

Converting the Black-Scholes PDE
to
The Heat Equation

The Black-Scholes partial differential equation and boundary value problem is

$$L(V) = \frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + rS \frac{\partial V}{\partial S} - rV = 0, \quad 0 \leq S, \quad 0 \leq t \leq T$$

$$V(S, T) = f(S), \quad 0 \leq S, \quad V(0, t) = 0, \quad 0 \leq t \leq T.$$

If V is the price of a call option, then the boundary condition $f(S) = \max(S - E, 0)$, where E denotes the strike price of the call option.

The following change of variables transforms the Black-Scholes boundary value problem into a standard boundary value problem for the heat equation.

$$S = e^x, \quad t = T - \frac{2\tau}{\sigma^2},$$

$$V(S, t) = v(x, \tau) = v\left(\ln(S), \frac{\sigma^2}{2}(T - t)\right).$$

The partial derivatives of V with respect to S and t expressed in terms of partial derivatives of v in terms of x and τ are:

$$\frac{\partial V}{\partial t} = -\frac{\sigma^2}{2} \frac{\partial v}{\partial \tau}$$

$$\frac{\partial V}{\partial S} = \frac{1}{S} \frac{\partial v}{\partial x}$$

$$\frac{\partial^2 V}{\partial S^2} = -\frac{1}{S^2} \frac{\partial v}{\partial x} + \frac{1}{S^2} \frac{\partial^2 v}{\partial x^2}$$

Placing these expressions into the Black-Scholes partial differential equation and simplifying we have

$$\frac{\partial v}{\partial \tau} = \frac{\partial^2 v}{\partial x^2} + \left(\frac{2r}{\sigma^2} - 1\right) \frac{\partial v}{\partial x} - \frac{2r}{\sigma^2} v.$$

Setting $\kappa = 2r/\sigma^2$ and $t = \tau$, the Black-Scholes boundary value problem becomes

$$\frac{\partial v}{\partial t} = \frac{\partial^2 v}{\partial x^2} + (\kappa - 1) \frac{\partial v}{\partial x} - \kappa v, \quad -\infty < x < \infty, \quad 0 \leq t \leq \frac{\sigma^2}{2} T$$

$$v(x, 0) = V(e^x, T) = f(e^x), \quad -\infty < x < \infty$$

One more change of variables is needed in order to eliminate the last two terms on the right hand side of the last equation. To this end set

$$v(x, t) = e^{\alpha x + \beta t} u(x, t) = \phi u,$$

where we'll pick α and β later. Computing the partials of v in terms of x and t we have

$$\begin{aligned} \frac{\partial v}{\partial t} &= \beta \phi u + \phi \frac{\partial u}{\partial t} \\ \frac{\partial v}{\partial x} &= \alpha \phi u + \phi \frac{\partial u}{\partial x} \\ \frac{\partial^2 v}{\partial x^2} &= \alpha^2 \phi u + 2\alpha \phi \frac{\partial u}{\partial x} + \phi \frac{\partial^2 u}{\partial x^2} \end{aligned}$$

Placing these expressions into the partial differential equation which v satisfies, and setting

$$\begin{aligned} \alpha &= -\frac{1}{2}(k-1) = \frac{\sigma^2 - 2r}{2\sigma^2} \\ \beta &= -\frac{1}{4}(k+1)^2 = -\left(\frac{\sigma^2 + 2r}{2\sigma^2}\right)^2. \end{aligned}$$

we have

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, \quad -\infty < x < \infty, \quad 0 \leq t \leq \frac{\sigma^2}{2}T \quad (1)$$

$$u(x, 0) = e^{-\alpha x} v(x, 0) = e^{-\alpha x} f(e^x), \quad -\infty < x < \infty \quad (2)$$

If the option is a call option, with strike price E , then $f(x) = \max(x - E, 0)$, and

$$u(x, 0) = e^{-\alpha x} \max(e^x - E, 0).$$

It can be shown that the solution to the heat equation (1) and initial condition (2) is given by the following integral

$$u(x, t) = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} u(\xi, 0) e^{-\frac{(x-\xi)^2}{4t}} d\xi.$$

Find the value of an option, whose value at expiration equals $f(S)$, where

$$f(S) = \begin{cases} 0, & S < 1 \\ 3, & 1 \leq S \leq 2 \\ 0, & S > 3 \end{cases} .$$

$$\begin{aligned} V(S, 0) &= \nu \left(\ln S, \frac{\sigma^2 T}{2} \right) = e^{\alpha \ln S} e^{\beta \frac{\sigma^2 T}{2}} u \left(\ln S, \frac{\sigma^2 T}{2} \right) \\ &= e^{\alpha \ln S} e^{\beta \frac{\sigma^2 T}{2}} \frac{1}{\sqrt{4\pi \frac{\sigma^2 T}{2}}} \int_{-\infty}^{\infty} u(\xi, 0) e^{-\frac{(\ln S - \xi)^2}{4 \frac{\sigma^2 T}{2}}} d\xi \\ &= e^{\alpha \ln S} e^{\beta \frac{\sigma^2 T}{2}} \frac{1}{\sqrt{2\pi \sigma^2 T}} \int_{-\infty}^{\infty} e^{-\alpha \xi} f(e^\xi) e^{-\frac{(\ln S - \xi)^2}{2\sigma^2 T}} d\xi \\ &= e^{\alpha \ln S} e^{\beta \frac{\sigma^2 T}{2}} \frac{3}{\sqrt{2\pi \sigma^2 T}} \int_0^{\ln 2} e^{-\alpha \xi} e^{-\frac{(\ln S - \xi)^2}{2\sigma^2 T}} d\xi \\ &= e^{\alpha \ln S} e^{\beta \frac{\sigma^2 T}{2}} \frac{3S^{-\alpha}}{\sqrt{2\pi \sigma^2 T}} e^{\frac{\alpha^2 \sigma^2 T}{2}} \int_{\lambda_1}^{\lambda_2} e^{-\lambda^2/2} d\lambda \quad \begin{cases} \lambda_1 = \frac{\ln(S/2) + (r - \sigma^2/2)T}{\sigma\sqrt{T}} \\ \lambda_2 = \frac{\ln S + (r - \sigma^2/2)T}{\sigma\sqrt{T}} \end{cases} \\ &= 3e^{-\frac{\sigma^2 + 8r}{8}T} [N(\lambda_2) - N(\lambda_1)] . \end{aligned}$$