

Math 412-501

Theory of Partial Differential Equations

**Lecture 3-6: Nonhomogeneous problems.  
Method of eigenfunction expansion.**

## Nonhomogeneous boundary conditions

Initial-boundary value problem for the heat equation:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} \quad (0 < x < L),$$

$$u(0, t) = A,$$

$$u(L, t) = B,$$

$$u(x, 0) = f(x) \quad (0 < x < L).$$

Our goal is to reduce this problem to an analogous one with homogeneous boundary conditions.

Suppose  $u_0(x, t)$  is a solution of the equation that satisfies the boundary conditions.

Consider  $w(x, t) = u(x, t) - u_0(x, t)$ .

$u$  and  $u_0$  are solutions of the heat equation

$\implies w = u - u_0$  is also a solution

$u$  and  $u_0$  satisfy the same boundary conditions

$\implies w = u - u_0$  satisfies homogeneous BCs

$w(x, t)$  is a solution of the following problem:

$$\frac{\partial w}{\partial t} = k \frac{\partial^2 w}{\partial x^2} \quad (0 < x < L),$$

$$w(0, t) = w(L, t) = 0,$$

$$w(x, 0) = g(x), \quad \text{where } g(x) = f(x) - u_0(x, 0).$$

*How do we find a solution  $u_0(x, t)$ ?*

**Steady-state** (or **equilibrium**) solution:

$$u_0(x, t) = u_0(x).$$

$u_0(x)$  is a solution of the boundary value problem

$$\frac{d^2 u_0}{dx^2} = 0, \quad u_0(0) = A, \quad u_0(L) = B.$$

The general solution of the equation is

$$u_0(x) = c_1 x + c_2, \quad \text{where } c_1, c_2 \text{ are constants.}$$

Boundary conditions are satisfied if

$$u_0(x) = A + \frac{B - A}{L} x.$$

## Initial-boundary value problem:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} \quad (0 < x < L),$$

$$u(0, t) = A, \quad u(L, t) = B,$$

$$u(x, 0) = f(x) \quad (0 < x < L).$$

**Solution:**  $u(x, t) = u_0(x) + w(x, t)$

$$= A + \frac{B - A}{L} x + \sum_{n=1}^{\infty} b_n \exp\left(- (n\pi/L)^2 kt\right) \sin \frac{n\pi x}{L},$$

where  $b_n$  are coefficients of the Fourier sine series on  $[0, L]$  of the function  $g(x) = f(x) - A - \frac{B-A}{L}x$ .

Heat equation with sources:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} + Q(x) \quad (0 < x < L),$$

$$u(0, t) = A, \quad u(L, t) = B,$$

$$u(x, 0) = f(x) \quad (0 < x < L).$$

Suppose  $u_0(x, t) = u_0(x)$  is an equilibrium solution.

Then  $w(x, t) = u(x, t) - u_0(x)$  is a solution of the homogeneous heat equation with homogeneous boundary conditions.

**Solution:**  $u(x, t) = u_0(x) + w(x, t)$ ,

where  $u_0$  is the solution of the boundary value problem

$$k \frac{d^2 u_0}{dx^2} + Q(x) = 0, \quad u_0(0) = A, \quad u_0(L) = B;$$

and  $w(x, t)$  is the solution of the initial-boundary value problem

$$\frac{\partial w}{\partial t} = k \frac{\partial^2 w}{\partial x^2} \quad (0 < x < L),$$

$$w(0, t) = w(L, t) = 0,$$

$$w(x, 0) = f(x) - u_0(x).$$

For some boundary conditions, there is no equilibrium solution. Besides, there is no equilibrium solution if the boundary conditions or the sources depend on time.

Time-dependent problem:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} + Q(x, t) \quad (0 < x < L),$$

$$u(0, t) = A(t), \quad u(L, t) = B(t) \quad (t > 0),$$

$$u(x, 0) = f(x) \quad (0 < x < L).$$

We can still reduce the problem to an analogous one with homogeneous boundary conditions.



Suppose  $u_0(x, t)$  is a smooth function satisfying the boundary conditions (not necessarily a solution of the PDE). For example,

$$u_0(x, t) = A(t) + \frac{B(t) - A(t)}{L} x.$$

Then  $w(x, t) = u(x, t) - u_0(x, t)$  satisfies homogeneous boundary conditions.

