of the first kind

$$\int_0^1 x^m n_1(x) dx = \frac{1}{2} \phi_0 \left(\frac{m+1}{2} \right), \tag{30}$$

and

$$\int_0^1 x^m n_2(x) dx = \psi_0 \left(\frac{m+1}{2} \right), \tag{31}$$

for $m = 0, 1, 2,$
 for $z > 0$ in E
 above algorithm
 equations.

4 APPLICAT

In order to illu
 procedure desc
 as described be
 considered to b
 are determined

4.1 Solutions

Example 1 He
 function.

where n_0, n_1 are
 We find that
 satisfying the p
 given by the rel

where

$$B_1 = \frac{1}{2} k^2 h_0 \times \left[\frac{\alpha k^2}{\alpha k^2} \right]$$

$$C_1 = \left[\frac{\alpha k^2}{\alpha k^2} \right]$$

with $H(x)$ deno
 $p(x, y, z)$ as defi

for $m = 0, 1, 2, 3, \dots$. These are obtained by setting $x = \exp(-z)$, $n_1(x) = n(-\ln x)$ for $z > 0$ in Eq. (23), and $x = \exp(z)$, $n_2(x) = n(\ln x)$ for $z < 0$ in Eq. (24). The above algorithm is based on the existence of unique solution of the above two integral equations.

4 APPLICATIONS

In order to illustrate the proposed method of solution of the inverse problem, by the procedure described above, we consider two special examples of the direct problem as described below. The known data for the corresponding inverse problems can be considered to be the ones obtained from the solutions of the direct problems which are determined by the standard techniques.

4.1 Solutions to the Direct Problems

Example 1 Here the refractive index is chosen to be the following piecewise constant function.

$$n(z) = \begin{cases} n_0, & z < 0 \\ n_1, & z > 0 \end{cases} \quad (32)$$

where n_0, n_1 are constants: α and h_0 are also assumed to be known constants.

We find that the solution of this direct problem can be written down immediately, satisfying the partial differential Eq. (4) and the boundary conditions (5) and (6) as given by the relations:

$$\begin{aligned} \tilde{p}(z) &= B_1 \exp[-(\lambda^2 - k^2 n_1)^{1/2} z] + \frac{1}{2} k^2 h_0 (\lambda^2 - k^2 n_1)^{-1/2} \\ &\quad \times [\exp\{(\lambda^2 - k^2 n_1)^{1/2} (z - a)\} H(a - z) \\ &\quad + \exp\{-(\lambda^2 - k^2 n_1)^{1/2} (z - a)\} H(z - a)], \quad \text{for } z > 0, \\ \tilde{p}(z) &= C_1 \exp[(\lambda^2 - k^2 n_0)^{1/2} z], \quad \text{for } z < 0, \end{aligned} \quad (33)$$

where

$$\begin{aligned} B_1 &= \frac{1}{2} k^2 h_0 \exp[-(\lambda^2 - k^2 n_1)^{1/2} a] (\lambda^2 - k^2 n_1)^{-1/2} \\ &\quad \times \left[\frac{\alpha k^2 (\lambda^2 - k^2 n_1)^{1/2} - k^2 (\lambda^2 - k^2 n_0)^{1/2} + \epsilon \lambda^2 (\lambda^2 - k^2 n_0)^{1/2} (\lambda^2 - k^2 n_1)^{1/2}}{\alpha k^2 (\lambda^2 - k^2 n_1)^{1/2} + k^2 (\lambda^2 - k^2 n_0)^{1/2} + \epsilon \lambda^2 (\lambda^2 - k^2 n_0)^{1/2} (\lambda^2 - k^2 n_1)^{1/2}} \right], \\ C_1 &= \left[\frac{\alpha h_0 k^4 \exp[-(\lambda^2 - k^2 n_1)^{1/2} a]}{\alpha k^2 (\lambda^2 - k^2 n_1)^{1/2} + k^2 (\lambda^2 - k^2 n_0)^{1/2} + \epsilon \lambda^2 (\lambda^2 - k^2 n_0)^{1/2} (\lambda^2 - k^2 n_1)^{1/2}} \right], \end{aligned} \quad (30)$$

with $H(x)$ denoting the Heaviside function. Here $\tilde{p}(z)$ is the double Fourier transform of $p(x, y, z)$ as defined by the relation (7) [assuming that $(\lambda^2 - k^2 n_0)^{1/2}$ and $(\lambda^2 - k^2 n_1)^{1/2}$

(30)

(31)

both possess a positive real part, (this choice is required by the condition of vanishing of the transformed pressure function $\tilde{p}(z)$ as $|z| \rightarrow \infty$ (see, for instance, [7])).

