Math 412-501 Theory of Partial Differential Equations Lecture 4-1: Green's functions. Dirac delta function.

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Heat equation on the infinite interval

Initial value problem:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} \qquad (-\infty < x < \infty),$$

$$u(x,0) = f(x).$$
Solution:
$$u(x,t) = \int_{-\infty}^{\infty} G(x,\tilde{x},t) f(\tilde{x}) d\tilde{x},$$
where
$$G(x,\tilde{x},t) = \frac{1}{\sqrt{4\pi kt}} e^{-\frac{(x-\tilde{x})^2}{4kt}}.$$

The solution is in the integral operator form. The function G is called the **kernel** of the operator. Also, $G(x, \tilde{x}, t)$ is called **Green's function** of the problem.

$$G(x, ilde{x},t)=rac{1}{\sqrt{4\pi kt}}\,e^{-rac{(x- ilde{x})^2}{4kt}}$$

 $G(x, \tilde{x}, t)$ is:

- defined for t > 0,
- positive,
- infinitely smooth.

$$\int_{-\infty}^{\infty} G(x, \tilde{x}, t) \, dx = \int_{-\infty}^{\infty} \frac{1}{\sqrt{4\pi kt}} \, e^{-\frac{(x-\tilde{x})^2}{4kt}} \, dx$$
$$= \int_{-\infty}^{\infty} \frac{1}{\sqrt{4\pi kt}} \, e^{-\frac{x^2}{4kt}} \, dx = \int_{-\infty}^{\infty} \frac{1}{\sqrt{\pi}} \, e^{-y^2} \, dy = 1$$

$$\begin{aligned} \frac{\partial G(x,\tilde{x},t)}{\partial x} &= \frac{1}{\sqrt{4\pi kt}} e^{-\frac{(x-\tilde{x})^2}{4kt}} \cdot \frac{-2(x-\tilde{x})}{4kt} \\ &= -\frac{x-\tilde{x}}{2kt} G(x,\tilde{x},t); \\ \frac{\partial^2 G(x,\tilde{x},t)}{\partial x^2} &= \left(\frac{(x-\tilde{x})^2}{4k^2t^2} - \frac{1}{2kt}\right) G(x,\tilde{x},t); \\ \frac{\partial G(x,\tilde{x},t)}{\partial t} &= -\frac{1}{2}t^{-3/2} \cdot \frac{1}{\sqrt{4\pi k}} e^{-\frac{(x-\tilde{x})^2}{4kt}} \\ &+ \frac{1}{\sqrt{4\pi kt}} e^{-\frac{(x-\tilde{x})^2}{4kt}} \cdot \frac{(x-\tilde{x})^2}{4kt^2} &= \left(\frac{(x-\tilde{x})^2}{4kt^2} - \frac{1}{2t}\right) G(x,\tilde{x},t). \end{aligned}$$

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Lemma For any $\tilde{x} \in \mathbb{R}$ the function $u(x, t) = G(x, \tilde{x}, t)$ is a solution of the heat equation for t > 0.

Besides,
$$\lim_{t\to 0} G(x, \tilde{x}, t) = \begin{cases} \infty & \text{if } x = \tilde{x}, \\ 0 & \text{if } x \neq \tilde{x}. \end{cases}$$

We say that $G(x, \tilde{x}, t)$ is the solution of the initial value problem

$$rac{\partial u}{\partial t} = k rac{\partial^2 u}{\partial x^2} \qquad (-\infty < x < \infty),$$

 $u(x,0) = \delta(x - \tilde{x}),$

where $\delta(x)$ is the **Dirac delta function**.

Dirac delta function

$\delta(x)$ is a function on $\mathbb R$ such that

•
$$\delta(x) = 0$$
 for all $x \neq 0$,

•
$$\delta(0) = \infty$$
,

•
$$\int_{-\infty}^{\infty} \delta(x) \, dx = 1.$$

For any continuous function f and any $x_0 \in \mathbb{R}$, $\int_{-\infty}^{\infty} f(x)\delta(x-x_0) \, dx = f(x_0).$

 $\delta(x)$ is a generalized function (or distribution). That is, δ is a linear functional on a space of test functions f such that $\delta[f] = f(0)$.

Delta sequence/family

Regular function g can be regarded as a generalized function f^{∞}

$$f\mapsto \int_{-\infty}^{\infty}f(x)g(x)\,dx.$$

A **delta sequence** is a sequence of (sharply peaked) functions g_1, g_2, \ldots such that

$$\lim_{n\to\infty}\int_{-\infty}^{\infty}f(x)g_n(x)\,dx=f(0)$$

for any test function f (e.g., infinitely smooth and rapidly decaying). That is, $g_n \to \delta$ as $n \to \infty$ (as generalized functions).