Substitute  $u = u_0 + w$  into the equation:

$$\frac{\partial(u_0 + w)}{\partial t} = k \frac{\partial^2(u_0 + w)}{\partial x^2} + Q(x, t),$$

$$\frac{\partial w}{\partial t} = k \frac{\partial^2 w}{\partial x^2} + \left( Q(x, t) - \frac{\partial u_0}{\partial t} + k \frac{\partial^2 u_0}{\partial x^2} \right).$$

Also,  $w(x, 0) = f(x) - u_0(x, 0)$ .

**Solution:**  $u(x, t) = u_0(x, t) + w(x, t)$ ,

where

$$u_0(x, t) = A(t) + \frac{B(t) - A(t)}{L} x$$

and  $w(x, t)$  is the solution of the initial-boundary value problem

$$\frac{\partial w}{\partial t} = k \frac{\partial^2 w}{\partial x^2} + \tilde{Q}(x, t) \quad (0 < x < L),$$

$$w(0, t) = w(L, t) = 0,$$

$$w(x, 0) = g(x) \quad (0 < x < L),$$

in which  $\tilde{Q}(x, t) = Q(x, t) - \frac{\partial u_0}{\partial t} + k \frac{\partial^2 u_0}{\partial x^2}$ ,  
 $g(x) = f(x) - u_0(x, 0)$ .

## Method of eigenfunction expansion

Initial-boundary value problem:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} + Q(x, t) \quad (0 < x < L),$$

$$u(0, t) = u(L, t) = 0,$$

$$u(x, 0) = f(x) \quad (0 < x < L).$$

Consider the related eigenvalue problem

$$\phi'' = -\lambda\phi, \quad \phi(0) = \phi(L) = 0.$$

The eigenvalues are  $\lambda_n = \left(\frac{n\pi}{L}\right)^2$ ,  $n = 1, 2, \dots$ ,  
and the corresponding eigenfunctions are

$$\phi_n(x) = \sin \frac{n\pi x}{L}.$$

Any piecewise smooth function on  $[0, L]$  can be expanded into the Fourier sine series. In particular,

$$u(x, t) = \sum_{n=1}^{\infty} a_n(t) \phi_n(x),$$

where  $a_1(t), a_2(t), \dots$  are some functions.

Let us assume that the series can be differentiated term-by-term (one can show this; an important reason is that  $u(x, t)$  satisfies homogeneous boundary conditions). Then

$$\frac{\partial u}{\partial t} = \sum_{n=1}^{\infty} a'_n(t) \phi_n(x),$$

$$\frac{\partial^2 u}{\partial x^2} = \sum_{n=1}^{\infty} a_n(t) \phi_n''(x) = - \sum_{n=1}^{\infty} \lambda_n a_n(t) \phi_n(x).$$

Also, 
$$Q(x, t) = \sum_{n=1}^{\infty} q_n(t)\phi_n(x),$$

$$f(x) = \sum_{n=1}^{\infty} b_n\phi_n(x).$$

Substitute all series into the equation:

$$\begin{aligned} \sum_{n=1}^{\infty} a'_n(t)\phi_n(x) &= \\ &= -k \sum_{n=1}^{\infty} \lambda_n a_n(t)\phi_n(x) + \sum_{n=1}^{\infty} q_n(t)\phi_n(x), \end{aligned}$$

$$\sum_{n=1}^{\infty} (a'_n(t) + k\lambda_n a_n(t) - q_n(t))\phi_n(x) = 0.$$

Initial condition is satisfied if

$$\sum_{n=1}^{\infty} a_n(0)\phi_n(x) = \sum_{n=1}^{\infty} b_n\phi_n(x).$$

It follows that for any  $n \geq 1$  we have

$$a'_n(t) + k\lambda_n a_n(t) = q_n(t), \quad a_n(0) = b_n.$$

This is an initial value problem for  $a_n$ .

Multiply both sides of the ODE by  $e^{\lambda_n kt}$ :

$$e^{\lambda_n kt} a'_n(t) + k\lambda_n e^{\lambda_n kt} a_n(t) = e^{\lambda_n kt} q_n(t),$$

$$(e^{\lambda_n kt} a_n(t))' = e^{\lambda_n kt} q_n(t),$$

$$e^{\lambda_n kt} a_n(t) - a_n(0) = \int_0^t e^{\lambda_n k\tau} q_n(\tau) d\tau,$$

$$a_n(t) = b_n e^{-\lambda_n kt} + e^{-\lambda_n kt} \int_0^t e^{\lambda_n k\tau} q_n(\tau) d\tau.$$

