

1. (45) Consider the following problem

$$\begin{aligned}\frac{\partial^2 u}{\partial t^2} &= \frac{\partial^2 u}{\partial x^2} & 0 < x < \pi, & \quad 0 < t \\ u(0, t) &= 0, \quad u_x(\pi, t) = 0, & & \quad 0 < t \\ u(x, 0) &= \sin x, \quad u_t(x, 0) = e^x, & & \quad 0 < x < \pi\end{aligned}$$

Let ASTP refer to the Associated Sturm-Liouville Problem that arises from this problem when the separation of variables technique is applied.

- (a) What is the ASTP ?

$$\begin{aligned}\frac{d^2 \phi}{dx^2} + \lambda \phi &= 0 \\ \phi(0) &= 0 \\ \phi'(\pi) &= 0\end{aligned}$$

- (b) Derive the Rayleigh quotient expression for the eigenvalues of the ASTP. What can you say about the eigenvalues of ASTP?

Multiply the differential equation that λ and ϕ satisfy by ϕ and then integrate by parts

$$\begin{aligned}0 &= \lambda \int_0^\pi \phi^2 dx + \int_0^\pi \phi'' \phi dx \\ &= \lambda \int_0^\pi \phi^2 dx + \left\{ \phi' \phi \Big|_0^\pi - \int_0^\pi \phi' \phi' dx \right\} \\ &= \lambda \int_0^\pi \phi^2 dx - \int_0^\pi (\phi')^2 dx\end{aligned}$$

The boundary conditions, which ϕ satisfies enable us to conclude that $\phi' \phi|_0^\pi = 0$. Solving for λ we have

$$\lambda = \frac{\int_0^\pi (\phi')^2}{\int_0^\pi \phi^2 dx}$$

This equality enables us to conclude that any eigenvalue must satisfy $\lambda \geq 0$. Moreover if 0 is an eigenvalue, then the corresponding eigenfunction must equal $\phi = ax + b$ for some constants a and b . The boundary conditions then force a and b to both equal 0. Thus, there are no non-trivial solutions to the Sturm-Liouville problem when $\lambda = 0$, and we may conclude that 0 is not an eigenvalue. Thus, every eigenvalue must be positive.

(c) What are the eigenvalues and eigenfunctions of the ASTP?

Since the eigenvalues are positive we know that any solution of $\frac{d^2\phi}{dx^2} + \lambda\phi = 0$ is of the form $\phi = a \cos \sqrt{\lambda}x + b \sin \sqrt{\lambda}x$. The boundary condition $\phi(0) = 0$ forces $a = 0$. Thus, $\phi = \sin \sqrt{\lambda}x$, and $\phi'(\pi) = 0$ only if

$$\begin{aligned}\sqrt{\lambda}\pi &= (2k-1)\frac{\pi}{2} \text{ for } k = 1, 2, \dots \\ \sqrt{\lambda} &= \frac{2k-1}{2} \text{ and } \lambda = \left(\frac{2k-1}{2}\right)^2\end{aligned}$$

Thus, the eigenvalues are $\lambda = \left(\frac{2k-1}{2}\right)^2$ and the corresponding eigenfunctions are $\sin \frac{2k-1}{2}x$.

(d) How would you solve the original problem?

When separating variables with $u(x, t) = g(x)h(t)$ the function h is found to satisfy the differential equation $h'' + \lambda h = 0$. Thus,

$$h(t) = a_k \cos\left(\frac{2k-1}{2}t\right) + b_k \sin\left(\frac{2k-1}{2}t\right)$$

We look for a solution of the original problem in the form

$$u(x, t) = \sum_{k=1}^{\infty} \left(a_k \cos\left(\frac{2k-1}{2}t\right) + b_k \sin\left(\frac{2k-1}{2}t\right) \right) \sin\left((2k-1)\frac{\pi}{2L}x\right)$$

The constants a_k and b_k are determined from the initial conditions.

$$\begin{aligned}\sin x &= u(x, 0) = \sum_{k=1}^{\infty} a_k \sin\left(\frac{2k-1}{2}x\right) \\ e^x &= u_t(x, 0) = \sum_{k=1}^{\infty} \frac{2k-1}{2} b_k \sin\left(\frac{2k-1}{2}x\right)\end{aligned}$$

Thus, we must have

$$\begin{aligned}a_k &= \frac{2}{\pi} \int_0^{\pi} \sin x \sin\left(\frac{2k-1}{2}x\right) dx \\ b_k &= \frac{4}{(2k-1)\pi} \int_0^{\pi} e^x \sin\left((2k-1)\frac{\pi}{2L}x\right) dx\end{aligned}$$

2. (45) Solve the following problem

$$\begin{aligned}\frac{\partial u}{\partial t} &= \frac{\partial^2 u}{\partial x^2} + \frac{x}{\pi} (\cos t + \sin t) \quad 0 < x < \pi, \quad 0 < t \\ u(0, t) &= \cos t, \quad u(\pi, t) = 1 + \sin t \quad 0 < t \\ u(x, 0) &= 1, \quad 0 < x < \pi\end{aligned}$$

