

Math 412-501

Theory of Partial Differential Equations

Lecture 2-11: Review for Exam 2.

Heat conduction in a rectangle

Initial-boundary value problem:

$$\frac{\partial u}{\partial t} = k \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) \quad (0 < x < L, 0 < y < H),$$

$$u(x, y, 0) = f(x, y) \quad (0 < x < L, 0 < y < H),$$

$$\frac{\partial u}{\partial x}(0, y, t) = \frac{\partial u}{\partial x}(L, y, t) = 0 \quad (0 < y < H),$$

$$\frac{\partial u}{\partial y}(x, 0, t) = \frac{\partial u}{\partial y}(x, H, t) = 0 \quad (0 < x < L).$$

We search for the solution $u(x, y, t)$ as a superposition of solutions with separated variables that satisfy the boundary conditions.

Separation of variables: $u(x, y, t) = \phi(x)h(y)G(t)$.
Substitute this into the heat equation:

$$\phi(x)h(y)\frac{dG}{dt} = k \left(\frac{d^2\phi}{dx^2}h(y)G(t) + \phi(x)\frac{d^2h}{dy^2}G(t) \right).$$

Divide both sides by $k \cdot \phi(x)h(y)G(t) = k \cdot u(x, y, t)$:

$$\frac{1}{kG} \cdot \frac{dG}{dt} = \frac{1}{\phi} \cdot \frac{d^2\phi}{dx^2} + \frac{1}{h} \cdot \frac{d^2h}{dy^2}.$$

It follows that $\frac{1}{\phi} \cdot \frac{d^2\phi}{dx^2} = -\lambda$, $\frac{1}{h} \cdot \frac{d^2h}{dy^2} = -\mu$,

$\frac{1}{kG} \cdot \frac{dG}{dt} = -\lambda - \mu$, where λ and μ are **separation constants**.

The variables have been separated:

$$\begin{aligned}\frac{dG}{dt} &= -(\lambda + \mu)kG, \\ \frac{d^2\phi}{dx^2} &= -\lambda\phi, & \frac{d^2h}{dy^2} &= -\mu h.\end{aligned}$$

Proposition Suppose G , ϕ , and h are solutions of the above ODEs for the same values of λ and μ .

Then $u(x, y, t) = \phi(x)h(y)G(t)$ is a solution of the heat equation.

Boundary conditions $\frac{\partial u}{\partial x}(0, y, t) = \frac{\partial u}{\partial x}(L, y, t) = 0$ hold if $\phi'(0) = \phi'(L) = 0$.

Boundary conditions $\frac{\partial u}{\partial y}(x, 0, t) = \frac{\partial u}{\partial y}(x, H, t) = 0$ hold if $h'(0) = h'(H) = 0$.

1st eigenvalue problem: $\phi'' = -\lambda\phi$, $\phi'(0) = \phi'(L) = 0$.

Eigenvalues: $\lambda_n = \left(\frac{n\pi}{L}\right)^2$, $n = 0, 1, 2, \dots$

Eigenfunctions: $\phi_0 = 1$, $\phi_n(x) = \cos \frac{n\pi x}{L}$, $n \geq 1$.

2nd eigenvalue problem: $h'' = -\mu h$, $h'(0) = h'(H) = 0$.

Eigenvalues: $\mu_m = \left(\frac{m\pi}{H}\right)^2$, $m = 0, 1, 2, \dots$

Eigenfunctions: $h_0 = 1$, $h_m(y) = \cos \frac{m\pi y}{H}$, $m \geq 1$.

Dependence on t :

$$G'(t) = -(\lambda + \mu)kG(t) \implies G(t) = C_0 e^{-(\lambda + \mu)kt}$$

Solution of the boundary value problem:

$$\begin{aligned} u(x, y, t) &= e^{-(\lambda_n + \mu_m)kt} \phi_n(x) h_m(y) \\ &= \exp\left(-\left(\left(\frac{n\pi}{L}\right)^2 + \left(\frac{m\pi}{H}\right)^2\right)kt\right) \cdot \cos \frac{n\pi x}{L} \cdot \cos \frac{m\pi y}{H}. \end{aligned}$$

We are looking for the solution of the initial-boundary value problem as a superposition of solutions with separated variables.

$$\begin{aligned} u(x, y, t) &= \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} e^{-(\lambda_n + \mu_m)kt} \phi_n(x) h_m(y) \\ &= \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} \exp\left(-\left(\left(\frac{n\pi}{L}\right)^2 + \left(\frac{m\pi}{H}\right)^2\right)kt\right) \cdot \cos \frac{n\pi x}{L} \cdot \cos \frac{m\pi y}{H} \end{aligned}$$

How do we find coefficients $C_{n,m}$?

