### A Central Element of the Quantum Group $U_q(\mathfrak{so}_{2n})$

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### Outline of the talk

- Define basic algebraic structures and research problem
- Apply main formula for simple cases n = 3, 4
- ullet Describe additional progress for general n
- Show some probabilistic applications

## The underlying Lie algebra

#### **Definition**

The Lie algebra  $\mathfrak{so}_{2n}(\mathbb{C})$  is the set of  $2n \times 2n$  matrices

$$\left\{ \begin{bmatrix} A & B \\ C & D \end{bmatrix} : A = -D^T, B^T = -B, C^T = -C \right\},\,$$

where  $A, B, C, D \in \mathbb{C}^{n \times n}$ .

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- Often study operators by analyzing eigenvalues and eigenspaces.
- Analogously, there are two types of "eigenvalues" we'll consider
  - Weights (denoted  $\mu$  or  $\lambda$ ) for 2*n*-dim. **fundamental** representation
  - Roots (denoted  $\alpha_i$  or  $-\alpha_i$ ) for  $(2n)^2$ -dim. **adjoint** representation
- Let  $L_i$  be a function which sends a matrix M to the diagonal entry  $M_{ii}$ . The weights and roots for  $\mathfrak{so}_{2n}$  are

$$\mu = \pm L_i, \quad \alpha = \pm L_i \pm L_j$$

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  - Generated by  $E_i, F_i, H_i$   $(1 \le i \le n)$ . Example (n = 2)

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  - Generated by  $E_i, F_i, q^{\pm H_i}$  with q-deformed relations
  - Example of an element

$$(q^2+1)E_1^2F_1q^{H_1-H_2}$$
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$$\textit{E}_{1} = \begin{bmatrix} \vdots & 1 & \ddots & \vdots \\ \vdots & \vdots & \ddots & \vdots \\ \vdots & \ddots & -1 & \vdots \end{bmatrix}, \;\; \textit{F}_{2} = \begin{bmatrix} \vdots & \ddots & \ddots & \vdots \\ \vdots & \vdots & \ddots & \vdots \\ -1 & \ddots & \ddots & \vdots \end{bmatrix}, \;\; \textit{H}_{1} = \begin{bmatrix} 1 & \ddots & \ddots & \vdots \\ \vdots & -1 & \ddots & \vdots \\ \vdots & \ddots & -1 & \vdots \\ \vdots & \ddots & \ddots & 1 \end{bmatrix}$$

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### Motivation: the Casimir element

- In  $U(\mathfrak{so}_{2n})$  (the symmetric case), the **quadratic Casimir element** is a distinguished element of the center.
- This Casimir element can be procedurally represented as a generator matrix of a Markov process.
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## The main problem and formula

#### **Problem**

Find an explicit form for a central element of  $U_q(\mathfrak{so}_{2n})$  in terms of the generators  $E_i, F_i, q^{H_i}$ .

Recall: in  $U(\mathfrak{so}_{2n})$ , we find dual elements and compute  $\sum_i X_i X^i$ .

### Proposition (Kuan '16)

For each weight  $\mu$ , let  $v_{\mu}$  be a vector in its weight space. Given weights  $\mu, \lambda$ , suppose  $e_{\mu\lambda}$  sends  $v_{\lambda}$  to  $v_{\mu}$  and  $f_{\lambda\mu}$  sends  $v_{\mu}$  to  $v_{\lambda}$ . If  $e_{\mu\lambda}^*$  and  $f_{\mu\lambda}^*$  are their q-pairing dual elements, and  $\rho$  is half the sum of the positive roots of  $\mathfrak{g}$ , then

$$\sum_{\mu} q^{(-2\rho,\mu)} q^{H_{-2\mu}} + \sum_{\mu > \lambda} q^{(\mu-\lambda,\mu)} q^{(-2\rho,\mu)} e_{\mu\lambda}^* q^{H_{-\mu-\lambda}} f_{\lambda\mu}^*$$

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#### We wish to compute

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- $(-2\rho, \mu)$  and  $(\mu \lambda, \mu)$  are ordinary dot products, so the corresponding terms are just powers of q.
- $q^H$ s are products of  $q^{\pm H_i}$ s, which are also simple to compute.