Example 2 Here the refractive index is chosen to be

$$n(z) = \begin{cases} n_0 & z > 0 \\ \exp(2z) & z < 0 \end{cases} \quad (34)$$

where n_0 is a constant, α and h_0 are also assumed to be known constants as in the first example.

The differential equation for the transformed pressure function $\tilde{p}(z)$ for the region $z < 0$ can be easily converted to a Bessel's differential equation by a standard transformation (see [8]). Using this idea, we find that the solution of the direct problem, which satisfies the partial differential Eq. (4) and the boundary conditions (5) and (6), can be expressed in the following form:

$$\begin{aligned} \tilde{p}(z) = & B_2 \exp[-(\lambda^2 - k^2 n_0)^{1/2} z] + \frac{1}{2} k^2 h_0 (\lambda^2 - k^2 n_0)^{-1/2} \\ & \times [\exp\{(\lambda^2 - k^2 n_0)^{1/2} (z - a)\} H(a - z) \\ & + \exp\{-(\lambda^2 - k^2 n_0)^{1/2} (z - a)\} H(z - a)], \quad \text{for } z > 0 \end{aligned} \quad (35)$$

$$\tilde{p}(z) = C_2 \sum_{n=0}^{\infty} \frac{(-1)^n (k/2)^{2n+\lambda} \exp[(2n + \lambda)z]}{n! \Gamma(n + \lambda + 1)}, \quad \text{for } z < 0 \quad (36)$$

where

$$\begin{aligned} B_2 = & \frac{1}{2} k^2 h_0 \exp[-(\lambda^2 - k^2 n_0)^{1/2} a] (\lambda^2 - k^2 n_0)^{-1/2} \\ & \times \left[\frac{\alpha k^2 (\lambda^2 - k^2 n_0)^{1/2} - [k^2 - \varepsilon \lambda^2 (\lambda^2 - k^2 n_0)^{1/2}] J'_\lambda(k) / J_\lambda(k)}{\alpha k^2 (\lambda^2 - k^2 n_0)^{1/2} + [k^2 + \varepsilon \lambda^2 (\lambda^2 - k^2 n_0)^{1/2}] J'_\lambda(k) / J_\lambda(k)} \right], \\ C_2 = & \left[\frac{\alpha h_0 k^4 \exp[-(\lambda^2 - k^2 n_0)^{1/2} a]}{J_\lambda(k) \{ \alpha k^2 (\lambda^2 - k^2 n_0)^{1/2} + [k^2 + \varepsilon \lambda^2 (\lambda^2 - k^2 n_0)^{1/2}] J'_\lambda(k) / J_\lambda(k) \}} \right], \end{aligned}$$

$J_\lambda(k)$ denoting the Bessel function of the first kind, and prime denoting derivative with respect to the argument.

4.2 Solutions to the Inverse Problems

The quantities $\tilde{p}(0^\pm)$ for the above two examples are determined from (33) and (35) respectively. The above Algorithm-I is then applied to recover the triple unknowns $\{n(z), \alpha, h_0\}$ for each of these two examples from corresponding known pressure data $p(x, y, 0^\pm)$ or equivalently Fourier transformed data $\tilde{p}(0^\pm)$. We find that the recovered function $n(z)$ from the application of the above inverse algorithm agrees completely with the choices (32) and (34) for the two examples respectively. The other two

unknown constants are determined at various steps of the solution.

I. For example

II. For example

It is easily verified that the solution is recovered from the above for both examples obtained above (see (32) and (34)) of solution of the direct problem.

Acknowledgements

It is a pleasure to thank the

unknown constants α and h_0 are also recovered completely. Some results obtained at various steps of the Algorithm-I when $a \rightarrow 0^+$ are given below.