From the initial condition $u(x, y, 0) = f(x, y)$.

$$f(x, y) = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} \cos \frac{n\pi x}{L} \cos \frac{m\pi y}{H}$$

(double Fourier cosine series)

How do we solve the heat conduction problem?

$$\frac{\partial u}{\partial t} = k \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) \quad (0 < x < L, 0 < y < H),$$

$$u(x, y, 0) = f(x, y) \quad (0 < x < L, 0 < y < H),$$

$$\frac{\partial u}{\partial x}(0, y, t) = \frac{\partial u}{\partial x}(L, y, t) = 0 \quad (0 < y < H),$$

$$\frac{\partial u}{\partial y}(x, 0, t) = \frac{\partial u}{\partial y}(x, H, t) = 0 \quad (0 < x < L).$$

- Expand f into the double Fourier cosine series:

$$f(x, y) = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} \cos \frac{n\pi x}{L} \cos \frac{m\pi y}{H}.$$

- Write the solution: $u(x, y, t) =$

$$= \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} \exp\left(-\left(\left(\frac{n\pi}{L}\right)^2 + \left(\frac{m\pi}{H}\right)^2\right)kt\right) \cos \frac{n\pi x}{L} \cos \frac{m\pi y}{H}.$$

How do we expand f into the double Fourier cosine series?

$$f(x, y) = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} \cos \frac{n\pi x}{L} \cos \frac{m\pi y}{H}.$$

If $f(x, y)$ is smooth then the expansion **exists**.

How do we determine coefficients?

Multiply both sides by $\phi_N(x)h_M(y)$ and integrate over the rectangle.

$$\begin{aligned}
& \int_0^L \int_0^H f(x, y) \phi_N(x) h_M(y) dx dy \\
= & \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} \int_0^L \int_0^H \phi_N(x) h_M(y) \phi_n(x) h_m(y) dx dy \\
= & \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} \int_0^L \phi_N(x) \phi_n(x) dx \int_0^H h_M(y) h_m(y) dy.
\end{aligned}$$

We have the orthogonality relations:

$$\begin{aligned}
\int_0^L \phi_N(x) \phi_n(x) dx &= 0, & n \neq N, \\
\int_0^H h_M(y) h_m(y) dy &= 0, & m \neq M.
\end{aligned}$$

Hence

$$\begin{aligned} & \int_0^L \int_0^H f(x, y) \phi_N(x) h_M(y) dx dy \\ &= C_{N,M} \int_0^L \phi_N^2(x) dx \int_0^H h_M^2(y) dy. \end{aligned}$$

It remains to recall that

$$\begin{aligned} \int_0^L \phi_0^2(x) dx &= L, & \int_0^L \phi_N^2(x) dx &= \frac{L}{2}, & N \geq 1, \\ \int_0^H h_0^2(x) dx &= H, & \int_0^H h_M^2(x) dx &= \frac{H}{2}, & M \geq 1. \end{aligned}$$

How do we expand $f(x, y)$ into the double series?

$$f(x, y) = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} C_{n,m} \cos \frac{n\pi x}{L} \cos \frac{m\pi y}{H},$$

where

$$C_{0,0} = \frac{1}{LH} \int_0^L \int_0^H f(x, y) dx dy,$$

$$C_{n,0} = \frac{2}{LH} \int_0^L \int_0^H f(x, y) \cos \frac{n\pi x}{L} dx dy, \quad n \geq 1,$$

$$C_{0,m} = \frac{2}{LH} \int_0^L \int_0^H f(x, y) \cos \frac{m\pi y}{H} dx dy, \quad m \geq 1,$$

$$C_{n,m} = \frac{4}{LH} \int_0^L \int_0^H f(x, y) \cos \frac{n\pi x}{L} \cos \frac{m\pi y}{H} dx dy.$$

Laplace's equation in a semicircle

Solve Laplace's equation inside a semicircle of radius a ($0 < r < a$, $0 < \theta < \pi$) subject to the boundary conditions: $u = 0$ on the diameter and $u(a, \theta) = g(\theta)$.