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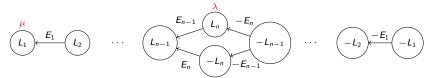
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# Computing $e_{\mu\lambda}$ and $f_{\lambda\mu}$

The generators  $E_i$ s and  $F_i$ s are operators that move us between different weight spaces.



Here,  $e_{\mu\lambda}$  and  $f_{\lambda\mu}$  send us from  $\lambda$  to  $\mu$  and vice versa. In this case,  $e_{\mu\lambda}=E_1\cdots E_{n-1}$ , and  $f_{\lambda\mu}=F_{n-1}\cdots F_1$ .

## A q-deformed pairing

We introduce a function  $\langle , \rangle$ , which takes in (product of Fs and  $q^Hs$ ) and (product of Es and  $q^Hs$ ), outputting (rational function in q). More formally:  $U_q(\mathfrak{b}-)\times U_q(\mathfrak{b}+)\to \mathbb{Q}(q)$ .

For the generators, the only nonzero pairings are

$$\langle q^{H_{\alpha}}, q^{H_{\beta}} \rangle = q^{-(\alpha \cdot \beta)}, \quad \langle F_i, E_i \rangle = -\frac{1}{q - q^{-1}},$$

where  $\alpha$  and  $\beta$  are linear combinations of the  $\alpha_i$ s.

There is also an inductive way to compute things like

$$\langle q^{H_1}F_2F_1, q^{H_2}E_1E_2\rangle,$$

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Take n = 4. Here are some example computations:

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### Example: find **dual element under** $\langle , \rangle$ of $\underline{F_2F_1}$ for n=3.

- $\{F_1F_2, \underline{F_2F_1}\}$  both have nonzero pairing with both of  $\{E_1E_2, E_2E_1\}$ . (Call these  $\{f_1, f_2\}$  and  $\{e_1, e_2\}$ .)
- **Dual elements**  $f_i^*$  are combinations of the  $e_i$ s, such that  $\langle f_i, f_i^* \rangle = \delta_{ii}$ .
- Form matrix of pairings M such that  $M_{ij} = \langle f_i, e_j \rangle$ :

$$M = (q - q^{-1})^2 \begin{bmatrix} 1 & 1/q \\ 1/q & 1 \end{bmatrix}$$

Invert the matrix and look at corresponding (second) row.

$$M^{-1} = (q - q^{-1}) \begin{bmatrix} q & -1 \\ -1 & q \end{bmatrix}$$



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## Computational difficulties

Two main reasons this is more complicated than the other steps:

- Matrix *M* needs to be **invertible**.
  - Need to make sure different f<sub>i</sub>s and e<sub>i</sub>s linearly independent
  - Serre relation makes this hard: for example,

$$E_1^2 E_2 + E_2 E_1^2 = (1+q)E_1 E_2 E_1$$

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# The central element of $U_q(\mathfrak{so}_6)$

Let  $r = q - \frac{1}{q}$ , and write (for example)  $E_1 E_2 E_3$  as  $E_{123}$ .

#### Theorem (L.)

The following element of the quantum group  $U_q(\mathfrak{so}_6)$  is central:

