## Exam 1: Solutions

Problem 1 (40 pts.) Let $f(x)=2 x$ for $0 \leq x \leq \pi$.
(i) Find the Fourier sine series of $f$ (with $[0, \pi]$ as the basic interval).

Take $L=\pi$ in the usual formulas:

$$
2 x \sim \sum_{n=1}^{\infty} B_{n} \sin n x
$$

where

$$
\begin{aligned}
B_{n} & =\frac{2}{\pi} \int_{0}^{\pi} 2 x \sin n x d x=\frac{4}{\pi} \int_{0}^{\pi} x \sin n x d x=-\frac{4}{n \pi} \int_{0}^{\pi} x d(\cos n x) \\
& =-\frac{4}{n \pi}\left(\left.x \cos n x\right|_{0} ^{\pi}-\int_{0}^{\pi} \cos n x d x\right)=-\frac{4}{n \pi}\left(\pi \cos n \pi-\left.\frac{\sin n x}{n}\right|_{0} ^{\pi}\right) \\
& =-\frac{4}{n} \cos n \pi=(-1)^{n+1} \frac{4}{n} .
\end{aligned}
$$

(ii) Over the interval $[-2.5 \pi, 2.5 \pi]$, sketch the function to which the series converges.

(iii) Roughly sketch the 12 th partial sum of the series.

(iv) Briefly describe how the answers to (ii) and (iii) would change if we studied the Fourier cosine series instead.

The Fourier cosine series converges to the even $2 \pi$-periodic extension of $f$. This extension is continuous, so there is no Gibbs phenomenon in the partial sums.

Problem 2 ( 30 pts.) Solve the initial-boundary value problem for the heat equation

$$
\begin{aligned}
& \frac{\partial u}{\partial t}=\frac{\partial^{2} u}{\partial x^{2}} \quad(0<x<\pi, \quad t>0) \\
& u(x, 0)=2 x \quad(0<x<\pi) \\
& u(0, t)=u(\pi, t)=0 \quad(t>0) .
\end{aligned}
$$

You may stop when you can say "And now continue as in Problem 1, above."
First look for solutions with separated variables: $u(x, t)=\phi(x) G(t)$. Substituting this into the equation, we obtain that

$$
\frac{G^{\prime}}{G}=\frac{\phi^{\prime \prime}}{\phi}=-\lambda=\text { const. }
$$

Hence $G^{\prime}=-\lambda G$ and $\phi^{\prime \prime}=-\lambda \phi$. The boundary conditions are satisfied provided $\phi(0)=\phi(\pi)=0$.
The eigenvalue problem

$$
\phi^{\prime \prime}=-\lambda \phi, \quad \phi(0)=\phi(\pi)=0
$$

has eigenvalues $\lambda_{n}=n^{2}, n=1,2, \ldots$. The corresponding eigenfunctions are $\phi_{n}(x)=\sin n x$. Further, the general solution of the equation $G^{\prime}=-\lambda G$ is $G(t)=c e^{-\lambda t}$, where $c$ is an arbitrary constant. Therefore the functions $u(x, t)=e^{-n^{2} t} \sin n x, n=1,2, \ldots$, are solutions of the boundary value problem.

We are looking for the solution of the initial-boundary value problem as a superposition of solutions with separated variables:

$$
u(x, t)=\sum_{n=1}^{\infty} b_{n} e^{-n^{2} t} \sin n x .
$$

The initial condition is satisfied if

$$
2 x=\sum_{n=1}^{\infty} b_{n} \sin n x \quad \text { for } 0<x<\pi .
$$

Hence $b_{n}$ are coefficients of the Fourier sine series of the function $f(x)=2 x$ on the interval $[0, \pi]$. It remains to use the solution of Problem 1(i).

Problem 3 ( 30 pts.) Consider the initial value problem for the wave equation

$$
\begin{aligned}
& \frac{\partial^{2} u}{\partial t^{2}}=4 \frac{\partial^{2} u}{\partial x^{2}} \quad(-\infty<x<\infty, \quad-\infty<t<\infty) \\
& u(x, 0)=3 \sin \pi x \quad(-\infty<x<\infty) \\
& \frac{\partial u}{\partial t}(x, 0)=8 \pi \sin 3 \pi x \quad(-\infty<x<\infty)
\end{aligned}
$$

(i) Solve the problem (try to obtain a simple formula).