Laplace's equation in polar coordinates (r, θ) :

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = 0.$$

$u = 0$ on the diameter \implies

$$u(r, 0) = u(r, \pi) = 0 \quad (0 < r < a),$$

$$u(0, \theta) = 0 \quad (0 < \theta < \pi).$$

Separation of variables: $u(r, \theta) = h(r)\phi(\theta)$.

Substitute this into Laplace's equation:

$$\frac{d^2 h}{dr^2} \phi(\theta) + \frac{1}{r} \frac{dh}{dr} \phi(\theta) + \frac{1}{r^2} h(r) \frac{d^2 \phi}{d\theta^2} = 0.$$

Divide both sides by $r^{-2}h(r)\phi(\theta) = r^{-2}u(r, \theta)$:

$$\frac{1}{h} \cdot \left(r^2 \frac{d^2 h}{dr^2} + r \frac{dh}{dr} \right) = -\frac{1}{\phi} \cdot \frac{d^2 \phi}{d\theta^2}.$$

It follows that

$$\frac{1}{h} \cdot \left(r^2 \frac{d^2 h}{dr^2} + r \frac{dh}{dr} \right) = -\frac{1}{\phi} \cdot \frac{d^2 \phi}{d\theta^2} = \lambda = \text{const.}$$

The variables have been separated:

$$r^2 \frac{d^2 h}{dr^2} + r \frac{dh}{dr} = \lambda h, \quad \frac{d^2 \phi}{d\theta^2} = -\lambda \phi.$$

Proposition Suppose h and ϕ are solutions of the above ODEs for the same value of λ . Then $u(r, \theta) = h(r)\phi(\theta)$ is a solution of Laplace's equation.

Boundary conditions $u(r, 0) = u(r, \pi) = 0$ hold if

$$\phi(0) = \phi(\pi) = 0.$$

Boundary condition $u(0, \theta) = 0$ holds if

$$h(0) = 0.$$

Euler's (or equidimensional) equation

$$r^2 \frac{d^2 h}{dr^2} + r \frac{dh}{dr} - \lambda h = 0 \quad (r > 0)$$

$$\lambda > 0 \implies h(r) = C_1 r^p + C_2 r^{-p} \quad (\lambda = p^2, p > 0)$$

$$\lambda = 0 \implies h(r) = C_1 + C_2 \log r$$

Eigenvalue problem: $\phi'' = -\lambda\phi$, $\phi(0) = \phi(\pi) = 0$.

Eigenvalues: $\lambda_n = n^2$, $n = 1, 2, \dots$

Eigenfunctions: $\phi_n(\theta) = \sin n\theta$.

Dependence on r :

$$r^2 h'' + rh' = \lambda h, \quad h(0) = 0.$$

$$\implies h(r) = C_0 r^p \quad (p = \sqrt{\lambda})$$

Solution of Laplace's equation:

$$u(r, \theta) = r^n \sin n\theta, \quad n = 1, 2, \dots$$

A superposition of the solutions with separated variables is a series

$$u(r, \theta) = \sum_{n=1}^{\infty} c_n r^n \sin n\theta,$$

where c_1, c_2, \dots are constants.

Substituting the series into the boundary condition $u(a, \theta) = g(\theta)$, we get

$$g(\theta) = \sum_{n=1}^{\infty} c_n a^n \sin n\theta.$$

Hence $c_n = b_n a^{-n}$, $n = 1, 2, \dots$, where

$$g(\theta) = \sum_{n=1}^{\infty} b_n \sin n\theta.$$

Solution.

$$u(r, \theta) = \sum_{n=1}^{\infty} b_n \left(\frac{r}{a}\right)^n \sin n\theta,$$

where

$$\sum_{n=1}^{\infty} b_n \sin n\theta$$

is the Fourier sine series of the function $g(\theta)$ on $[0, \pi]$, that is,

$$b_n = \frac{2}{\pi} \int_0^{\pi} g(\theta) \sin n\theta \, d\theta, \quad n = 1, 2, \dots$$