$$\begin{array}{l} q^{-4-2H_1-H_2-H_3} + q^{-2-H_2-H_3} + q^{H_2-H_3} + q^{H_3-H_2} + q^{2+H_2+H_3} + q^{4+2H_1+H_2+H_3} + \frac{r^2}{q^3}F_1q^{-H_1-H_2-H_3}E_1 \\ \\ + \frac{r^2}{q}F_2q^{-H_3}E_2 - \frac{r^2}{q^2}F_3q^{-H_2}E_3 + r^2qF_2q^{H_3}E_2 - r^2qF_3q^{H_2}E_3 + r^2q^3F_1q^{H_1+H_2+H_3}E_1 \\ \\ + \frac{r^2}{q^3}(qF_{12} - F_{21})q^{-H_1-H_3}(qE_{21} - E_{12}) - \frac{r^2}{q^3}(qF_{13} - F_{31})q^{-H_1-H_2}(qE_{31} - E_{13}) \\ \\ + r^2q(qF_{21} - F_{12})q^{H_1+H_3}(qE_{12} - E_{21}) - r^2q(qF_{31} - F_{13})q^{H_1+H_2}(qE_{13} - E_{31}) \\ \\ - \frac{r^2}{q^3}(q^2F_{123} - qF_{213} - qF_{213} - qF_{213})q^{-H_1}(q^2E_{231} - qE_{213} - qE_{213} + E_{123}) \\ \\ - \frac{r^4}{q^2}((q^2+1)F_{1231} - qF_{1312} - qF_{2131})((q^2+1)E_{1231} - qE_{1312} - qE_{2131}) \\ \\ - r^4F_2F_3E_2E_3. \end{array}$$

This element acts as a constant  $(q^6 + q^2 + 2 + q^{-2} + q^{-6})$  times the identity matrix in the fundamental representation.

# The central element of $U_q(\mathfrak{so}_8)$

#### Theorem (L.)

The following element of the quantum group  $U_q(\mathfrak{so}_8)$  is central:

$$q^{-6-2H_1-2H_2-H_3-H_4} + q^{-4-2H_2-H_3-H_4} + q^{-2-H_3-H_4} + q^{H_3-H_4} + q^{H_3-H_4} + q^{H_4-H_3} + q^{2+H_3+H_4} + q^{4+2H_2+H_3+H_4} + q^{6+2H_1+2H_2+H_3+H_4} + q^{H_4-H_3} + q^{2+H_3+H_4} + q^{4+2H_2+H_3+H_4} + q^{4+2H_2+H_3+H_4} + q^{2+2H_3-H_4} +$$

# The central element of $U_q(\mathfrak{so}_8)$ , continued

#### **Theorem**

(Here is the rest of the element.)

$$\cdots - r^4 F_3 F_4 E_4 E_3 - \frac{r^2}{q} (q^2 F_{432} - q F_{324} - q F_{423} + F_{234}) q^{H_2} (q^2 E_{234} - q E_{324} - q E_{423} + E_{432})$$

$$- \frac{r^2}{q} A_3 q^{H_1 + H_2} A_2 - r^2 q F_4 q^{H_3} E_4 - r^2 q (q F_{42} - F_{24}) q^{H_2 + H_3} (q E_{24} - E_{42})$$

$$- r^2 q (q^2 F_{421} - q F_{241} - q F_{142} + F_{124}) q^{H_1 + H_2 + H_3} (q^2 E_{124} - q E_{142} - q E_{241} + E_{421})$$

$$+ r^2 q F_3 q^{H_4} E_3 + r^2 q (q F_{32} - F_{32}) q^{H_2 + H_4} (q E_{23} - E_{32})$$

$$+ r^2 q (q^2 F_{321} - q F_{213} - q F_{132} + F_{123}) q^{H_1 + H_2 + H_4} (q^2 E_{123} - q E_{132} - q E_{213} + E_{321})$$

$$+ r^2 q^3 F_2 q^{H_2 + H_3 + H_4} E_2 + r^2 q^3 (q F_{21} - F_{12}) q^{H_1 + H_2 + H_3 + H_4} (q E_{12} - E_{21}) + r^2 q^5 F_1 q^{H_1 + 2H_2 + H_3 + H_4} E_1,$$

where the 10 boxed  $A_i$  s are omitted for brevity. This element acts as  $q^8 + q^4 + q^2 + 2 + q^{-2} + q^{-4} + q^{-8}$  times the identity matrix in the fundamental representation.

A strategy for computing certain dual elements:

#### Proposition (L.)

Suppose each index only shows up once in an element of  $e_{\mu\lambda}$  or  $f_{\lambda\mu}$ . Then the matrix  $M^{-1}$  can be inductively computed by tensoring the inverse matrix from a smaller set of indices repeatedly with  $\begin{bmatrix} q & -1 \\ -1 & q \end{bmatrix}$ .