I. For example 1:

$$\begin{aligned} \bar{p}_2(0^+) &= \frac{h_0}{\lambda}, \\ \bar{p}_4(0^+) &= \frac{h_0}{\lambda^3} \left[\frac{n_1}{2} - \frac{1}{\epsilon\lambda} \right], \\ \bar{p}_4(0^-) &= \frac{\alpha h_0}{\epsilon\lambda^4}, \\ \bar{p}_6(0^-) &= \frac{\alpha h_0}{\epsilon\lambda^6} \left[-\frac{(\alpha + 1)}{\epsilon\lambda} + \frac{(n_0 + n_1)}{2} \right], \\ \phi_0(\lambda) &= \frac{n_1}{\lambda}, \\ \psi_0(\lambda) &= \frac{n_0}{2\lambda}. \end{aligned}$$

II. For example 2:

$$\begin{aligned} \bar{p}_2(0^+) &= \frac{h_0}{\lambda}, \\ \bar{p}_4(0^+) &= \frac{h_0(\epsilon n_0 \lambda - 2)}{2\epsilon\lambda^4}, \\ \bar{p}_4(0^-) &= \frac{\alpha h_0}{\epsilon\lambda^4}, \\ \bar{p}_6(0^-) &= \frac{\alpha h_0}{\epsilon\lambda^4} \left[-\frac{(\alpha + 1)}{\epsilon\lambda^3} + \frac{n_0}{2\lambda^2} + \frac{1}{2\lambda(\lambda + 1)} \right], \\ \phi_0(\lambda) &= \frac{n_0}{\lambda}, \\ \psi_0(\lambda) &= \frac{1}{2(\lambda + 1)}. \end{aligned}$$

It is easily verified that original $n(z)$ for each of the two examples (see (32) and (34)) is recovered from taking the inverse Laplace transforms of $\phi_0(\lambda)$ and $\psi_0(\lambda)$ obtained above for both the examples. Alternatively, it can be seen that the $\phi_0(\lambda)$ and $\psi_0(\lambda)$ obtained above are consistent with (23) and (24) for each of the choices for $n(z)$ (see (32) and (34)) in the two examples. We have thus verified our proposed method of solution of the inverse problem through two examples.

Acknowledgements

It is a pleasure to thank the referees for their suggestions.

References

- [1] A.G. Ramm (1987). Inverse problems. *An Inverse Problem for the Helmholtz Equation in a Semi-infinite Medium*, Vol. 3, L19.
- [2] S.G. Mikhlin (1957). *Integral Equations*, Pergamon, London.
- [3] D.S. Jones (1977). The scattering of sound by a simple shear layer. *Phil. Trans. Roy. Soc., London* **284**, 287-328.
- [4] P.G. Drazin and W.H. Reid (1982). *Hydrodynamic Stability*, Cambridge, New York.
- [5] Horace Lamb (1932). *Hydrodynamics*, Dover, New York.
- [6] D.V. Evans (1968). The Effect of Surface Tension on the Waves Produced by a Heaving Circular Cylinder. *Proc. Camb. Phil. Soc.*, **64**, 833-847.
- [7] W. Ewing, W. Jardetzky and F. Press (1957). *Elastic Waves in Layered Media*. McGraw-Hill, New York.
- [8] G.M. Murphy (1960). *Ordinary Differential Equations and Their Solutions*. D. Van Nostrand, Inc., London.
- [9] P.A. Martin and A.G. Ramm (1985). Inverse scattering for geophysical problems. III. On the velocity-inversion problems of acoustics. *Proc. Roy. Soc., London, A* **399**, 153.

This work is concerned with near-field measurements of the inverse obstacle problem. The problem is ill-posed and a well-posed minimization process and nonlinear regularization are used to make the problem very easy and fast to solve. The unknown function is determined by an attractive feature of the problem: the field is known in the unbounded region. The computed fields due to the problem are limited by the aperture of the observation aperture.

Keywords: Inverse problem

1. INTRODUCTION

In inverse acoustics, where the boundary conditions are known and the field is measured, one can determine the shape of the obstacle. In medical testing, medical imaging, the inverse problem has been studied for two decades. So far, the problem has been solved for a limited number of cases.

We note that the problem is ill-posed. We note that the pattern of the scattered field is

*Corresponding author