A **delta family** is a family of functions h_{ε} , $0 < \varepsilon \leq \varepsilon_0$, such that $\lim_{\varepsilon \to 0} h_{\varepsilon} = \delta$.



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Initial value problem:

$$\begin{aligned} \frac{\partial u}{\partial t} &= k \frac{\partial^2 u}{\partial x^2} \qquad (-\infty < x < \infty), \\ u(x,0) &= \delta(x - \tilde{x}). \end{aligned}$$

The initial condition means that $g_t(x) = u(x + \tilde{x}, t)$ ought to be a delta family. Indeed, if $u(x, t) = G(x, \tilde{x}, t)$, then

$$g_t(x)=\frac{1}{\sqrt{4\pi kt}}\,e^{-x^2/(4kt)}$$

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Now suppose that

$$u(x,t) = \int_{-\infty}^{\infty} G(x,\tilde{x},t) f(\tilde{x}) d\tilde{x},$$

where f is a test function (smooth and rapidly decaying).

Then u(x, t) is infinitely smooth for t > 0 and solves the heat equation. Besides,

$$\lim_{t\to 0} u(x,t) = f(x).$$

Actually, it is sufficient that f be continuous and bounded.

Heat equation on a semi-infinite interval

Initial-boundary value problem:

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} \qquad (0 < x < \infty),$$
$$\frac{\partial u}{\partial x}(0, t) = 0,$$
$$u(x, 0) = f(x).$$
Solution:
$$u(x, t) = \int_0^\infty G_1(x, \tilde{x}, t) f(\tilde{x}) d\tilde{x},$$

where $G_1(x,\tilde{x},t) = \frac{1}{\sqrt{4\pi kt}} \Big(e^{-\frac{(x-x)}{4kt}} + e^{-\frac{(x+x)}{4kt}} \Big).$

Clearly, $G_1(x, \tilde{x}, t) = G(x, \tilde{x}, t) + G(-x, \tilde{x}, t)$. Since $G(x, \tilde{x}, t)$ is a solution of the heat equation for t > 0, so are $G(-x, \tilde{x}, t)$ and $G_1(x, \tilde{x}, t)$. By definition, $G_1(x, \tilde{x}, t)$ is even as a function of x. Therefore

$$\frac{\partial G_1}{\partial x}(0,\tilde{x},t)=0.$$

Let f_1 denote the even extension of the function f to \mathbb{R} , i.e., $f_1(x) = f_1(-x) = f(x)$ for all $x \ge 0$. Then $u(x,t) = \int_0^\infty G_1(x,\tilde{x},t)f(\tilde{x}) d\tilde{x} = \int_{-\infty}^\infty G(x,\tilde{x},t)f_1(\tilde{x}) d\tilde{x}.$

It follows that u(x, t) is indeed the solution of the initial-boundary value problem. Again, it is sufficient that f be continuous and bounded.

Heat equation on a finite interval

Initial-boundary value problem:

$$\begin{aligned} \frac{\partial u}{\partial t} &= k \frac{\partial^2 u}{\partial x^2} \qquad (0 < x < L), \\ u(0, t) &= u(L, t) = 0, \\ u(x, 0) &= f(x) \qquad (0 < x < L). \end{aligned}$$

Solution:
$$u(x, t) = \sum_{n=1}^{\infty} b_n e^{-\lambda_n k t} \sin \frac{n \pi x}{L}$$

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where $\lambda_n = (n\pi/L)^2$, $b_n = \frac{2}{L} \int_0^L f(\tilde{x}) \sin \frac{n\pi \tilde{x}}{L} d\tilde{x}.$

For
$$t > 0$$
 we obtain

$$u(x, t) = \frac{2}{L} \sum_{n=1}^{\infty} e^{-\lambda_n k t} \sin \frac{n \pi x}{L} \int_0^L f(\tilde{x}) \sin \frac{n \pi \tilde{x}}{L} d\tilde{x}$$

$$= \int_0^L G_2(x, \tilde{x}, t) f(\tilde{x}) d\tilde{x},$$

where

$$G_2(x, \tilde{x}, t) = \frac{2}{L} \sum_{n=1}^{\infty} e^{-\lambda_n k t} \sin \frac{n \pi x}{L} \sin \frac{n \pi \tilde{x}}{L}$$

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