Math 412-501 Theory of Partial Differential Equations Lecture 2-7: Sturm-Liouville eigenvalue problems.

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Sturm-Liouville differential equation:

$$\frac{d}{dx} \left(p \frac{d\phi}{dx} \right) + q\phi + \lambda \sigma \phi = 0 \quad (a < x < b),$$

where p = p(x), q = q(x), $\sigma = \sigma(x)$ are known functions on [a, b] and λ is an unknown constant.

The Sturm-Liouville equation is a linear homogeneous ODE of the second order.

Sturm-Liouville eigenvalue problem =

- = Sturm-Liouville differential equation +
- + linear homogeneous boundary conditions





J. C. F. Sturm (1803–1855) J. Liouville (1809–1882)

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The Sturm-Liouville equation usually arises after separation of variables in a linear homogeneous PDE of the second order.

Examples.

• $\phi'' + \lambda \phi = 0$ (heat, wave, Laplace's equations) • $r^2 \frac{d^2 h}{dr^2} + r \frac{dh}{dr} = \lambda h$ (Laplace's equation in polar coordinates) standard notation: $x^2 \phi'' + x \phi' - \lambda \phi = 0$ canonical form: $(x \phi')' - \lambda x^{-1} \phi = 0$

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Heat flow in a nonuniform rod:

$$c\rho \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(K_0 \frac{\partial u}{\partial x} \right) + Q,$$

$$K_0 = K_0(x), \ c = c(x), \ \rho = \rho(x), \ Q = Q(u, x, t).$$

The equation is linear homogeneous if $Q = \alpha(x, t)u$. We assume that $\alpha = \alpha(x)$.

$$c\rho\frac{\partial u}{\partial t} = \frac{\partial}{\partial x}\left(K_0\frac{\partial u}{\partial x}\right) + \alpha u$$

Separation of variables: $u(x, t) = \phi(x)G(t)$. Substitute this into the heat equation:

$$c\rho\phi\frac{dG}{dt}=\frac{d}{dx}\Big(K_0\frac{d\phi}{dx}\Big)G+\alpha\phi G.$$

Divide both sides by $c(x)\rho(x)\phi(x)G(t) = c\rho u$:

$$\frac{1}{G}\frac{dG}{dt} = \frac{1}{c\rho\phi}\frac{d}{dx}\left(K_0\frac{d\phi}{dx}\right) + \frac{\alpha}{c\rho} = -\lambda = \text{const.}$$

The variables have been separated:

$$\frac{dG}{dt} + \lambda G = 0,$$

$$\frac{d}{dx}\left(K_0\frac{d\phi}{dx}\right) + \alpha\phi + \lambda c\rho\phi = 0.$$

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Sturm-Liouville differential equation:

$$rac{d}{dx} \Big(p rac{d\phi}{dx} \Big) + q\phi + \lambda \sigma \phi = 0 \quad (a < x < b).$$

Examples of boundary conditions:

- $\phi(a) = \phi(b) = 0$ (Dirichlet conditions)
- $\phi'(a) = \phi'(b) = 0$ (von Neumann conditions)
- $\phi'(a) = 2\phi(a), \ \phi'(b) = -3\phi(b)$ (Robin conditions)
- $\phi(a) = 0$, $\phi'(b) = 0$ (mixed conditions) • $\phi(a) = \phi(b)$, $\phi'(a) = \phi'(b)$ (periodic conditions)
 - $|\phi(a)| < \infty$, $\phi(b) = 0$ (singular conditions)

$$\frac{d}{dx}\left(p\frac{d\phi}{dx}\right) + q\phi + \lambda\sigma\phi = 0 \quad (a < x < b).$$

The equation is **regular** if p, q, σ are real and continuous on [a, b], and $p, \sigma > 0$ on [a, b].

The Sturm-Liouville eigenvalue problem is **regular** if the equation is regular and boundary conditions are of the form

$$egin{aligned} η_1\phi(a)+eta_2\phi'(a)=0,\ η_3\phi(b)+eta_4\phi'(b)=0, \end{aligned}$$

where $\beta_i \in \mathbb{R}$, $|\beta_1| + |\beta_2| \neq 0$, $|\beta_3| + |\beta_4| \neq 0$.

This includes Dirichlet, Neumann, and Robin conditions but excludes periodic and singular ones.