Let $f(x)=3 \sin \pi x$ and $g(x)=8 \pi \sin 3 \pi x, x \in \mathbb{R}$. According to d'Alembert's formula, the solution is

$$
u(x, t)=\frac{1}{2}(f(x-2 t)+f(x+2 t)+G(x+2 t)-G(x-2 t))
$$

where $G$ is an arbitrary anti-derivative of the function $g / 2$. We can take $G(x)=-\frac{4}{3} \cos 3 \pi x$. Then

$$
\begin{aligned}
u(x, t) & =\frac{1}{2}\left(3 \sin (\pi(x-2 t))+3 \sin (\pi(x+2 t))-\frac{4}{3} \cos (3 \pi(x+2 t))+\frac{4}{3} \cos (3 \pi(x-2 t))\right) \\
& =3 \cos 2 \pi t \cdot \sin \pi x+\frac{4}{3} \sin 6 \pi t \cdot \sin 3 \pi x
\end{aligned}
$$

(ii) Determine which two of the following boundary conditions are satisfied by the solution:

$$
u(0, t)=0, \quad \frac{\partial u}{\partial x}(0, t)=0, \quad u(0.5, t)=0, \quad \frac{\partial u}{\partial x}(0.5, t)=0
$$

The solution satisfies the boundary conditions $u(0, t)=0$ and $\frac{\partial u}{\partial x}(0.5, t)=0$, which can be verified directly. Another way to show this is to observe that the initial data $f(x)=3 \sin \pi x$ and $g(x)=$ $8 \pi \sin 3 \pi x$ are odd around 0 and even around 0.5 , that is, $f(-x)=-f(x), g(-x)=-g(x), f(0.5+x)=$ $f(0.5-x)$, and $g(0.5+x)=g(0.5-x)$ for all $x \in \mathbb{R}$. This implies that the solution $u(x, t)$ is odd around 0 and even around 0.5 as a function of $x$, which, in turn, implies the two boundary conditions.

Bonus Problem 4 ( 35 pts.) Consider the initial-boundary value problem for the wave equation

$$
\begin{aligned}
& \frac{\partial^{2} u}{\partial t^{2}}=\frac{\partial^{2} u}{\partial x^{2}} \quad(x>0, \quad t>0) \\
& u(x, 0)=f(x), \quad \frac{\partial u}{\partial t}(x, 0)=0 \quad(x>0) \\
& \frac{\partial u}{\partial x}(0, t)=0 \quad(t>0)
\end{aligned}
$$

where $f(x)=4 \sin ^{2} \pi x$ for $2 \leq x \leq 3$ and $f(x)=0$ otherwise.
(i) Sketch the solution $u(x, t)$ as a function of $x$ for $t=0, t=1, t=3$, and $t=4$.

The solution $u(x, t)$ can be extended to a solution in the whole plane that satisfies the initial conditions $u(x, 0)=F(x)$ and $\frac{\partial u}{\partial t}(x, 0)=0$, where the function $F: \mathbb{R} \rightarrow \mathbb{R}$ is the even extension of $f$. By d'Alembert's formula, $u(x, t)=\frac{1}{2}(F(x-t)+F(x+t))$. This representation allows one to easily sketch $u(x, t)$ as a function of $x$ for any $t \geq 0$ outside intervals $(0,1 / 2)$ and $(2,3)$.

At $t=0$, there is a single pulse localized on the interval $(2,3)$. It is divided into two identical pulses of the same shape and half the amplitude which start moving in opposite directions at unit speed. At $t=1 / 2$, the two pulses become separated. At $t=2$, the pulse travelling to the left reaches the end of the interval. Between $t=2$ and $t=3$ the pulse is being reflected from the boundary. The reflection is completed by $t=3$ when the pulse assumes the same shape as at $t=2$. Then it starts travelling to the right at unit speed.

## $t=0$ <br> 






(ii) Describe how the answer to (i) would change if we considered the boundary condition $u(0, t)=0$ instead.

In this case, the solution is $u(x, t)=\frac{1}{2}\left(F^{-}(x-t)+F^{-}(x+t)\right)$, where $F^{-}$is the odd extension of the function $f$ to the real line. The two solutions coincide for $0 \leq t \leq 2$. Since $t=2$, the left pulse looks differently. In particular, after reflection from the boundary it turns upside down.