**Solution:**  $u(x, t) = \sum_{n=1}^{\infty} a_n(t) \sin \frac{n\pi x}{L},$

where

$$a_n(t) = b_n e^{-\lambda_n k t} + e^{-\lambda_n k t} \int_0^t e^{\lambda_n k \tau} q_n(\tau) d\tau,$$

$$\lambda_n = (n\pi/L)^2,$$

$$b_n = \frac{\langle f, \phi_n \rangle}{\langle \phi_n, \phi_n \rangle} = \frac{2}{L} \int_0^L f(\xi) \sin \frac{n\pi \xi}{L} d\xi,$$

$$q_n(\tau) = \frac{\langle Q(\cdot, \tau), \phi_n \rangle}{\langle \phi_n, \phi_n \rangle} = \frac{2}{L} \int_0^L Q(\xi, \tau) \sin \frac{n\pi \xi}{L} d\xi.$$

If  $Q = 0$  then  $u(x, t) = \sum_{n=1}^{\infty} b_n e^{-\lambda_n k t} \phi_n(x),$   
just as the separation of variables yields.

## Poisson's equation

Poisson's equation is a nonhomogeneous version of Laplace's equation:  $\nabla^2 u = Q$ .

Consider a boundary value problem:

$$\nabla^2 u = Q \quad \text{in the domain } D,$$

$$u|_{\partial D} = \alpha,$$

where  $Q$  is a function on  $D$  and  $\alpha$  is a function on the boundary  $\partial D$ .

There are two ways to solve this problem.



**Method 1.** Suppose  $u_0$  is a smooth function in the domain  $D$  such that  $\nabla^2 u_0 = Q$ .

Then  $u = u_0 + w$ , where  $w$  is the solution of the boundary value problem for Laplace's equation:

$$\begin{aligned}\nabla^2 u &= 0 && \text{in the domain } D, \\ u|_{\partial D} &= \beta,\end{aligned}$$

where  $\beta(\mathbf{x}) = \alpha(\mathbf{x}) - u_0(\mathbf{x})$  for any  $\mathbf{x} \in \partial D$ .

We know how to solve the latter problem if  $D$  is a rectangle or a circle. Method 1 applies only to certain functions  $Q$  as it relies on a lucky guess.

*Example:*  $\nabla^2 u = Q(x, y) = e^{-x} \cos 2y - \sin 3x \sin y$ .

Lucky guess:  $u_0(x, y) = -\frac{1}{3}e^{-x} \cos 2y + \frac{1}{10} \sin 3x \sin y$ .

**Method 2.**  $u = u_1 + u_2$ ,

where  $u_1$  is the solution of the problem

$$\begin{aligned}\nabla^2 u_1 &= Q && \text{in the domain } D, \\ u_1|_{\partial D} &= 0,\end{aligned}$$

while  $u_2$  is the solution of the problem

$$\begin{aligned}\nabla^2 u_2 &= 0 && \text{in the domain } D, \\ u_2|_{\partial D} &= \alpha.\end{aligned}$$

We know how to solve the latter problem if  $D$  is a rectangle or a circle. The former problem is solved by method of eigenfunction expansion.

Boundary value problem:

$$\begin{aligned}\nabla^2 u_1 &= Q && \text{in the domain } D, \\ u_1|_{\partial D} &= 0.\end{aligned}$$

Let  $\lambda_1 < \lambda_2 \leq \dots$  be eigenvalues of the negative Dirichlet Laplacian in  $D$  (counting with multiplicities), and  $\phi_1, \phi_2, \dots$  be the corresponding (orthogonal) eigenfunctions.

We have that

$$u_1(x, y) = \sum_{n=1}^{\infty} c_n \phi_n(x, y),$$

$$Q(x, y) = \sum_{n=1}^{\infty} q_n \phi_n(x, y).$$

The Laplacian  $\nabla^2 u_1$  can be evaluated term-by-term (it follows from the fact that  $u_1|_{\partial D} = 0$ ). Hence

$$\nabla^2 u_1 = \sum_{n=1}^{\infty} c_n \nabla^2 \phi_n = - \sum_{n=1}^{\infty} \lambda_n c_n \phi_n.$$

Thus  $-\lambda_n c_n = q_n$  for  $n = 1, 2, \dots$

**Solution:**

$$u_1 = - \sum_{n=1}^{\infty} \lambda_n^{-1} q_n \phi_n, \quad \text{where} \quad q_n = \frac{\langle Q, \phi_n \rangle}{\langle \phi_n, \phi_n \rangle}.$$