The first step is to reduce this problem to one with homogeneous boundary conditions. Set $v(x, t) = u(x, t) - r(x, t)$, where $r(x, t) = \cos t + \frac{x}{\pi}(1 + \sin t - \cos t)$. Then v satisfies the following

$$\begin{aligned}\frac{\partial v}{\partial t} &= \frac{\partial^2 v}{\partial x^2} + \frac{x}{\pi} (\cos t + \sin t) - \left(-\sin t + \frac{x}{\pi} (\cos t + \sin t) \right) \text{ or} \\ \frac{\partial v}{\partial t} &= \frac{\partial^2 v}{\partial x^2} + \sin t, \quad 0 < x < \pi, \quad 0 < t \\ v(0, t) &= 0, \quad v(\pi, t) = 0, \quad 0 < t \\ v(x, 0) &= 0, \quad 0 < x < \pi\end{aligned}$$

The eigenvalues and eigenfunctions for the Sturm-Liouville problem, which arises from solving the associated homogeneous partial differential equation and boundary conditions are $\lambda = n^2$, $\phi_n(x) = \sin nx$, for $n = 1, 2, \dots$. So we look for the solution v in the form

$$v(x, t) = \sum_{n=1}^{\infty} a_n(t) \sin nx$$

Moreover, since the initial condition $v(x, 0) = 0$, we must have $a_n(0) = 0$ for each n . Placing this expression into the differential equation for v we have

$$\begin{aligned}\sum_{n=1}^{\infty} a'_n(t) \sin nx &= \sum_{n=1}^{\infty} (-n^2) a_n(t) \sin nx + \sin t \\ \sum_{n=1}^{\infty} (a'_n(t) + n^2 a_n(t)) \sin nx &= \sin t = \sum_{n=1}^{\infty} b_n(t) \sin nx \quad \text{or} \\ a'_n(t) + n^2 a_n(t) &= b_n(t), \quad a_n(0) = 0, \quad \text{with} \\ b_n(t) &= \frac{2 \sin t}{\pi} \int_0^{\pi} \sin nx \, dx, \quad \text{for } n = 1, 2, \dots \\ &= \frac{2 \sin t}{\pi} \left(\frac{1 - \cos(n\pi)}{n} \right) \\ &= \begin{cases} 0 & n \text{ even} \\ \frac{4 \sin t}{n\pi} & n \text{ odd} \end{cases}\end{aligned}$$

The solution of this initial value problem is

$$a_n(t) = \frac{4e^{-n^2 t}}{n^5 \pi + n\pi} + \frac{4(n^2 \sin t - \cos t)}{n\pi (n^4 + 1)}$$

when n is odd. Thus, we finally have

$$\begin{aligned}u(x, t) &= \sum_{n=1}^{\infty} \left(\frac{4e^{-(2n-1)^2 t}}{(2n-1)^5 \pi + (2n-1)\pi} + \frac{4((2n-1)^2 \sin t - \cos t)}{(2n-1)\pi ((2n-1)^4 + 1)} \right) \sin(2n-1)x \\ &\quad + \cos t + \frac{x}{\pi} (1 + \sin t - \cos t)\end{aligned}$$

3. (10) Consider the Sturm-Liouville problem

$$\begin{aligned} \frac{d}{dx} \left(p(x) \frac{d\phi}{dx} \right) + q(x) \phi + \lambda \sigma(x) \phi &= 0, \quad 0 < x < L \\ \phi(0) &= 0 \\ \phi'(L) &= 0 \end{aligned}$$

where p , q , and σ are continuous on the interval $[0, L]$, and both p and σ are positive and bounded away from zero. Let $L(\phi) = \frac{d}{dx} \left(p(x) \frac{d\phi}{dx} \right) + q(x) \phi$

The following differential equality may be useful

$$L(u)v - uL(v) = \frac{d}{dx} (pu'v - puv') .$$

- (a) Use the differential equality to show that $\int_0^L L(u)v \, dx = \int_0^L uL(v) \, dx$ for every pair of functions u and v which satisfy the boundary conditions of the Sturm-Liouville problem.

$$\begin{aligned} \int_0^L (L(u)v - uL(v)) \, dx &= (pu'v - puv') \Big|_0^L \\ &= p(L)u'(L)v(L) - p(L)u(L)v'(L) \\ &\quad - (p(0)u'(0)v(0) - p(0)u(0)v'(0)) \\ &= 0 \end{aligned}$$

Thus, we have $\int_0^L L(u)v \, dx = \int_0^L uL(v) \, dx$.

- (b) Show that if λ_1 and λ_2 are two different eigenvalues of the Sturm-Liouville problem with corresponding eigenfunctions ϕ_1 and ϕ_2 , then

$$\int_0^L \phi_1(x) \phi_2(x) \sigma(x) \, dx = 0 .$$

$$\begin{aligned} \lambda_1 \int_0^L \phi_1(x) \phi_2(x) \sigma(x) \, dx &= \int_0^L \lambda_1 \sigma(x) \phi_1(x) \phi_2(x) \, dx \\ &= - \int_0^L L(\phi_1) \phi_2 \, dx \\ &= - \int_0^L \phi_1 L(\phi_2) \, dx \\ &= \int_0^L \phi_1(x) \lambda_2 \phi_2(x) \sigma(x) \, dx \\ &= \lambda_2 \int_0^L \phi_1(x) \phi_2(x) \sigma(x) \, dx \end{aligned}$$

This implies that

$$(\lambda_1 - \lambda_2) \int_0^L \phi_1(x) \phi_2(x) \sigma(x) \, dx = 0$$

Since $\lambda_1 - \lambda_2 \neq 0$, we must have $\int_0^L \phi_1(x) \phi_2(x) \sigma(x) \, dx = 0$.