Regular Sturm-Liouville eigenvalue problem:

$$\begin{aligned} \frac{d}{dx} \left(p \frac{d\phi}{dx} \right) + q\phi + \lambda \sigma \phi &= 0 \quad (a < x < b), \\ \beta_1 \phi(a) + \beta_2 \phi'(a) &= 0, \\ \beta_3 \phi(b) + \beta_4 \phi'(b) &= 0. \end{aligned}$$

Eigenfunction: nonzero solution ϕ of the boundary value problem.

Eigenvalue: corresponding value of λ .

Eigenvalues and eigenfunctions of a regular Sturm-Liouville eigenvalue problem have **six** important properties. **Property 1.** All eigenvalues are real.

Property 2. All eigenvalues can be arranged in the ascending order

$$\lambda_1 < \lambda_2 < \ldots < \lambda_n < \lambda_{n+1} < \ldots$$

so that $\lambda_n \to \infty$ as $n \to \infty$.

This means that:

- there are infinitely many eigenvalues;
- there is a smallest eigenvalue;
- on any finite interval, there are only finitely many eigenvalues.

Remark. It is possible that $\lambda_1 < 0$.

Property 3. Given an eigenvalue λ_n , the corresponding eigenfunction ϕ_n is unique up to a multiplicative constant. The function ϕ_n has exactly n-1 zeros in (a, b).

We say that λ_n is a **simple** eigenvalue.

Property 4. Eigenfunctions belonging to different eigenvalues satisfy an integral identity:

$$\int_a^b \phi_n(x)\phi_m(x)\sigma(x)\,dx=0 \quad \text{ if } \quad \lambda_n\neq\lambda_m.$$

We say that ϕ_n and ϕ_m are **orthogonal** relative to the weight function σ .

Property 5. Any eigenvalue λ can be related to its eigenfunction ϕ as follows:

$$\lambda = \frac{-p\phi\phi' \Big|_a^b + \int_a^b (p(\phi')^2 - q\phi^2) dx}{\int_a^b \phi^2 \sigma dx}.$$

The right-hand side is called the **Rayleigh quotient**.

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Property 6. Any piecewise continuous function $f : [a, b] \to \mathbb{R}$ is assigned a series $f(x) \sim \sum_{n=1}^{\infty} c_n \phi_n(x)$,

where

$$c_n = \frac{\int_a^b f(x)\phi_n(x)\sigma(x)\,dx}{\int_a^b \phi_n^2(x)\sigma(x)\,dx}.$$

If f is piecewise smooth then the series converges for any a < x < b. The sum is equal to f(x) if f is continuous at x. Otherwise the series converges to $\frac{1}{2}(f(x+) + f(x-))$.

We say that the set of eigenfunctions ϕ_n is **complete**.

A regular Sturm-Liouville eigenvalue problem:

$$\phi'' + \lambda \phi = 0, \quad \phi(0) = \phi(L) = 0.$$

($p = \sigma = 1, q = 0, [a, b] = [0, L]$)
Eigenvalues: $\lambda_n = (\frac{n\pi}{L})^2, n = 1, 2, ...$
Eigenfunctions: $\phi_n(x) = \sin \frac{n\pi x}{L}$.

The zeros of ϕ_n divide the interval [0, L] into n equal parts.

Property 3a. Suppose $x_1 < x_2 < \ldots < x_{n-1}$ are zeros of the eigenfunction ϕ_n in (a, b). Then ϕ_{n+1} has exactly one zero in each of the following intervals: $(a, x_1), (x_1, x_2), (x_2, x_3), \ldots, (x_{n-2}, x_{n-1}), (x_{n-1}, b)$.



Eigenfunctions ϕ_n

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Orthogonality:
$$\int_{0}^{L} \sin \frac{n\pi x}{L} \sin \frac{m\pi x}{L} dx = 0, \quad n \neq m.$$

Rayleigh quotient:
$$\lambda = \frac{\int_{0}^{L} |\phi'(x)|^{2} dx}{\int_{0}^{L} |\phi(x)|^{2} dx}.$$

Fourier sine series:
$$f \sim \sum_{n=1}^{\infty} c_n \sin \frac{n\pi x}{L}$$
,

where
$$c_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$
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Note that
$$\int_0^L \left(\sin \frac{n\pi x}{L}\right)^2 dx = \frac{L}{2}.$$