- Dual of  $E_1$  is  $(q q^{-1})F_1$
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### Dual elements for general n, continued

- The above strategy doesn't work for repeated indices.
- However, the dimensions of *M* for small *n* show a pattern.
  - The dimension for indices (2, 2, 3, 4) in n = 4 is 5
  - The dimensions for (1, 2, 2, 3, 4) and (1, 1, 2, 2, 3, 4) are 15 and 20.

#### Conjecture

Suppose the index  $x_1-1$  is being added to a set of indices  $S=(x_1,\cdots,x_m)$  of dimension d, where  $x_1\leq\cdots\leq x_m$  and  $x_1\leq n-2$ .

- If  $x_1$  appears twice and we add  $(x_1 1)$  once, the dimension becomes 3d
- If  $x_1$  appears twice and we add  $(x_1 1)$  twice, the dimension becomes 4d.

Suppose M is the a pairing matrix for some basis for S. Then we can find a  $3 \times 3$  matrix  $M_1$  and a  $4 \times 4$  matrix  $M_2$ , such that the new pairing matrix M' is  $M \otimes M_1$  in the first case and  $M \otimes M_2$  in the second case.

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#### From central element to tensor representation

#### In order to extract the probabilistic interpretation:

• Replace each generator with its coproduct. For example

$$E_i \to E_i \otimes I + q^{H_i} \otimes E_i$$
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(This is similar to the symmetric case, where  $E_i \rightarrow E_i \otimes I + I \otimes E_i$ .)

- End up with a  $4n^2 \times 4n^2$  matrix with coefficients in terms of q.
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# Modifying to a generator matrix

The resulting  $4n^2 \times 4n^2$  matrix is not yet a generator matrix, just like in the symmetric case.

- Key idea: if Mv = 0, where  $v = (v_1, \dots, v_N)$ , we can conjugate by a diagonal matrix  $D = \text{diag}(v_1, \dots, v_N)$ .
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### Preliminary observations

Recall these properties of the generator matrix in the symmetric case:

- All nonzero off-diagonal entries equal
- 2n absorbing states, 2n maximal-choice states, all others pairwise.

Similar properties can be observed at least for  $U_q(\mathfrak{so}_6)$  and  $U_q(\mathfrak{so}_8)$ :

- The absorbing and pairwise states interact in the same ways (except the jump rates differ by a factor of  $q^2$ , causing **drift**).
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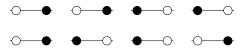
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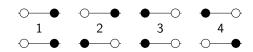
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#### Patterns in the coefficients

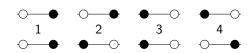


Here is the generator submatrix for n = 3:

$$\frac{1}{q^6}\begin{bmatrix} -1 - 2q^2 + q^6 - q^8 - q^{10} & q^2(2 - q^4 + q^6) & (q^4 - 1)^2 & q^4(2 - q^4 + q^6) \\ q^4(2 - q^4 + q^6) & -1 + 2q^4 + q^8 - 2q^{10} & 1 - q^2 + 2q^6 & q^2(q^4 - 1)^2 \\ q^4(q^4 - 1)^2 & q^2(1 - q^2 + 2q^6) & -q^2 - q^4 + q^6 - 2q^{10} - q^{12} & q^4(1 - q^2 + 2q^6) \\ q^6(2 - q^4 + q^6) & q^2(q^4 - 1)^2 & q^2(1 - q^2 + 2q^6) & -2q^2 + q^4 - 2q^8 - q^{12} \end{bmatrix}$$

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# Thank you!

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#### References



J. Jantzen.

Lectures on Quantum Groups.

DIMAC Series in Discrete Mathematics and Theoretical Computer Science. American Mathematical Society.



J. Kuan.

Stochastic duality of ASEP with two particle types via symmetry of quantum groups of rank two.

J. Phys. A, 49(11):115002, 29, 2016.